# Non-perturbative studies of gluons and gluinos on the lattice

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# $\mathcal{N} = 1$ supersymmetric Yang-Mills

#### Why study supersymmetric gauge theories on the lattice?

- Extensions of the Standard Model of particle physics: Non-perturbative SUSY effects important to introduce SUSY breaking at low energies.
- Theoretical concepts: SUSY has led to very powerful analytical approaches to understand strong interactions, which need to be complemented and extended by numerical methods.

#### Specific for $\mathcal{N} = 1$ supersymmetric Yang-Mills theory

- gauge sector of SUSY extension of Standard Model
- simplest model with SUSY and local gauge invariance
- Orientifold planar equivalence: SUSY Yang-Mills theory with  $N_c$  colours is equivalent to QCD with a single quark flavour,  $N_f = 1$  QCD, in the limit  $N_c \to \infty$ with Quarks in antisymmetric repr. of  $SU(N_c)$ .

### Non-perturbative Problems

- Spontaneous breaking of chiral symmetry  $Z_{2N_c} \rightarrow Z_2$  $\longleftrightarrow \text{Gluino condensate} \quad <\lambda\lambda > \neq 0$
- Spectrum of bound states  $\rightarrow$  Supermultiplets
- Confinement of static quarks
- Breaking of SUSY in the continuum limit?
- SUSY restauration on the lattice

• Check predictions from effective Lagrangeans (Veneziano, Yankielowicz, ...)

Extrapolation to the supersymmetric continuum limit of SU(3) SUSY Yang-Mills theory



• continuity to semiclassical regime

#### Vector supermultiplet:

• Gauge field  $A_{\mu}^{a}(x)$ ,  $a = 1, \ldots, N_{c}^{2} - 1$ , "Gluon" Gauge group SU( $N_{c}$ ) • Majorana-spinor field  $\lambda^a(x)$ ,  $\overline{\lambda} = \lambda^T C$ , "Gluino" adjoint representation:  $\mathcal{D}_{\mu}\lambda^{a} = \partial_{\mu}\lambda^{a} + g f_{abc}A^{b}_{\mu}\lambda^{c}$ • (auxiliary field  $D^a(x)$ ) Lagrangean:

 $\mathcal{L} = \int d^2 heta \operatorname{Tr}(W^A W_A) + \text{h. c.} = rac{1}{4} F^{\ a}_{\mu
u} F^{\ a}_{\mu
u} + rac{1}{2} \overline{\lambda}^a \gamma_\mu (\mathcal{D}_\mu \lambda)^a + rac{1}{2} D^a D^a$ 

• SUSY: (on-shell)  $\delta A^a_\mu = -2i\overline{\lambda}^a \gamma_\mu \varepsilon$ ,  $\delta \lambda^a = -\sigma_{\mu\nu} F^a_{\mu\nu} \varepsilon$ • Gluino mass term  $m_{\tilde{q}} \overline{\lambda}^a \lambda^a$  breaks SUSY softly. • Differences to QCD:

> $\lambda$ : 1.) Majorana, " $N_f = \frac{1}{2}$ " 2.) adjoint representation of  $SU(N_c)$

#### • Spontaneous breaking of chiral symmetry:

U(1)<sub> $\lambda$ </sub> (R-symmetry):  $\lambda' = e^{-i\varphi\gamma_5}\lambda$ ,  $\overline{\lambda}' = \overline{\lambda} e^{-i\varphi\gamma_5}$ Anomaly:  $U(1)_{\lambda} \rightarrow Z_{2Nc}$ Spontaneous breaking  $Z_{2N_c} \rightarrow Z_2$  by Gluino condensate  $\langle \lambda \lambda \rangle \neq 0$  $\leftrightarrow$  first order phase transition at  $m_{\tilde{q}} = 0$ 

#### Supersymmetric QCD

- additional quarks  $\psi$  and squarks  $\Phi_i$  in fundamental representation
- covariant derivatives, mass terms for  $(\psi, \Phi_i)$
- Yukawa interactions and scalar potential

 $i\sqrt{2}g\bar{\lambda}^{a}\left(\Phi_{1}^{\dagger}T^{a}P_{+}+\Phi_{2}T^{a}P_{-}\right)\psi$  $-i\sqrt{2}g\bar{\psi}\left(P_{-}T^{a}\Phi_{1}+P_{+}T^{a}\Phi_{2}^{\dagger}\right)\lambda^{a}$ 

### Spectrum of bound states

Expect colour neutral bound states of gluons and gluinos  $\rightarrow$  Supermultiplets

Predictions from effective Lagrangeans[1, 2]:

chiral supermultiplet (Veneziano, Yankielowicz) •  $0^-$  gluinoball a -  $\eta'$  $\overline{\lambda}\gamma_5\lambda$  $\sim$ •  $0^+$  gluinoball a -  $f_0$  $\overline{\lambda}\lambda$  $\sim$ • spin  $\frac{1}{2}$  gluino-glueball  $\sim \sigma_{\mu\nu} \operatorname{Tr} (F_{\mu\nu} \lambda)$ 

mixing of multiplets possible

Generalization (Farrar, Gabadadze, Schwetz) additional chiral supermultiplet •  $0^-$  glueball

•  $0^+$  glueball

• gluino-glueball

# Bound states on the lattice

- Glueballs:  $0^+$ ,  $0^- \cong \Box$ • Gluino-glueballs, Spin  $\frac{1}{2}$  Majorana:  $\chi_{\alpha} \simeq \frac{1}{2} F_{\mu\nu}^{\ a} (\sigma_{\mu\nu})_{\alpha\beta} \lambda_{\beta}^{a}$
- Gluino-balls:  $\overline{\lambda}\gamma_5\lambda$ : **a**  $\eta'$ , 0<sup>-</sup>,  $\overline{\lambda}\lambda$ : **a**  $f_0$ , 0<sup>+</sup>
- Mixing of Glueballs and Gluino-balls

#### **Effective masses**



#### Continuum limit



parameter ranges:  $0.2 < am_{a-\pi} < 0.7$ , lattice spacing  $0.053 \,\mathrm{fm} < a < 0.082 \,\mathrm{fm}$ , lattice sizes  $12^3 \times 24$  to  $24^3 \times 48$ 

Fit	$w_0 m_{g  ilde{g}}$	$w_0 m_{0^{++}}$	$w_0 m_{\mathrm{a}-\eta'}$
linear fit	0.917(91)	1.15(30)	1.05(10)
quadratic fit	0.991(55)	0.97(18)	0.950(63)
SU(2) SYM	0.93(6)	1.3(2)	0.98(6)

(in units of the gradient flow scale  $w_0$ )











#### SUSY breaking on the lattice:

- Local lattice theory: SUSY breaking unavoidable at any finite lattice spacing
- No general solution by Ginsparg-Wilson relation found (so far)
- Necessary to find specific solution for the model under consideration Approach for SUSY Yang-Mills theory (Curci, Veneziano) 1. Wilson action:

$$S = -\frac{\beta}{N_c} \sum_{p} \operatorname{Re} \operatorname{Tr} U_p$$

$$\left\{ \overline{\lambda}_x^a \lambda_x^a - K \sum_{\mu=1}^4 \left[ \overline{\lambda}_{x+\hat{\mu}}^a V_{ab,x\mu} (1+\gamma_{\mu}) \lambda_x^b + \overline{\lambda}_x^a V_{ab,x\mu}^t (1-\gamma_{\mu}) \lambda_{x+\hat{\mu}}^b \right] \right\}$$

$$K = \frac{1}{2\pi m_0} \quad \text{hopping parameter,} \quad m_0 : \text{bare gluino mass}$$

 $V_{ab,x\mu} = 2 \operatorname{Tr} \left( U_{x\mu}^{\dagger} T_a U_{x\mu} T_b \right),$ adjoint link variables

# Simulation details

#### Algorithms

• TS-PHMC and RHMC algorithm

#### Sign Problem

Fermion action:

$$S_f = \frac{1}{2}\overline{\lambda}Q\lambda = \frac{1}{2}\lambda M\lambda, \qquad M \equiv CQ$$
$$\int [d\lambda] e^{-S_f} = Pf(M) = \pm\sqrt{\det Q}$$

Effective gauge field action

$$S_{\text{eff}} = -\frac{\beta}{N_c} \sum_{p} \operatorname{Re} \operatorname{Tr} U_p - \frac{1}{2} \log \det Q[U]$$

Reweighting with sign Pf(M)

- Wilson fermions: mild sign problem, vanishes in continuum limit
- sign Pf(M): real negative eigenvalues of Q



- overlap fermions implement chirals symmetry on the lattice
- eigenvalue spectrum on circle
- RHMC + Overlap: stable polynomial approximation of sign-function to order N:

$$D_{\rm ov} = \frac{1}{2} + \frac{1}{2} \gamma_5 {\rm sign}(\gamma_5 D_W)$$

## Summary

- finalized analysis of SU(2) SUSY Yang-Mills particle spectrum [3]
- investigated phase transitions of SU(2) SUSY Yang-Mills [4, 5]
- most recently: final continuum extrapolations of the bound state spectrum of SU(3) SUSY Yang-Mills (gauge group of QCD) [6, 7]

#### SUSY breaking under control

- formation of bound state multiplets verified by numerical investigations
- SUSY restoration verified by Ward identities [8]

#### Outlook

- Overlap formulation as alternative, in particular for investigations of chiral symmetry breaking
- extensions to SQCD under investigation: mixed representations (fund.+adjoint), scalar fields, tuning

### References

- In recent studies of SU(3) SUSY Yang-Mills: one-loop O(a) improvement by clover term.
- 2. Tuning towards the chiral supersymmetric continuum limit:

• Wilson term breaks chiral symmetry and SUSY

 $2m_0 + 8$ 

- only tuning of gluino mass required to recover both symmetries in the continuum limit
- practical implementation: extrapolation to vanishing adjoint pion mass  $(m_{a-\pi})$ , cross check with SUSY Ward identities

Challenging extension towards supersymmetric QCD

• Yukawa couplings and scalar potential need to be fine tuned • order of 10 tuning parameters

• reduced tuning for chiral symmetric formulations (overlap fermions)

#### • not relevant for the current pa- -0.6 rameter range

#### Challenging measurement of bound state operators

#### • Flavour singlet meson operators with disconnected contributions



• Gluino-glue: variational methods (optimized ground state overlap)

• Glueballs, mixed with Gluino-balls: variational methods in large operator basis

• Baryon operators: spectacle contributions

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