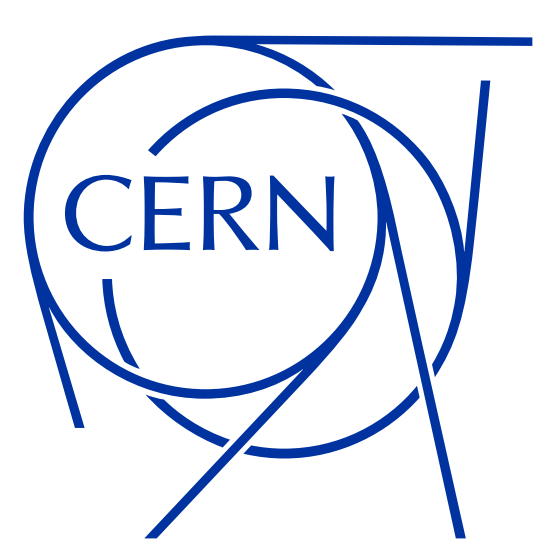


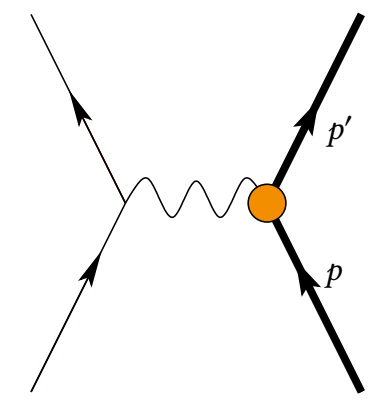
# Proton structure from lattice QCD

Jeremy Green, in collaboration with Michael Engelhardt, Nesreen Hasan, Stefan Krieg, Stefan Meinel, John Negele, Andrew Pochinsky, Giorgio Silvi, and Sergey Syritsyn



## Electromagnetic form factors

The distributions of charge and magnetization in a proton are probed in elastic electron-proton scattering.



Photon-proton vertex is parameterized by two form factors,

$$\langle p' | J_\mu | p \rangle = \bar{u}(p') \left[ \gamma_\mu F_1(Q^2) + \frac{i\sigma_{\mu\nu}(p' - p)^\nu}{2m} F_2(Q^2) \right] u(p), \quad Q^2 = -(p' - p)^2.$$

These combine to form the electric and magnetic form factors,

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{4m^2} F_2(Q^2) \quad \rightarrow \text{Fourier transform of charge density}$$

$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2) \quad \rightarrow \text{Fourier transf. of magnetization density.}$$

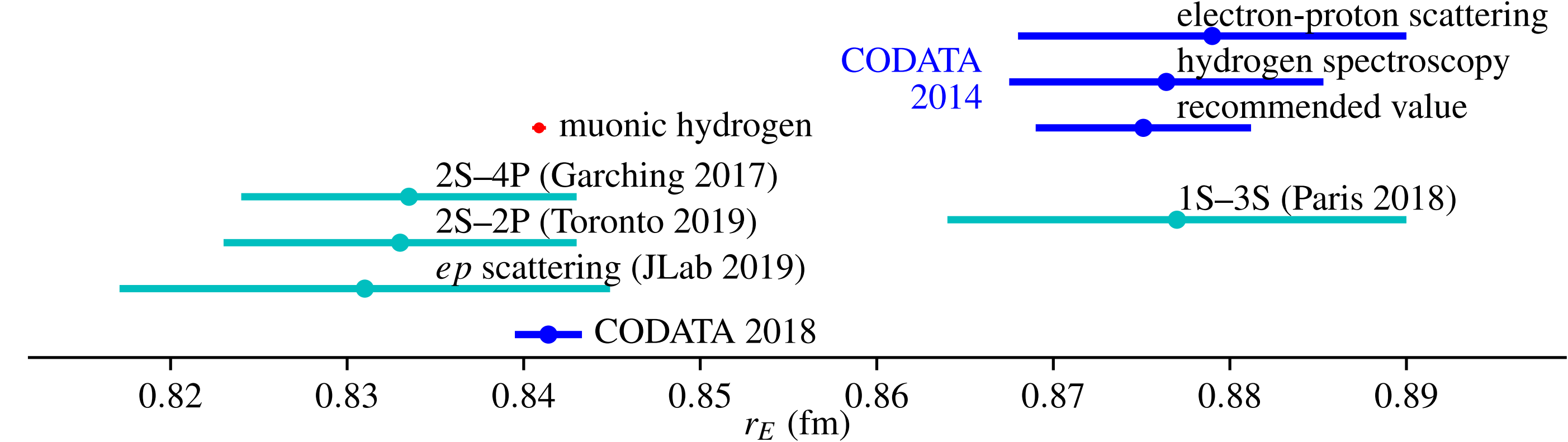
Near  $Q^2 = 0$  they contain key properties of the proton:

electric charge  $1 = G_E(0)$

rms charge radius  $r_E^2 = -6 \frac{dG_E}{dQ^2} \Big|_{Q^2=0}$

magnetic moment  $\mu = G_M(0)$ .

Muonic hydrogen experiment led to very precise  $r_E$  but also *radius puzzle*.



Newer *ep* scattering and electronic hydrogen spectroscopy experiments have cast doubt on precision of older measurements.

A. Beyer *et al.*, *Science* 358, 79–85 (2017), H. Fleurbaey *et al.*, *Phys. Rev. Lett.* 120, 183001 (2018), N. Bezginov *et al.*, *Science* 365, 1007–1012 (2019), W. Xiong *et al.*, *Nature* 575, 147–150 (2019).

## Finite-volume effects in form factors (preliminary)

Lattice calculations are done in periodic spatial volume of size  $L^3$ .

Generically effect on nucleon observables is suppressed  $\sim e^{-m_\pi L}$ .

Volume also constrains  $Q^2$ : momenta  $p_j = 2\pi n/L$ ,  $n \in \mathbb{Z}$ .

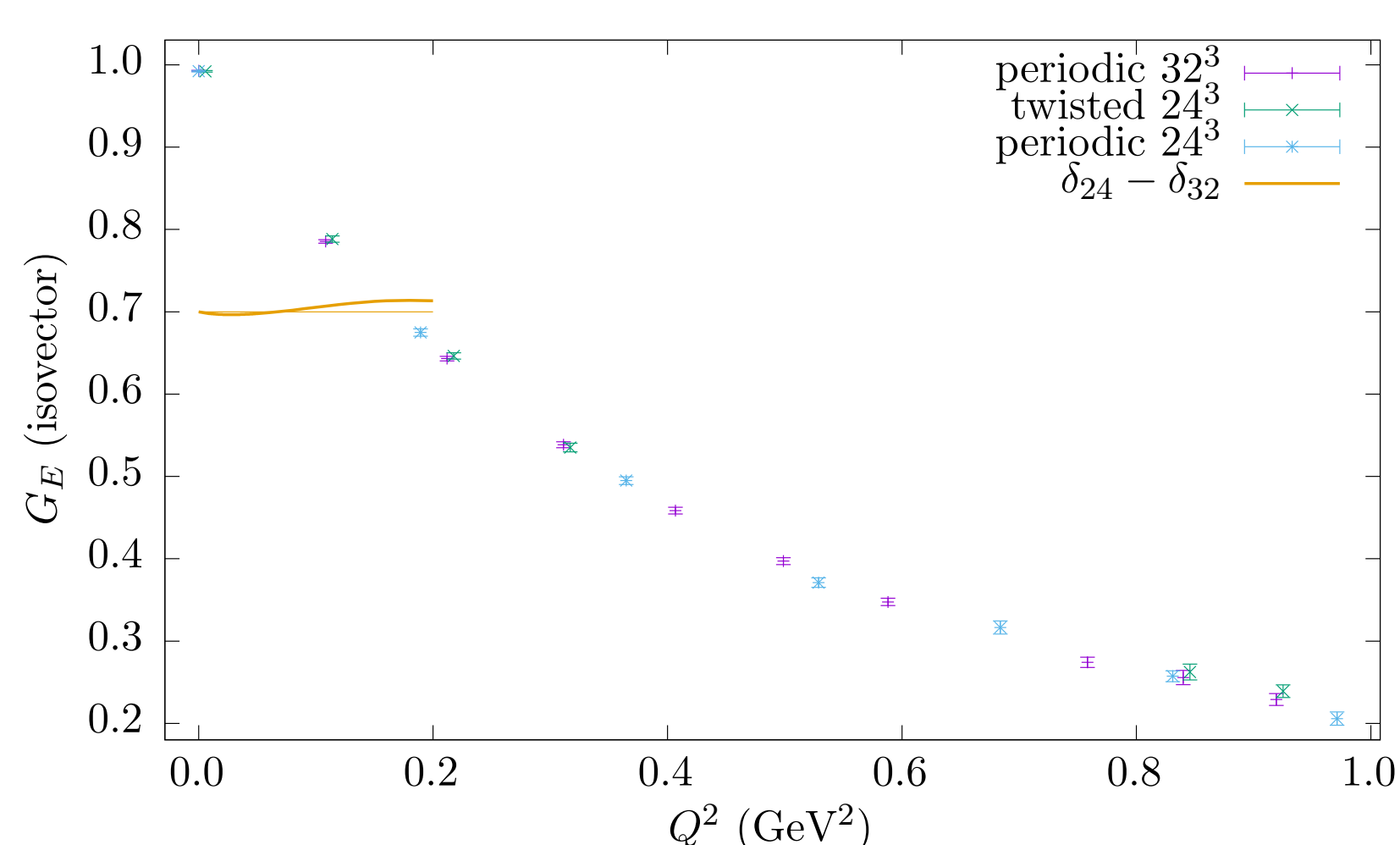
How to study  $L$  dependence at fixed  $Q^2$ ? Use twisted boundary cond.:

$$q(\vec{x}) = e^{i\theta} q(\vec{x} + \hat{j}L) \implies p_j = (2\pi n + \theta)/L, \quad n \in \mathbb{Z}.$$

Lattice setup:  $m_\pi \approx 250$  MeV,  $a = 0.116$  fm, two ensembles:

- $32^3 \times 48$ ,  $m_\pi L = 4.8$ ,
- $24^3 \times 48$ ,  $m_\pi L = 3.6$ , plus twisted B.C. to match momenta.

Preliminary results for isovector  $G_E$  and  $G_M$ :

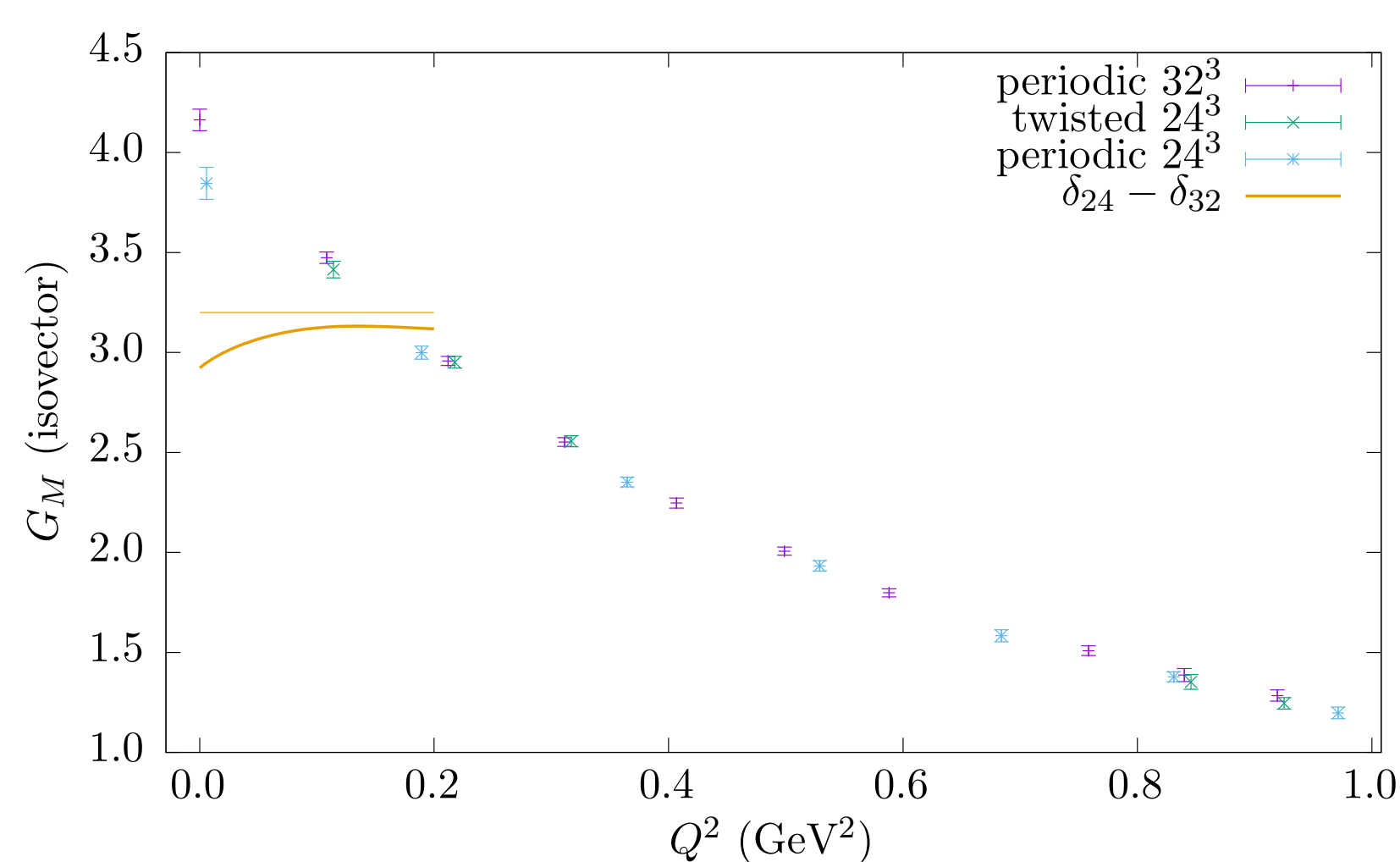


Orange curve: predicted finite-volume effect from partially quenched heavy baryon ChPT.

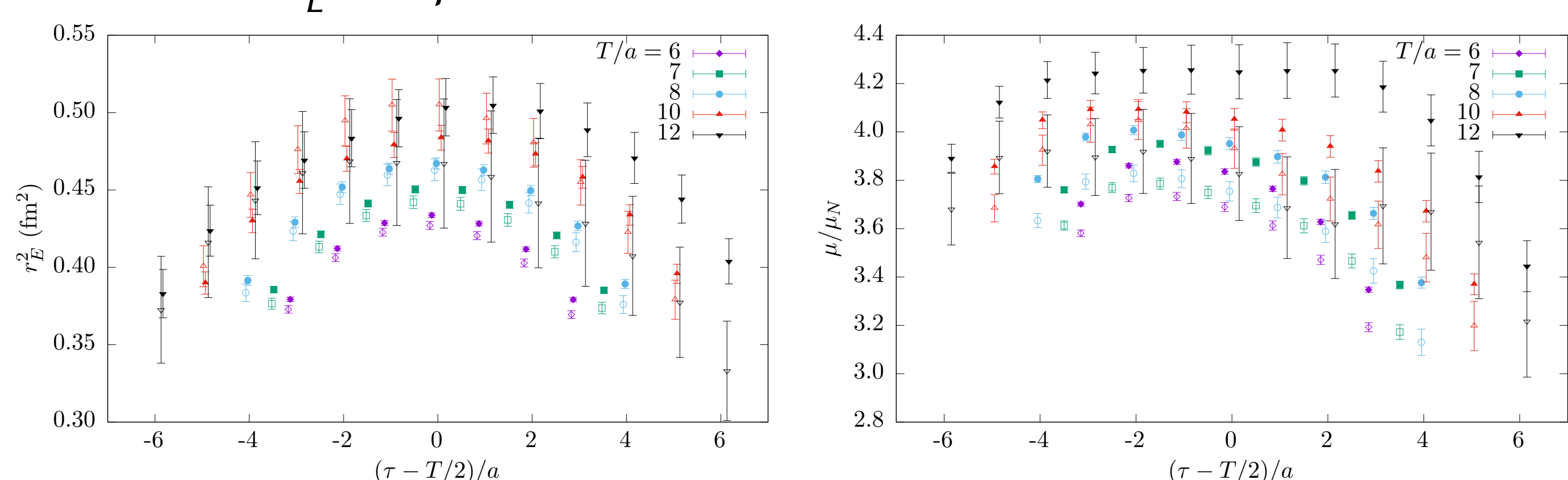
F.-J. Jiang and B. C. Tiburzi, *Phys. Rev. D* 78, 114505 (2008)

Exact  $\theta$ -derivative yields direct calculation of magnetic moment at  $Q^2 = 0$ .

N. Hasan, J.G., S. Meinel *et al.*, *Phys. Rev. D* 97, 034504 (2018); N. Hasan, talk at Lattice 2017



Plateaus for  $r_E^2$  and  $\mu$  from derivative method:



$\sim -2\%$  effect for  $r_E^2$ : smaller than ChPT and opposite sign.

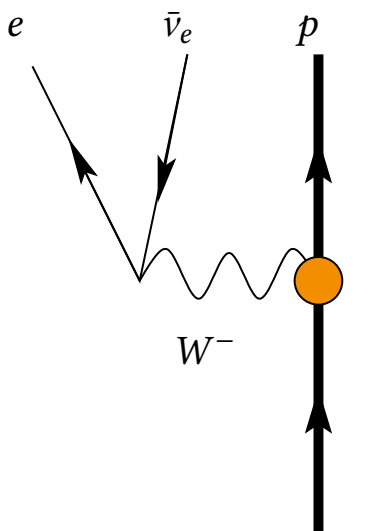
$\sim -5\%$  effect for  $\mu$ : similar to ChPT.

Excited-state contributions are significant.

## Neutron beta decay

In the Standard Model, neutrons decay by emitting a virtual  $W^-$  boson. The quark-level coupling to  $W$  bosons is of the form  $V - A$ . The baryon-level couplings are modified by QCD:

$$g_V \approx 1, \quad g_A/g_V = 1.2732(23) \text{ PDG 2019}$$



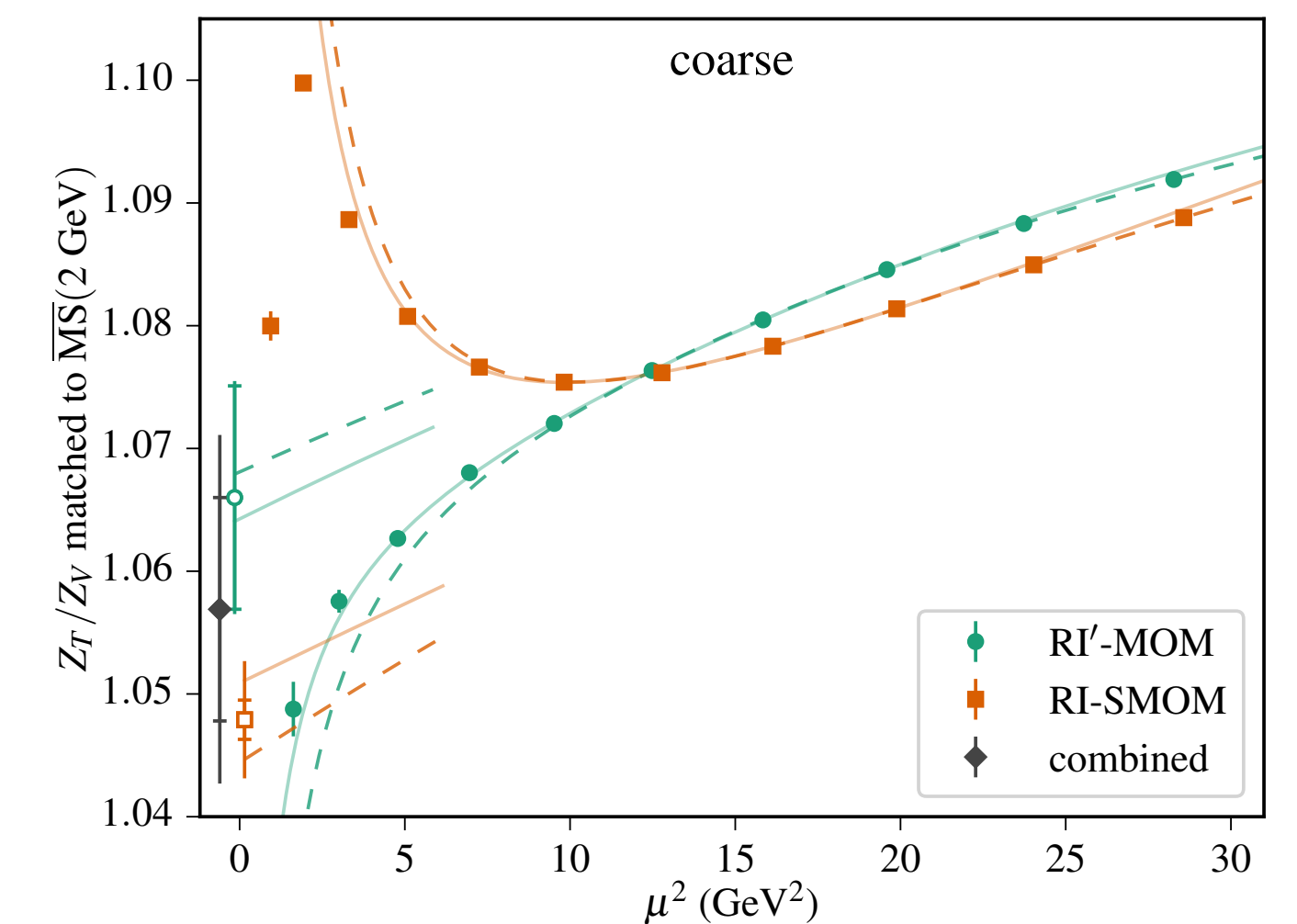
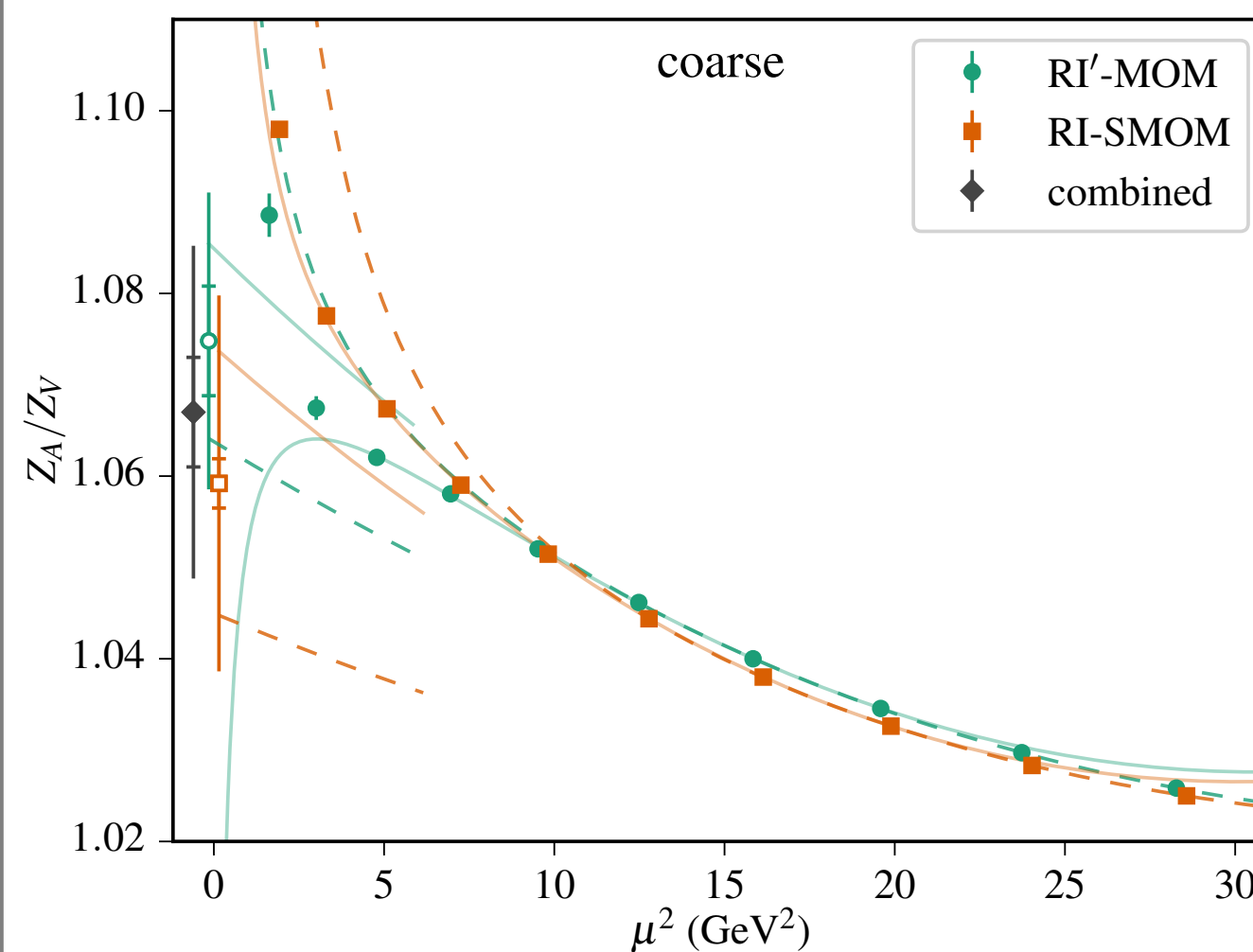
Precision  $\beta$ -decay experiments may be sensitive to beyond-the-Standard-Model physics by detecting scalar or tensor couplings.

Need to compute the corresponding “charges”,  $g_S$  and  $g_T$  on the lattice.

Lattice setup:  $m_\pi = m_\pi^{\text{phys}}$ , two ensembles: N. Hasan, J.G. *et al.*, *Phys. Rev. D* 99, 114505 (2019)

- $a = 0.116$  fm,  $48^4$ ,  $m_\pi L = 3.9$ ,
- $a = 0.093$  fm,  $64^4$ ,  $m_\pi L = 4.0$ .

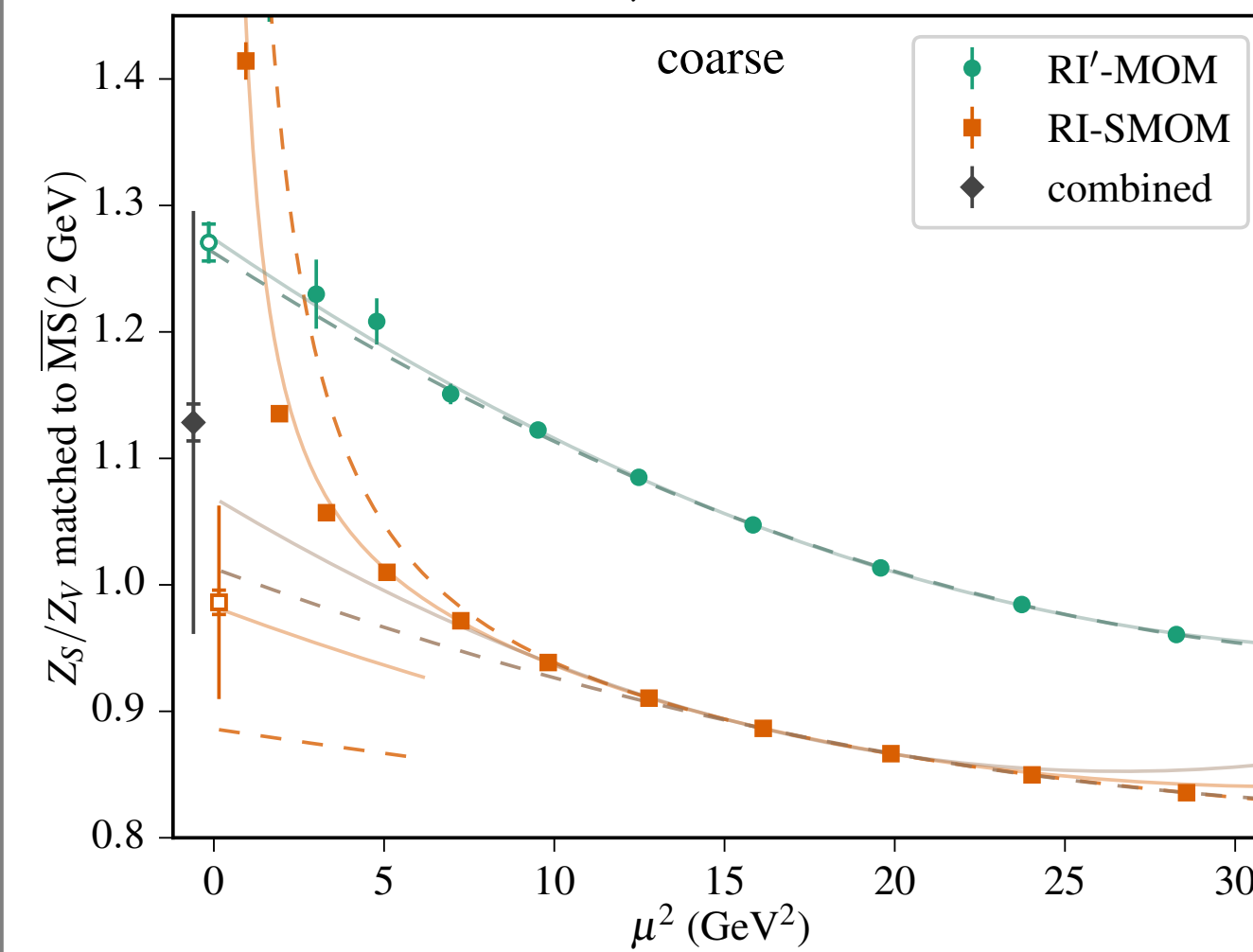
Need to compute renormalization factors  $Z_{A,S,T}$ .



Two different intermediate renormalization schemes used.

Reasonable agreement for  $Z_{A,T}$ .

Large difference for  $Z_S$   $\rightarrow$  large uncertainty for  $g_S$ .



Final results:

$$g_A = 1.265(49), \quad g_T = 0.972(41), \quad g_S = 0.927(303).$$

Large discrepancy between two renormalization schemes for scalar needs further study. Could affect results by other collaborations!

## Reducing excited-state effects

Challenge in lattice calculations: can't exactly create proton.

Must use Euclidean time evolution to suppress excited states  $\sim e^{-Et}$ .

Can the standard creation operator  $\chi_1^\dagger$  be improved?

J.G. *et al.*, *Phys. Rev. D* 100, 074510 (2019)

$$\chi_1 \sim (uud)_{\frac{1}{2}} \quad (\text{standard})$$

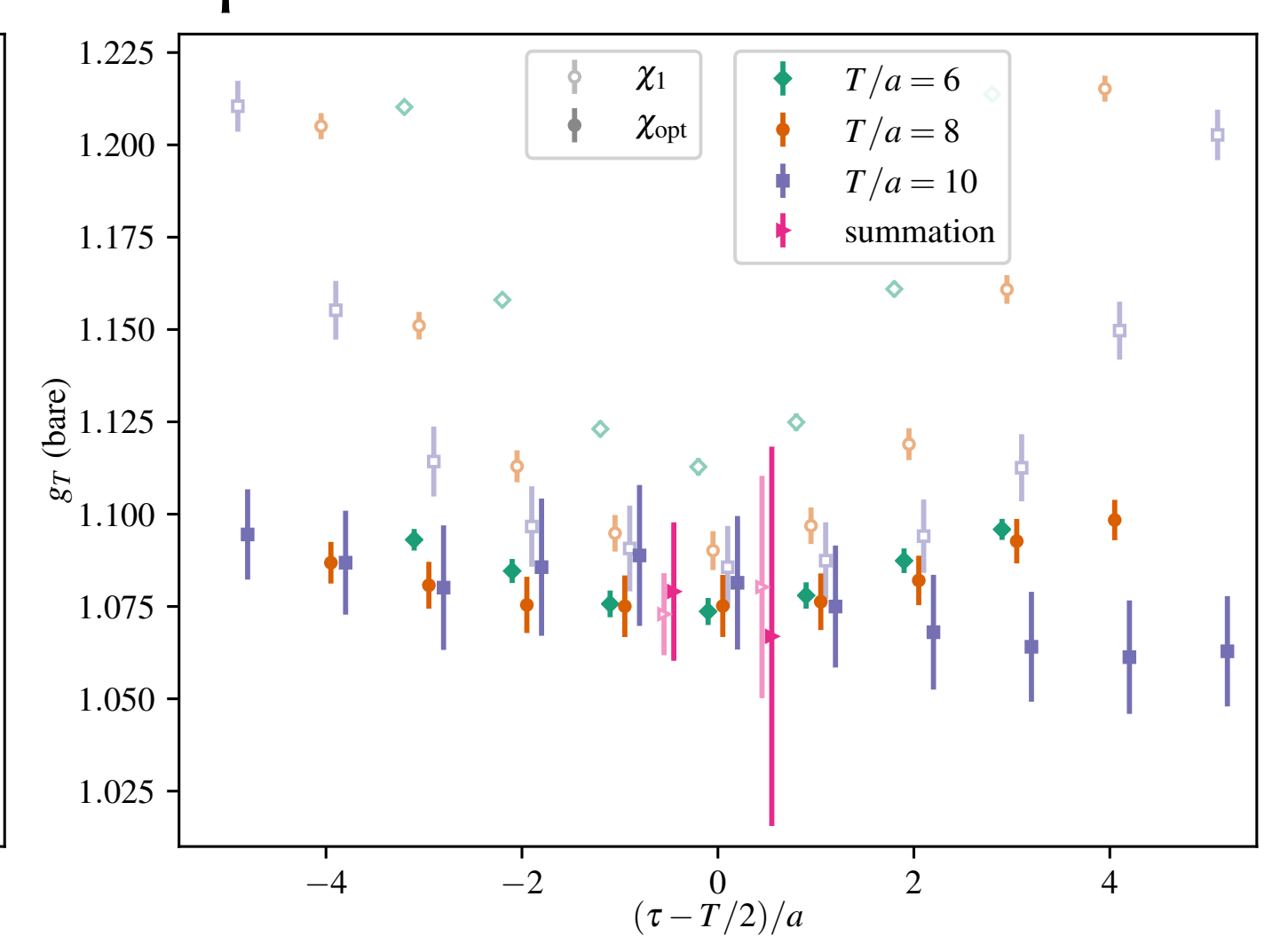
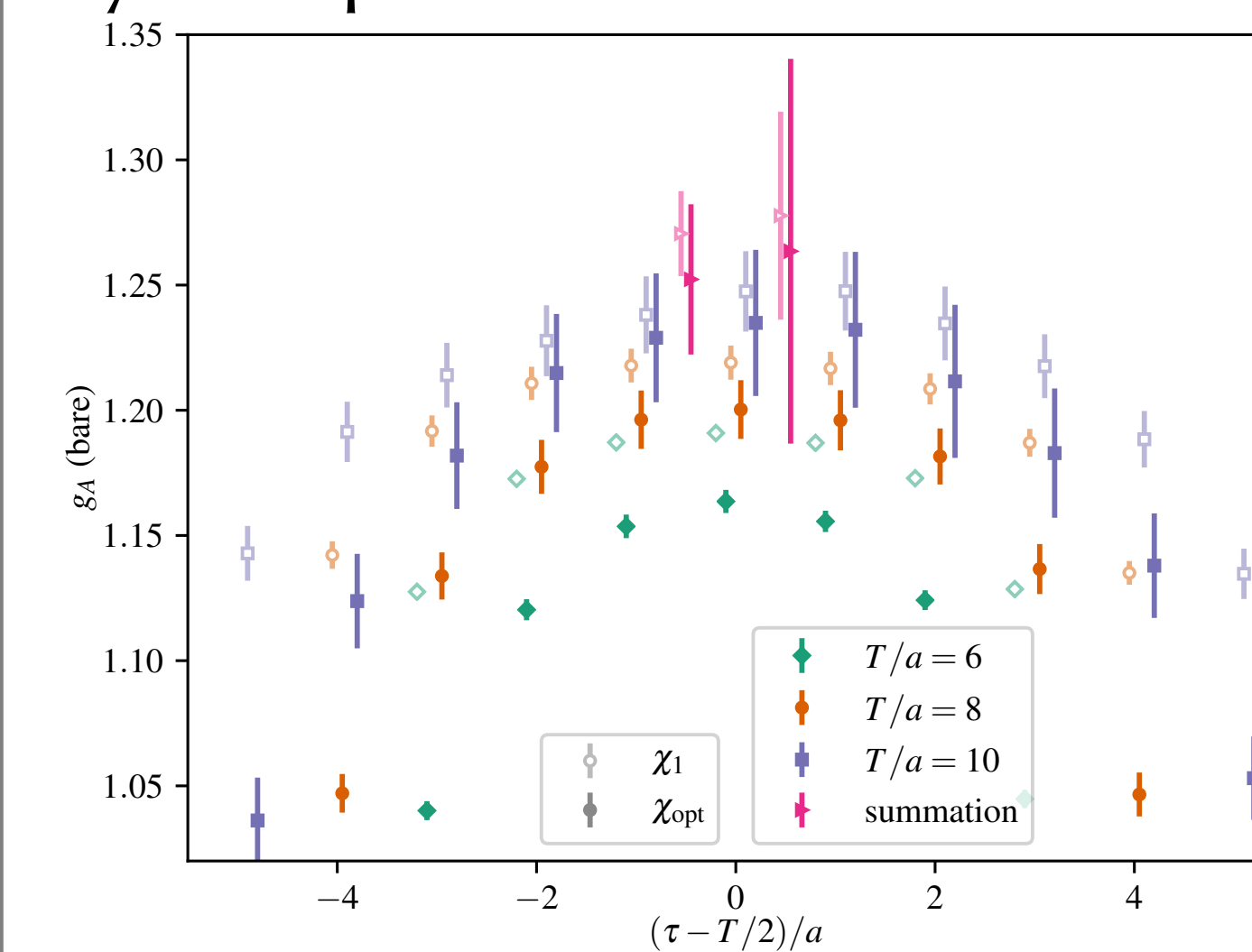
$$\chi_2 \sim (uud)_{\frac{1}{2}} g \quad (\text{hybrid})$$

$$\chi_3 \sim (uud)_{\frac{3}{2}} g \quad (\text{hybrid})$$

Use variational setup: optimize

$$\chi_{\text{opt}} = c_1 \chi_1 + c_2 \chi_2 + c_3 \chi_3$$

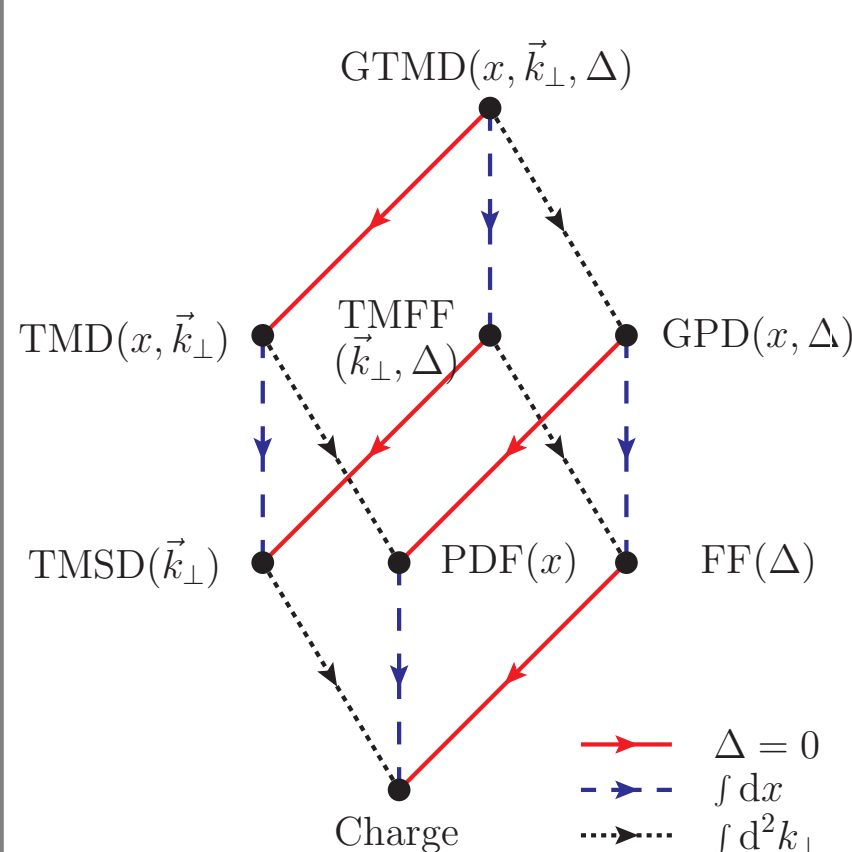
Hybrid operators chosen for their low computational cost.



Optimized operator  $\chi_{\text{opt}}$  has significantly reduced excited-state effects in  $g_T$  but increased effects in  $g_A$ . No universal improvement.

Need to systematically target lowest-lying excitations!

## Generalized TMDs (ongoing calculations)



Distributions of quarks in the proton that are multidimensional: longitudinal momentum, transverse momentum, and transverse position.

Focus on quark orbital angular momentum and spin-orbit correlation. Use momentum derivative method and chiral fermions.

C. Lorcé, B. Pasquini, and M. Vanderhaeghen, *JHEP* 2011 (05), 041

M. Engelhardt, *Phys. Rev. D* 95, 094505 (2017); M. Engelhardt *et al.*, *PoS Lattice* 2018, 115, *SPIN* 2018, 047 [1901.00843]