

Optical Pumping of Electron Spins in Quantum Dot Ensembles

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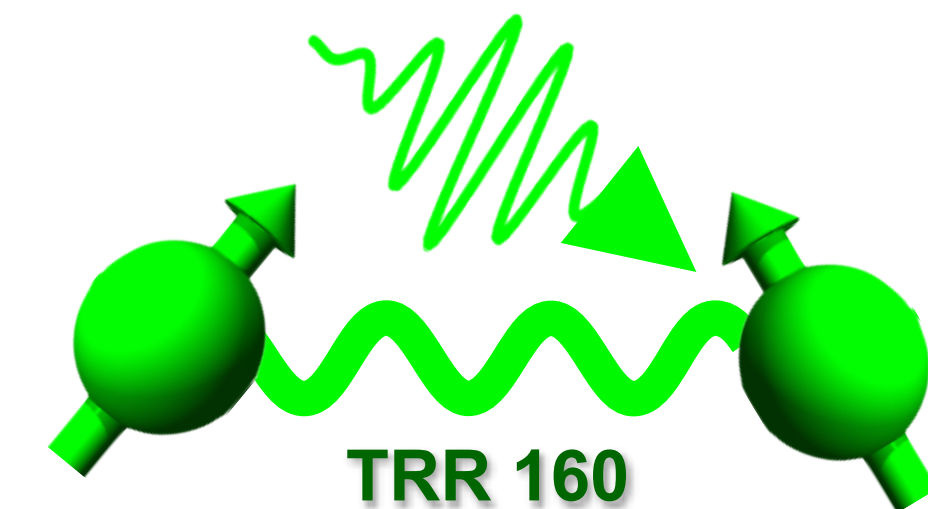
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Transregio I 60 Dortmund/St Petersburg: Coherent manipulation of interacting spin excitations in tailored semiconductors

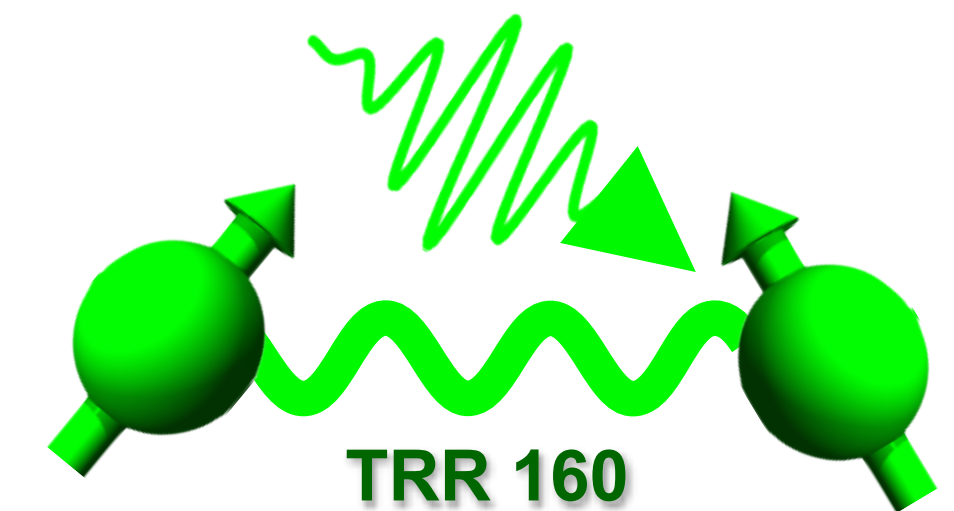


TU Dortmund

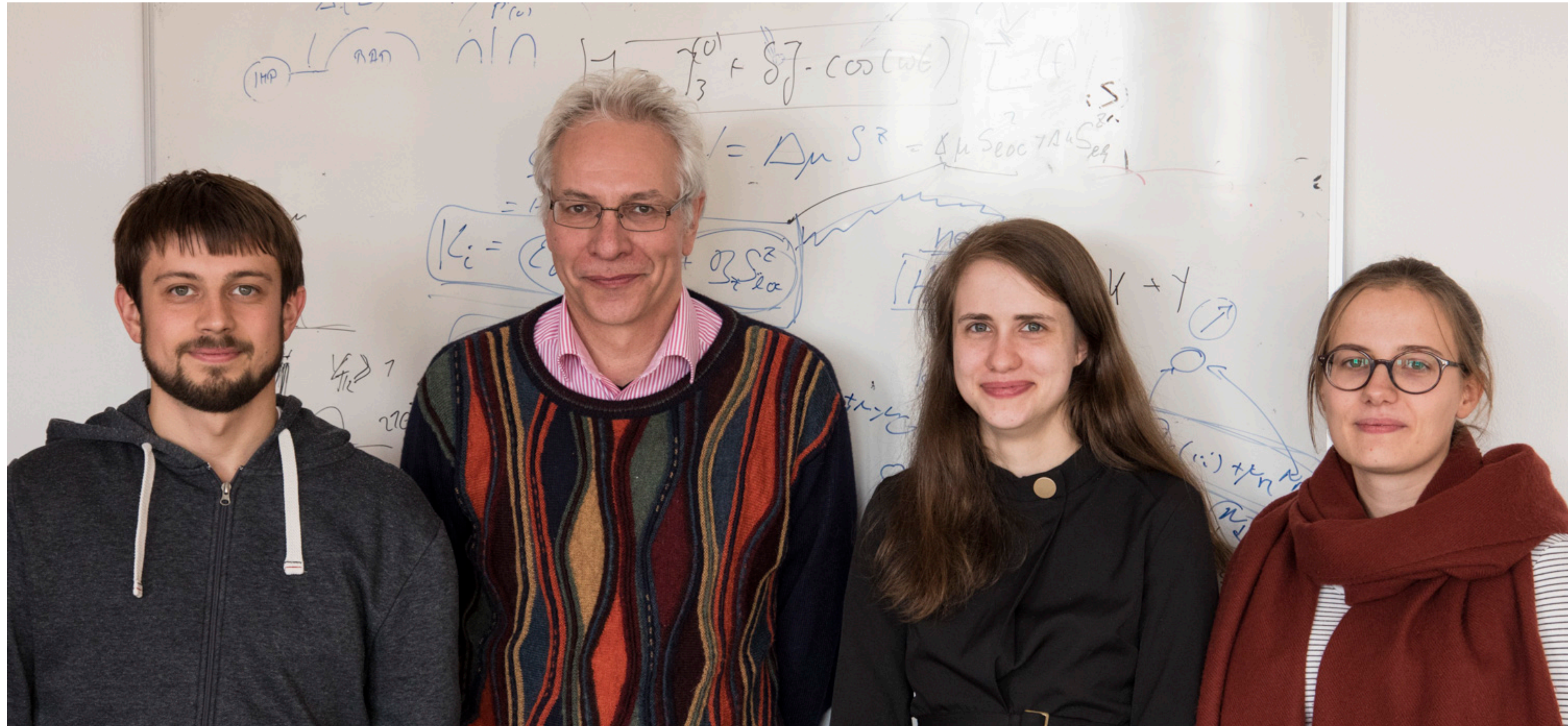
Ioffe Institute, St. Petersburg



St. Petersburg State University



TRR 160



Andreas Fischer

Dr. Natalie Jäschke

Iris Kleinjohann

MENU ▾

nature

Article | Published: 23 October 2019

Quantum supremacy using a programmable superconducting processor

Frank Arute, Kunal Arya, [...] John M. Martinis 

Nature 574, 505–510(2019) | [Cite this article](#)

We verify that the **quantum processor is working properly** using a method called cross-entropy benchmarking [11,12,14], which compares how often each bitstring is observed experimentally with its corresponding ideal probability computed **via simulation on a classical computer**.

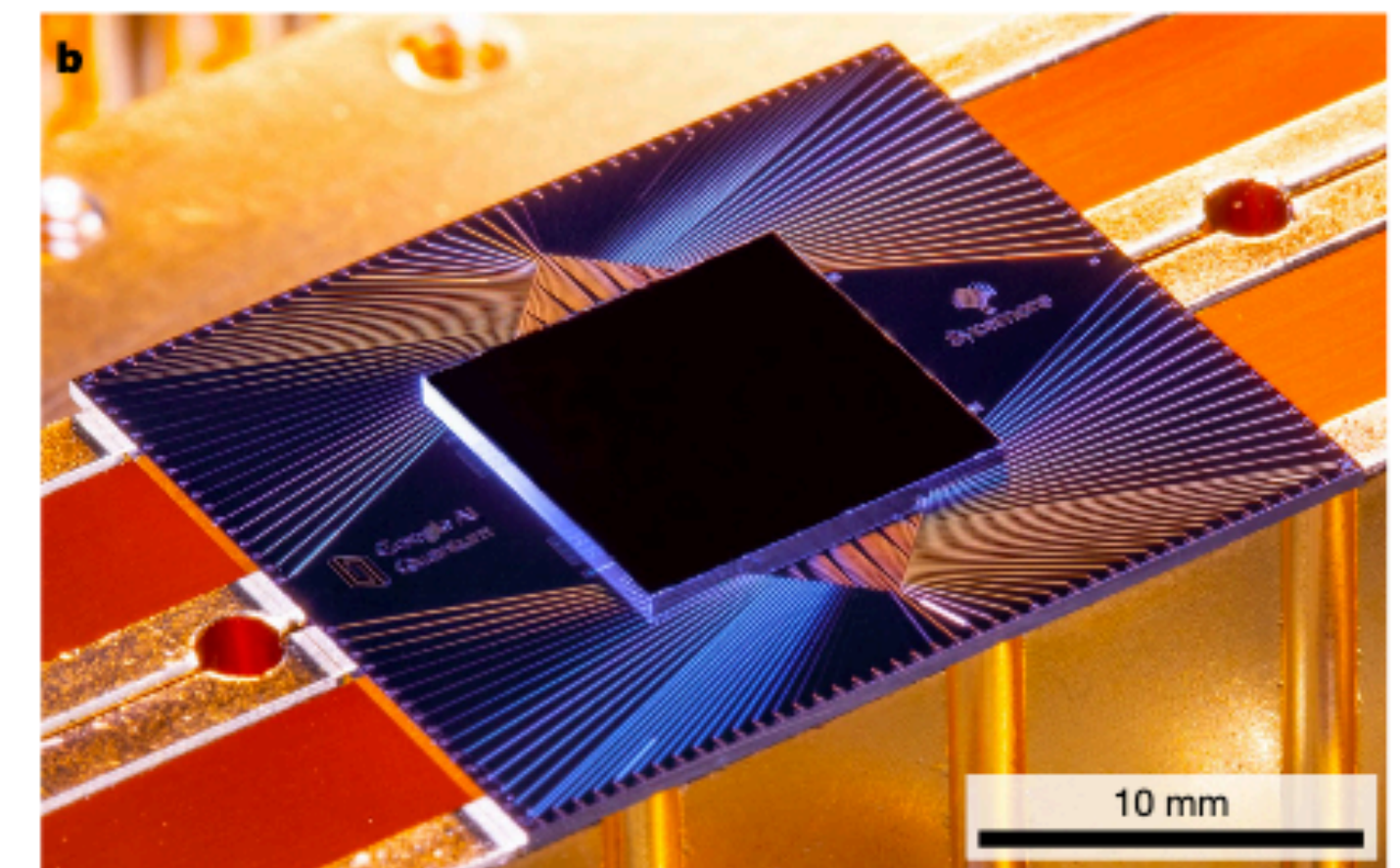
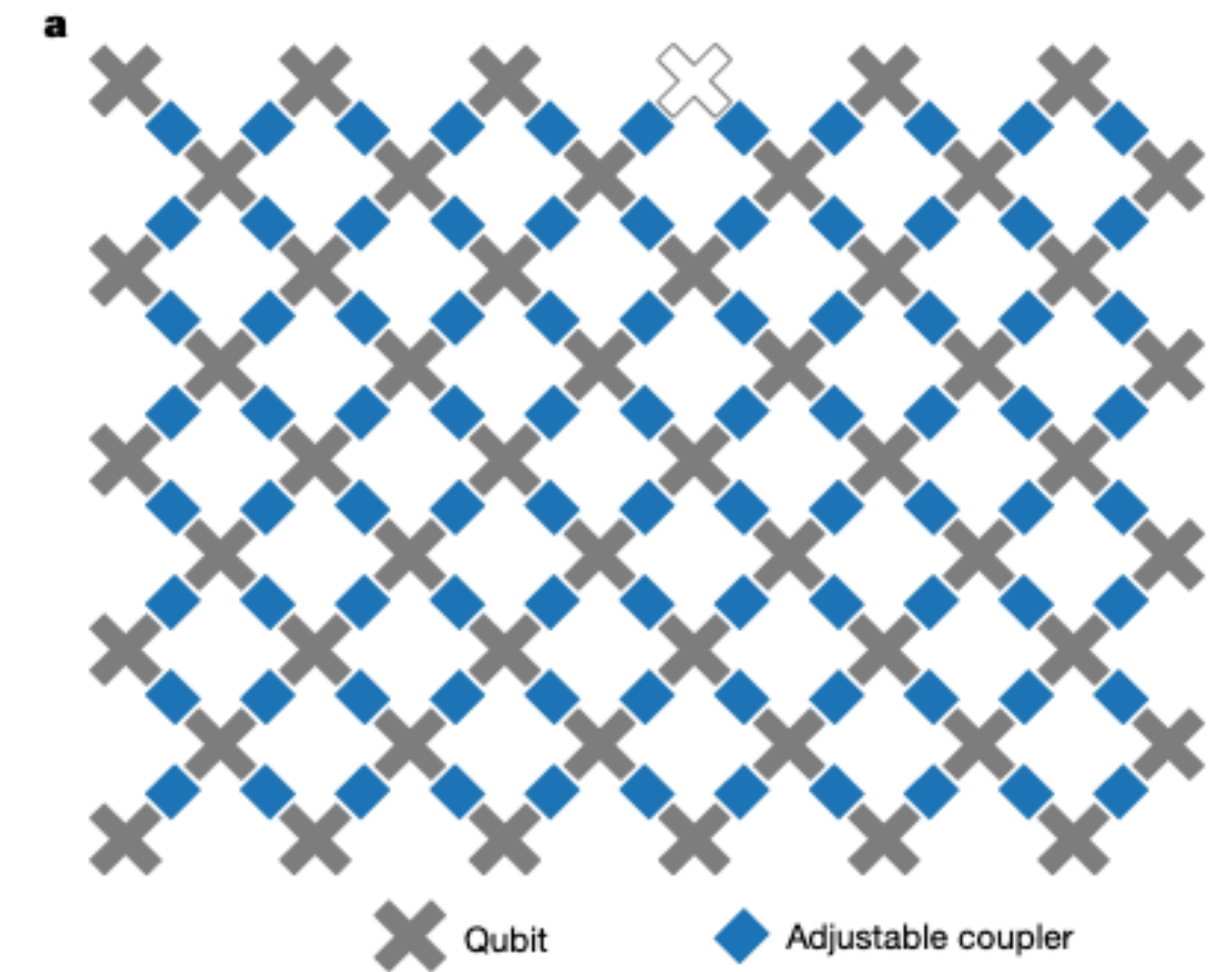


Fig. 1 | The Sycamore processor. **a**, Layout of processor, showing a rectangular array of 54 qubits (grey), each connected to its four nearest neighbours with couplers (blue). The inoperable qubit is outlined. **b**, Photograph of the Sycamore chip.

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Quantencomputer

Ein Quantum Überlegenheit?

Google verkündet die "Quantum Supremacy": Sein Sycamore-Chip soll dramatisch viel schneller rechnen als jeder Supercomputer bislang. IBM-Forscher zweifeln am Durchbruch.

Von **Eike Kühl**

23. Oktober 2019, 16:06 Uhr / Aktualisiert am 23. Oktober 2019, 16:10 Uhr / [102 Kommentare](#)

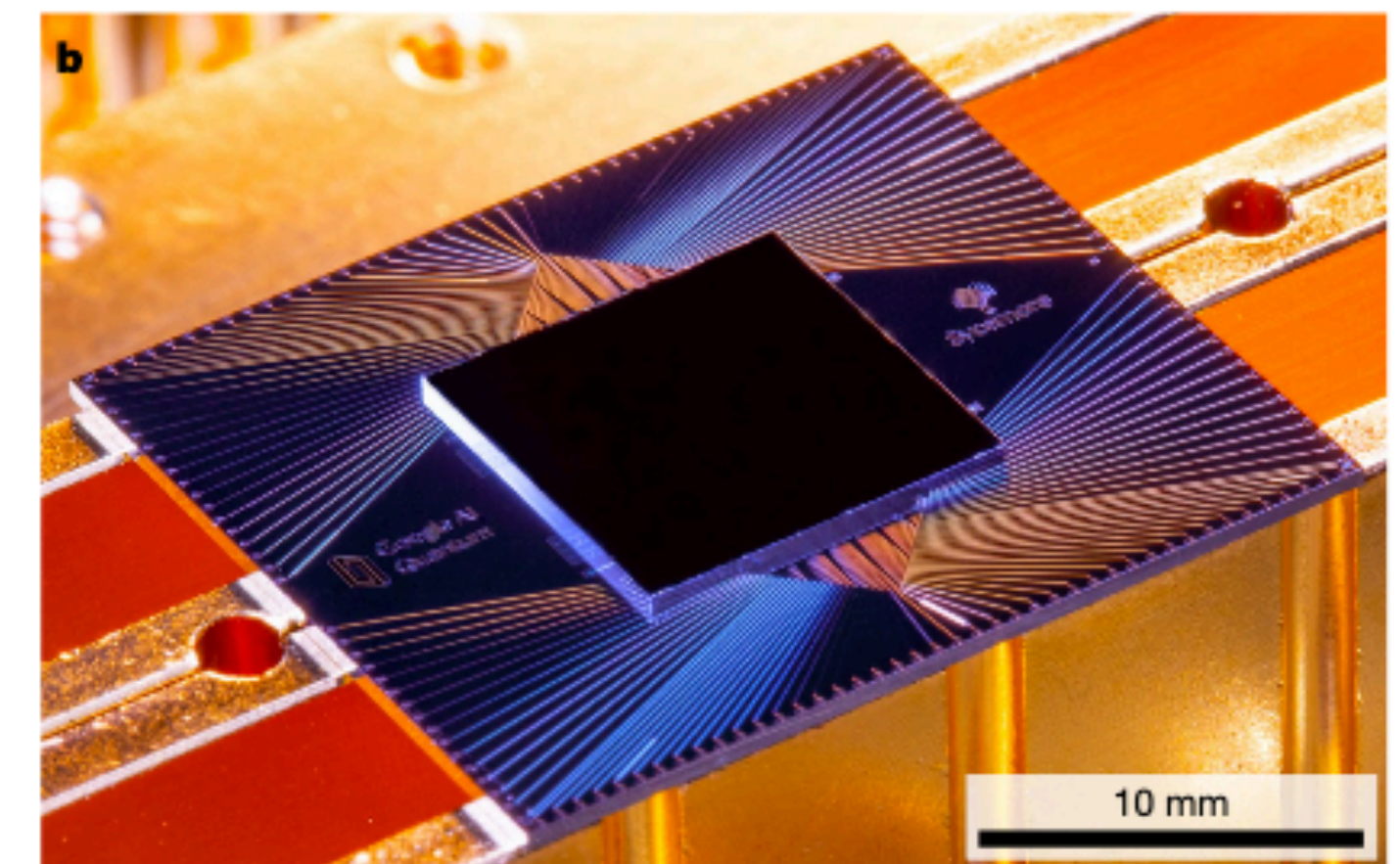
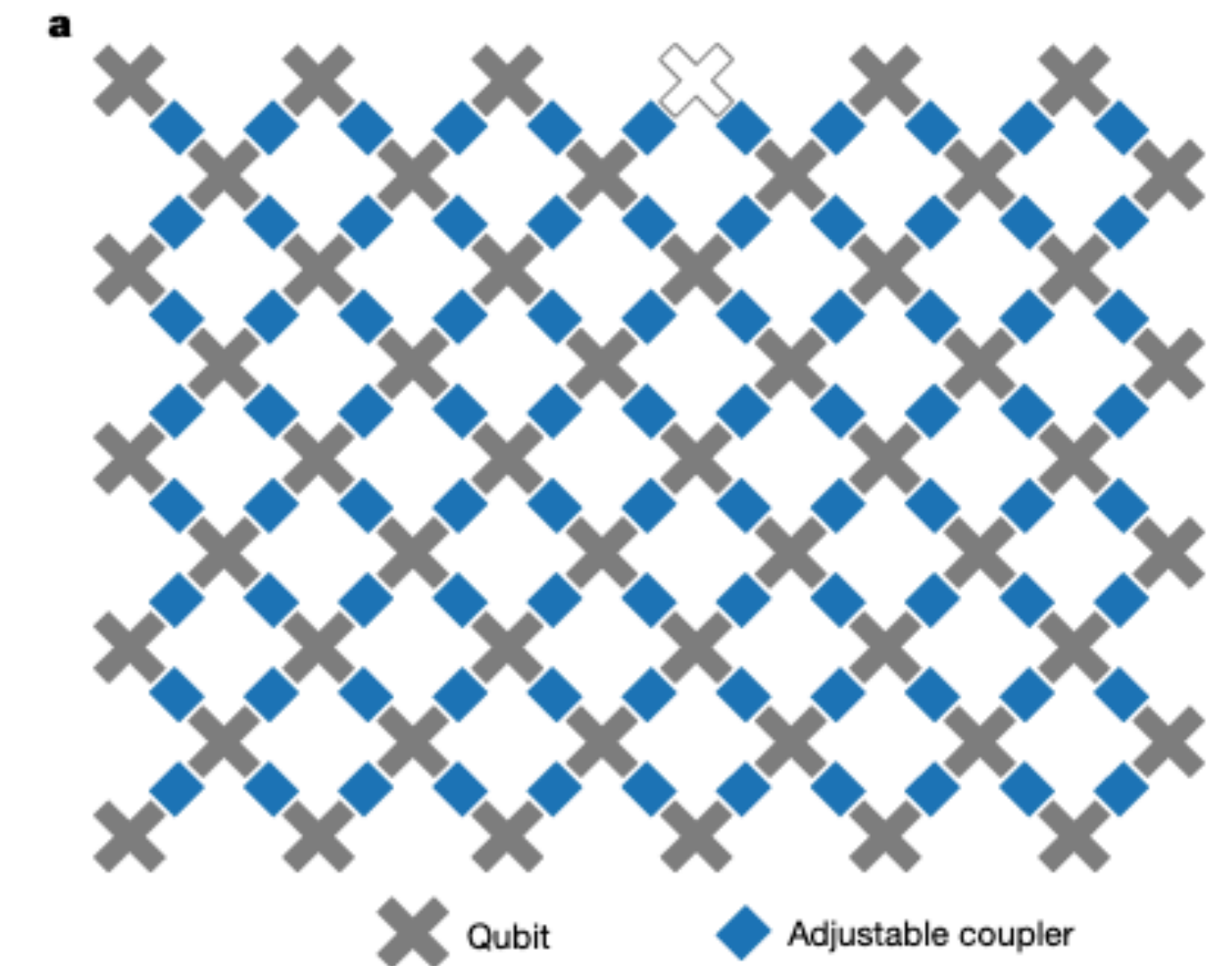
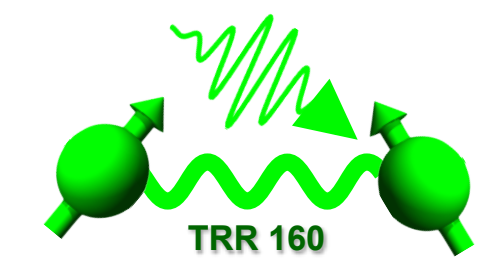
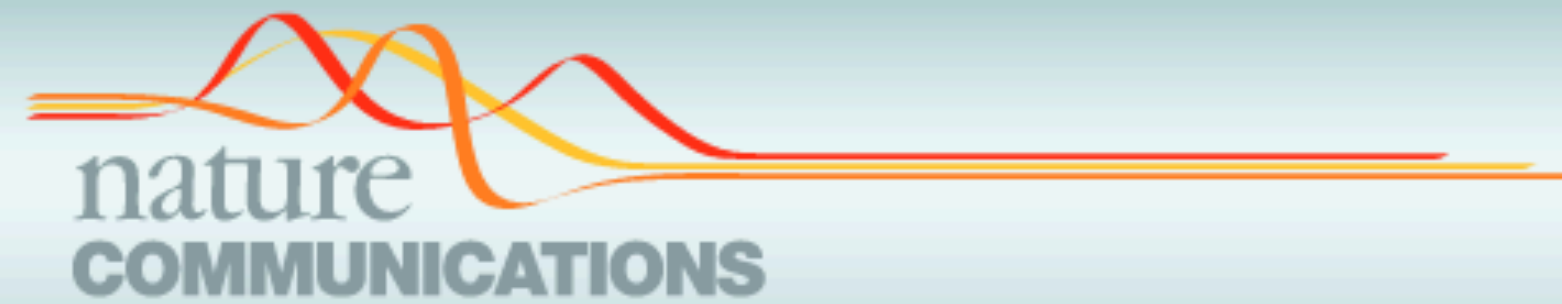
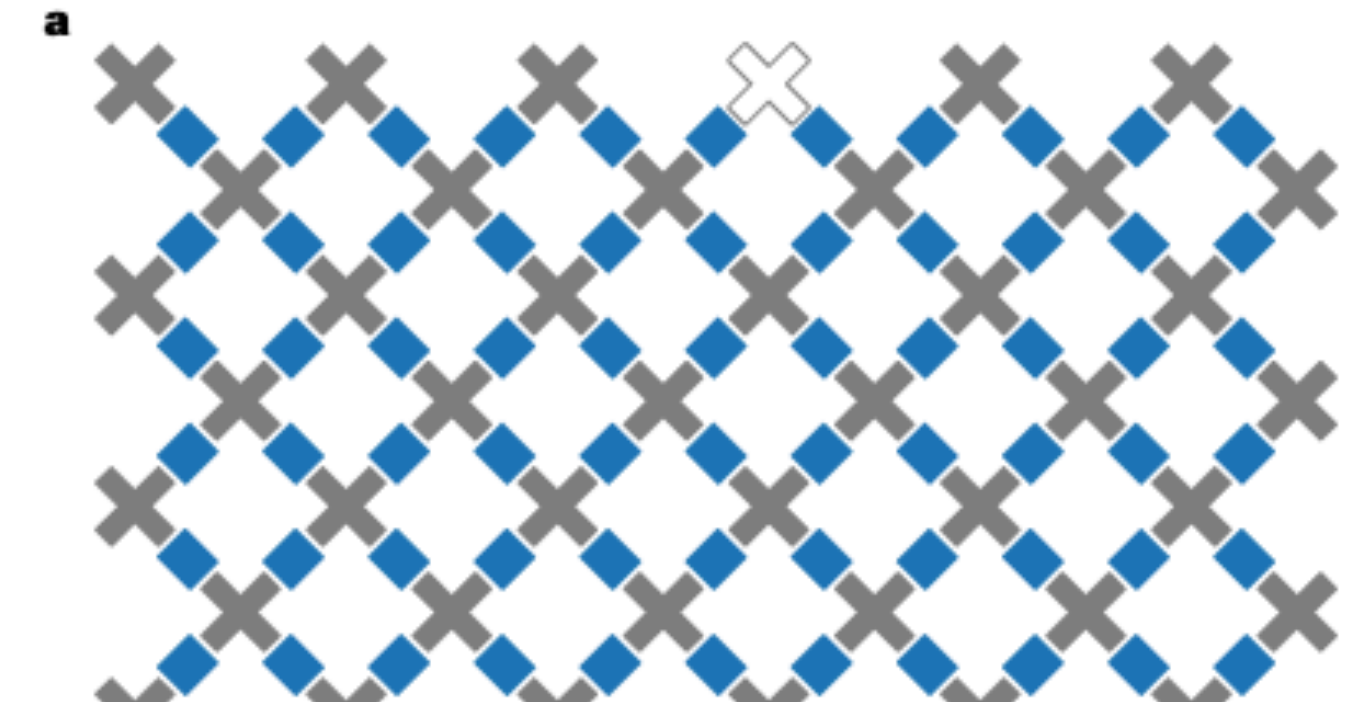


Fig. 1 | The Sycamore processor. a, Layout of processor, showing a rectangular array of 54 qubits (grey), each connected to its four nearest neighbours with couplers (blue). The inoperable qubit is outlined. **b**, Photograph of the Sycamore chip.



Motivation

- works only at very low temperature (superconducting qubits)
- only nearest neighbor coupling of qubit
- IBM Q: results depend on the daily calibration
- quantum mechanical simulations on classical computers are still needed



Feb 11.2020

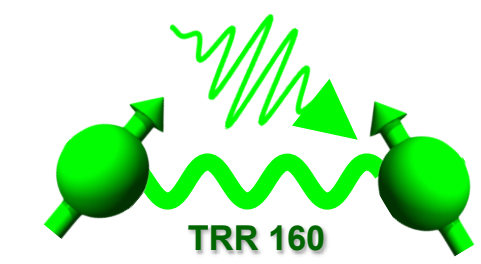
ARTICLE

<https://doi.org/10.1038/s41467-019-14053-w>

OPEN

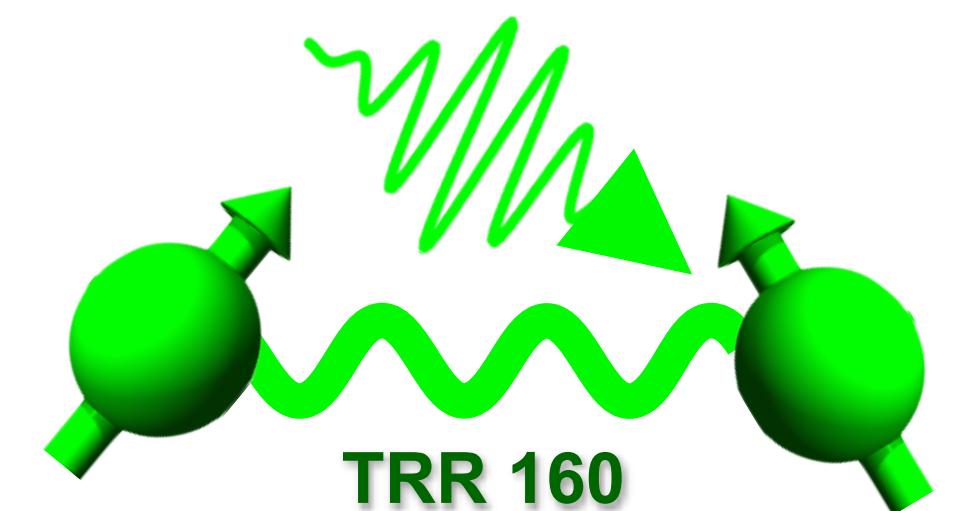
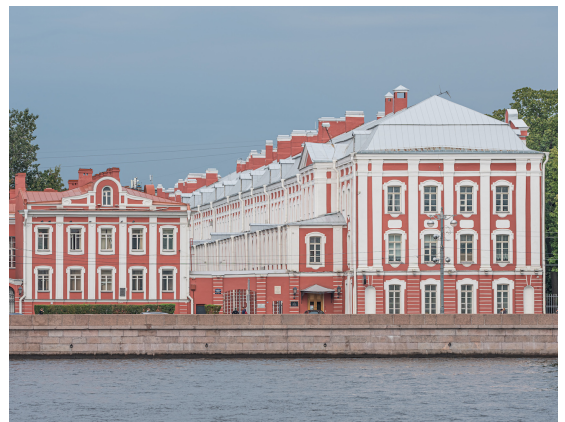
Coherent spin control of s-, p-, d- and f-electrons in a silicon quantum dot

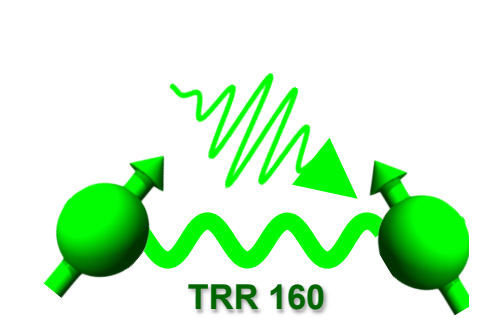
R.C.C. Leon^{1*}, C.H. Yang¹, J.C.C. Hwang^{1,6}, J. Camirand Lemyre², T. Tanttu¹, W. Huang¹,
K.W. Chan¹, K.Y. Tan³, F.E. Hudson¹, K.M. Itoh⁴, A. Morello¹, A. Laucht¹, M. Pioro-Ladrière^{2,5},
A. Saraiva^{1*} & A.S. Dzurak^{1*}



TRR 160: Coherent manipulation of interacting spin excitations in tailored semiconductors

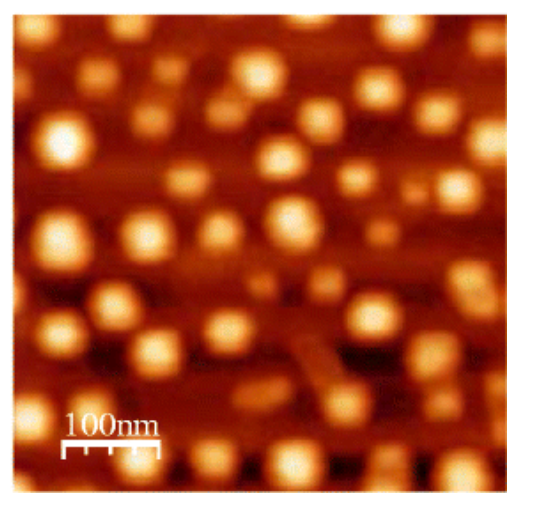
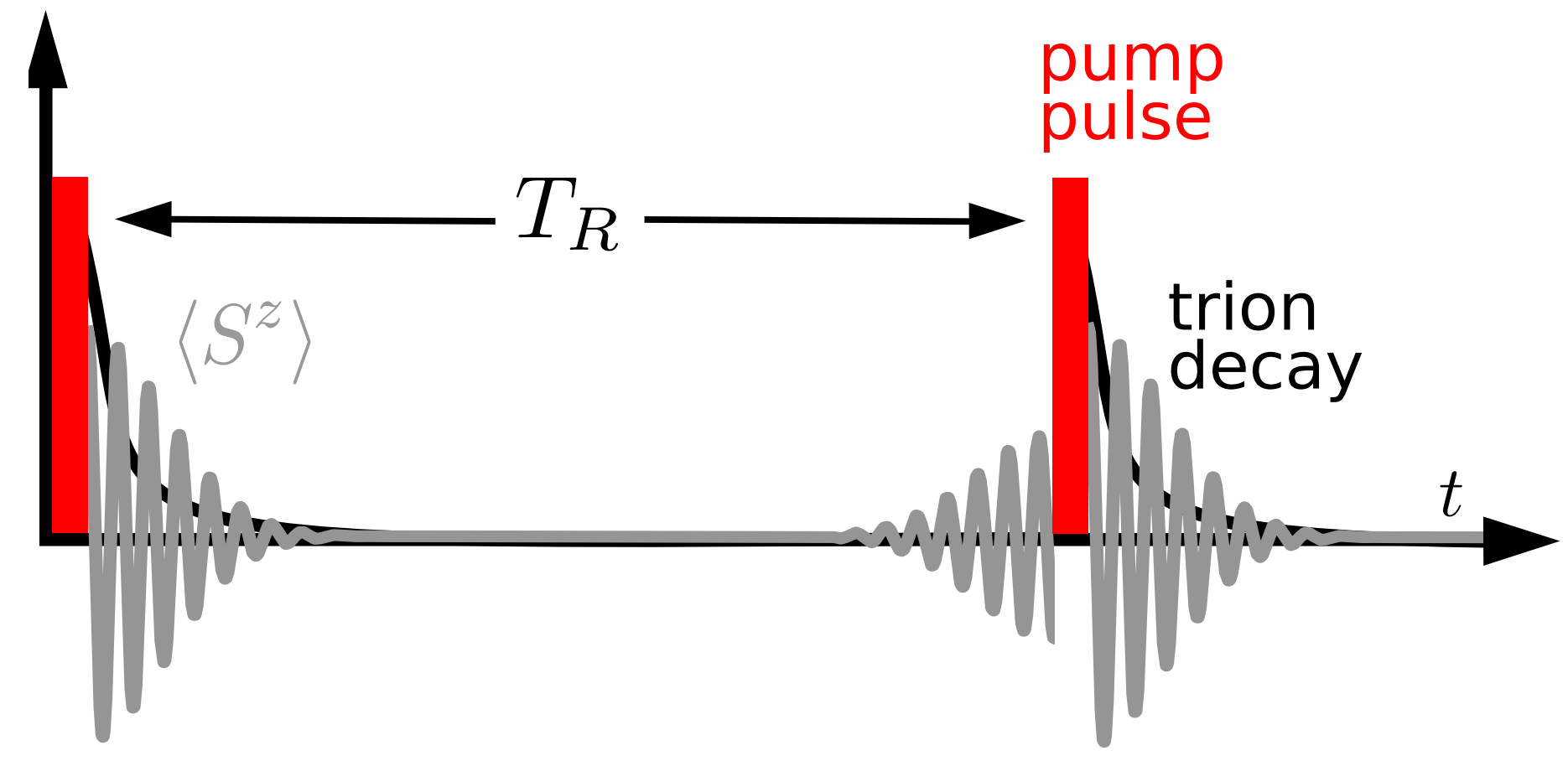
- application of spin excitation for future technology
- exploitation of spin coherence: reducing power consumption of charge driven electronics
- potential embedding in traditional semiconductor structures
- manipulation with light: opto-spintronic integration



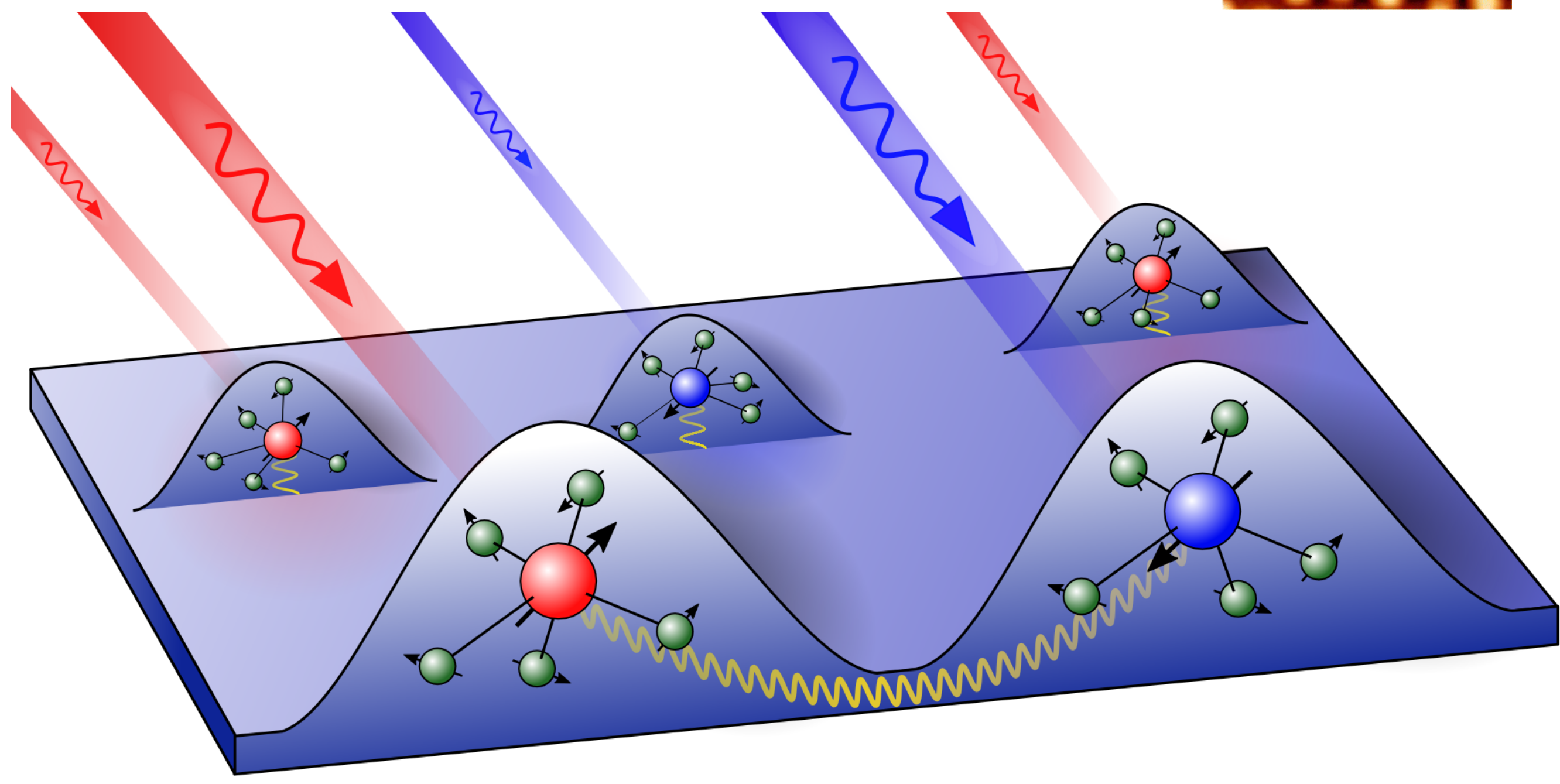


Optical pumping of electron spins in Quantum dot ensemble

1. Periodic pumping of quantum dot ensembles

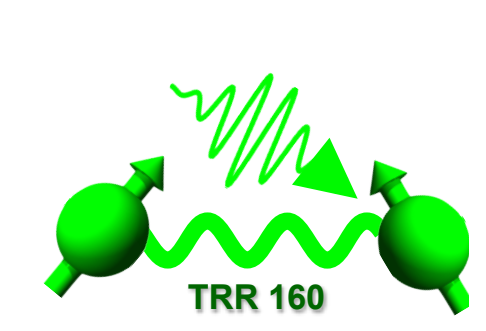


2. Multi-color pumping: addressing different QDs



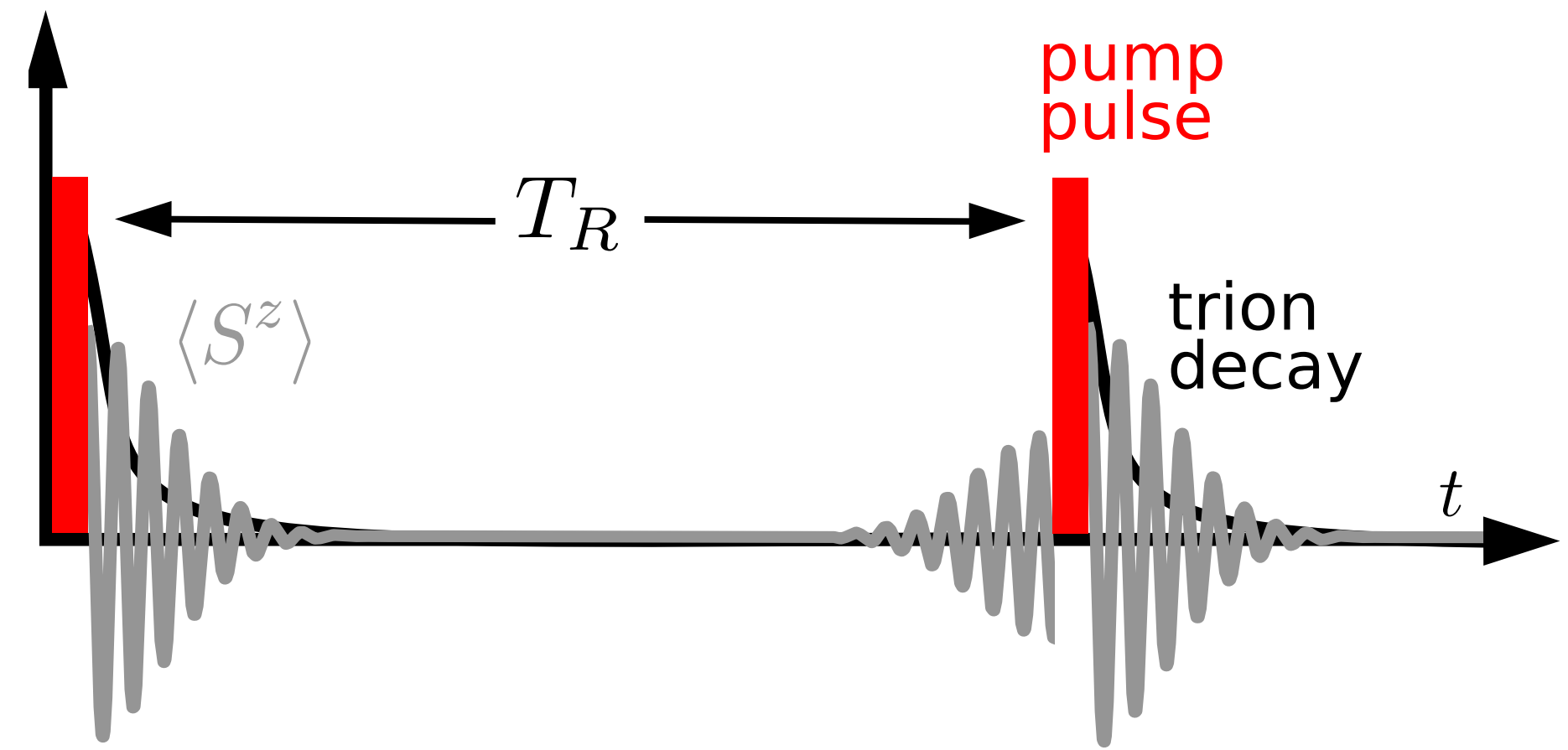
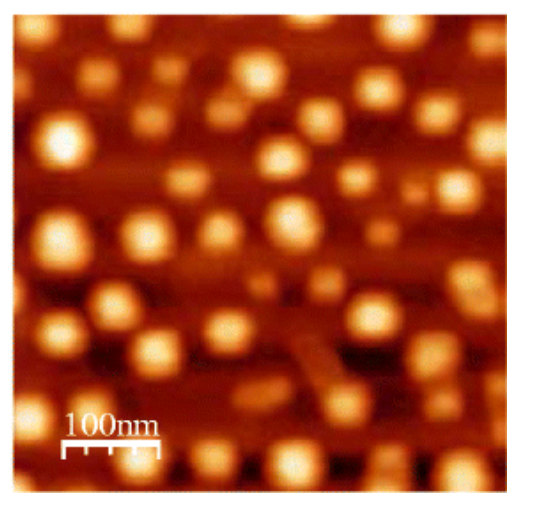
3. higher-order spin noise: information on the coupling to the environment

$\vec{B} = B\vec{e}_x$



Optical pumping of electron spins in Quantum dot ensemble

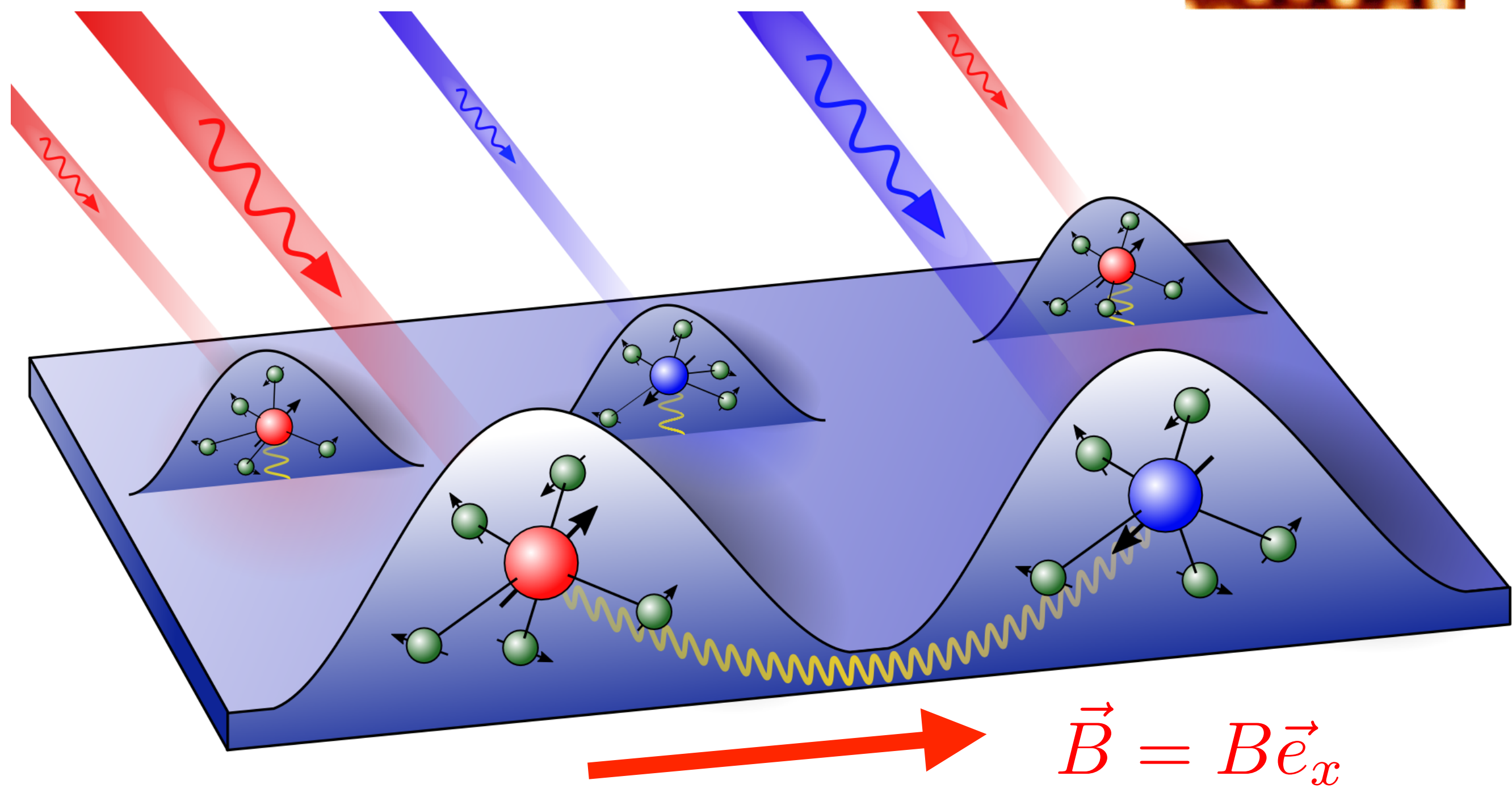
I. Periodic pumping of quantum dot ensembles



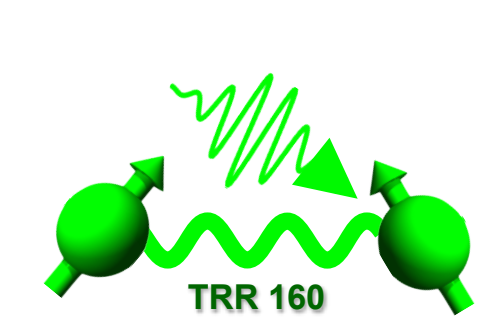
$$\frac{d}{dt} \vec{S}(t) = \left(\vec{B}_N + \vec{B}_{\text{ext}} \right) \times \vec{S}(t) + \gamma P_{T\mu}(0) \vec{e}_z e^{-2\gamma t}$$

nuclear spins

trion decay: S_z

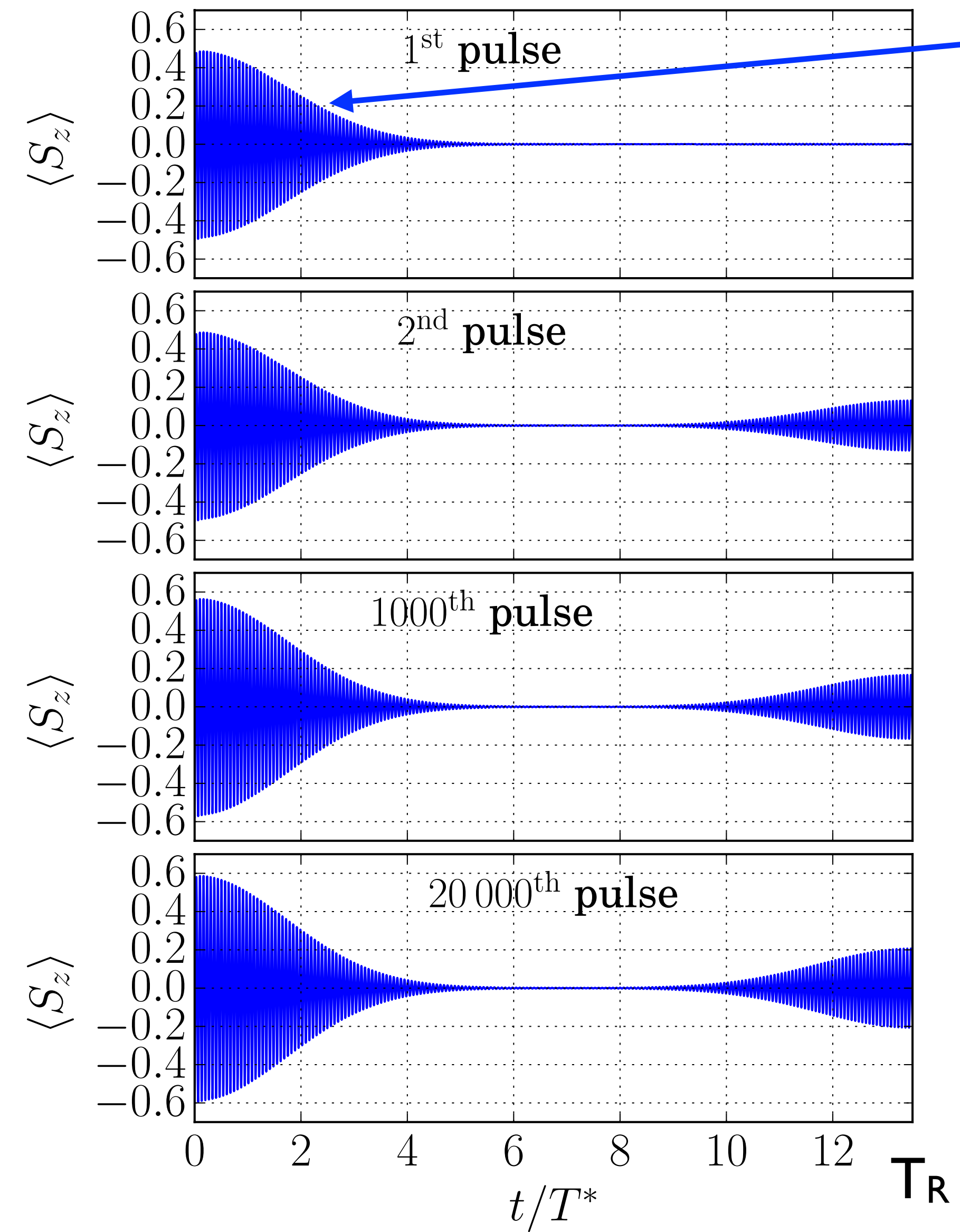


- Larmor precession around an effective magnetic field: **coherence**
- coupling to nuclear spins: **decoherence**
- challenge: combine light-matter interaction and electron-nuclear spin dynamics



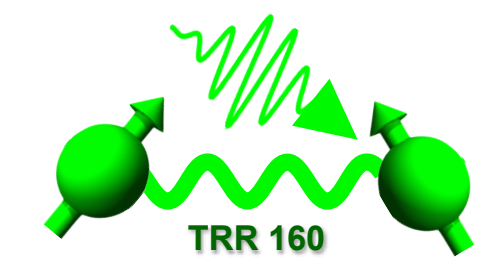
Revival of the spin polarization

decoherence: T^*

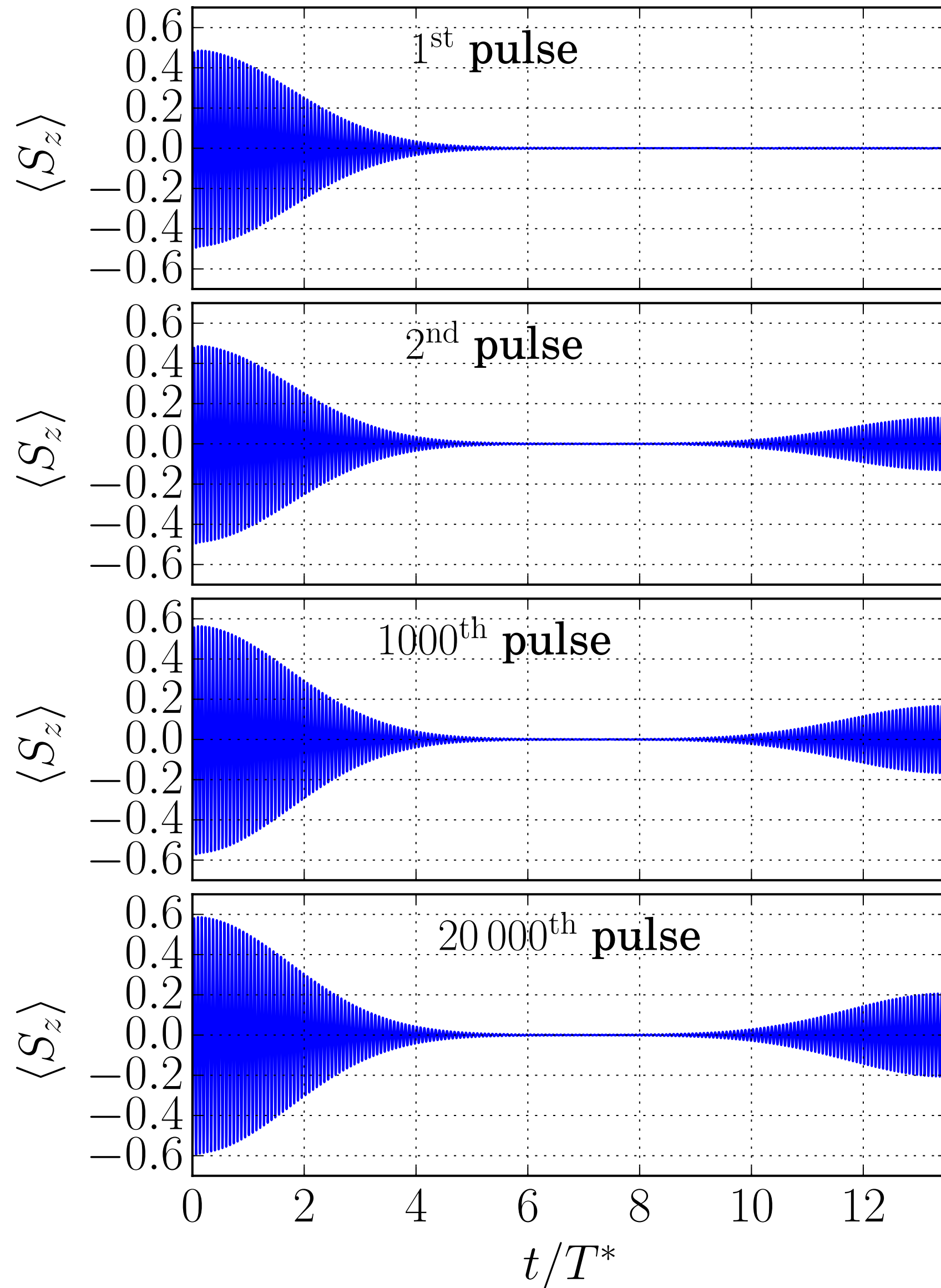


Electron spin synchronization

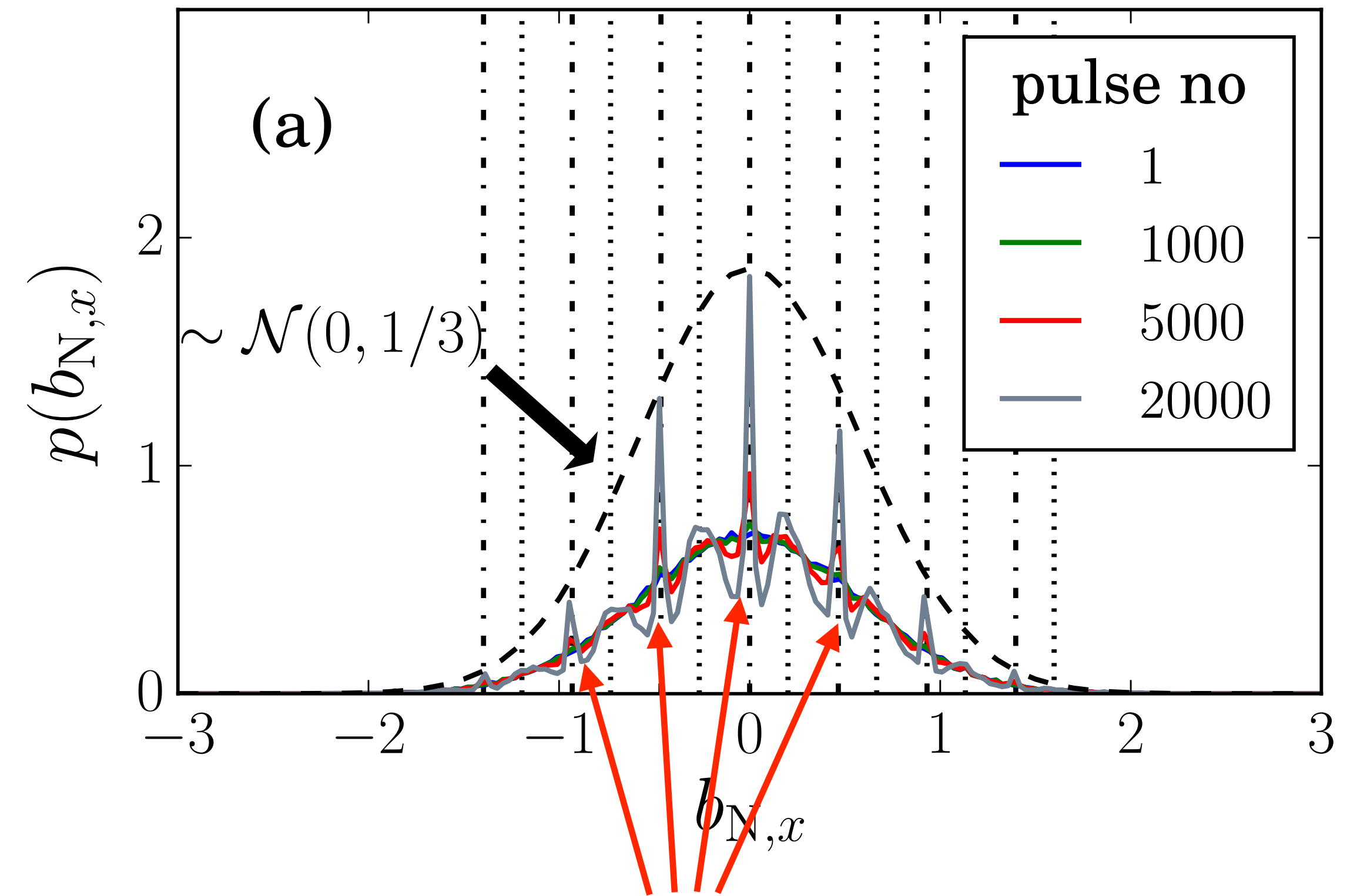
Nuclear spin bath synchronization



Revival of the spin polarization

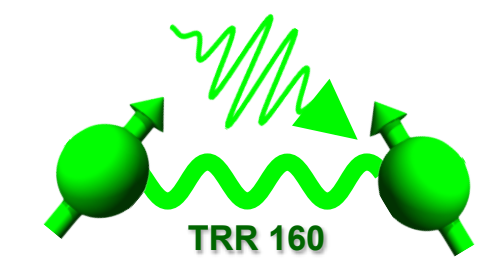


Nuclear magnetic field B_N distribution

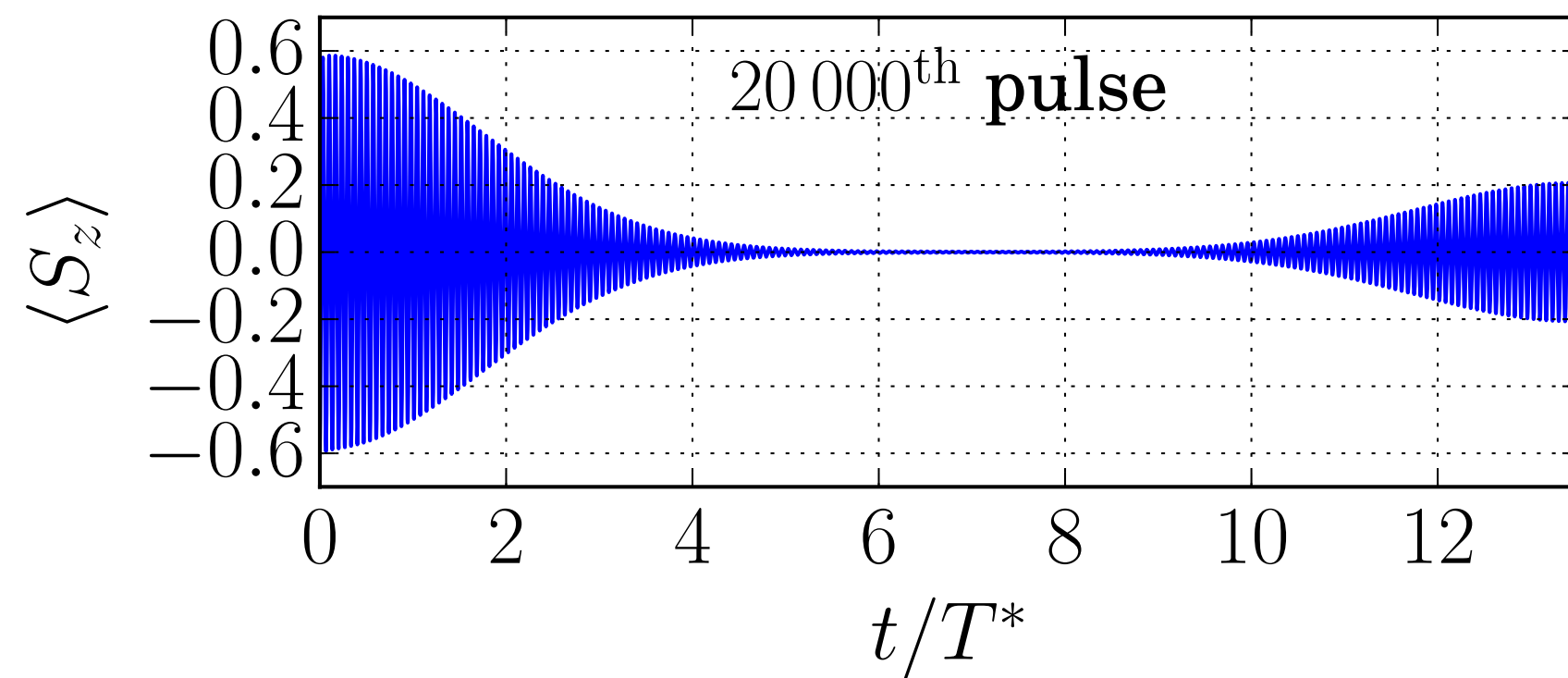


resonance conditions $\omega_L T_R = 2\pi n$
 $\omega_L T_R = \arctan(\omega_L/\gamma) + 2\pi n$

Nuclear spin bath synchronization



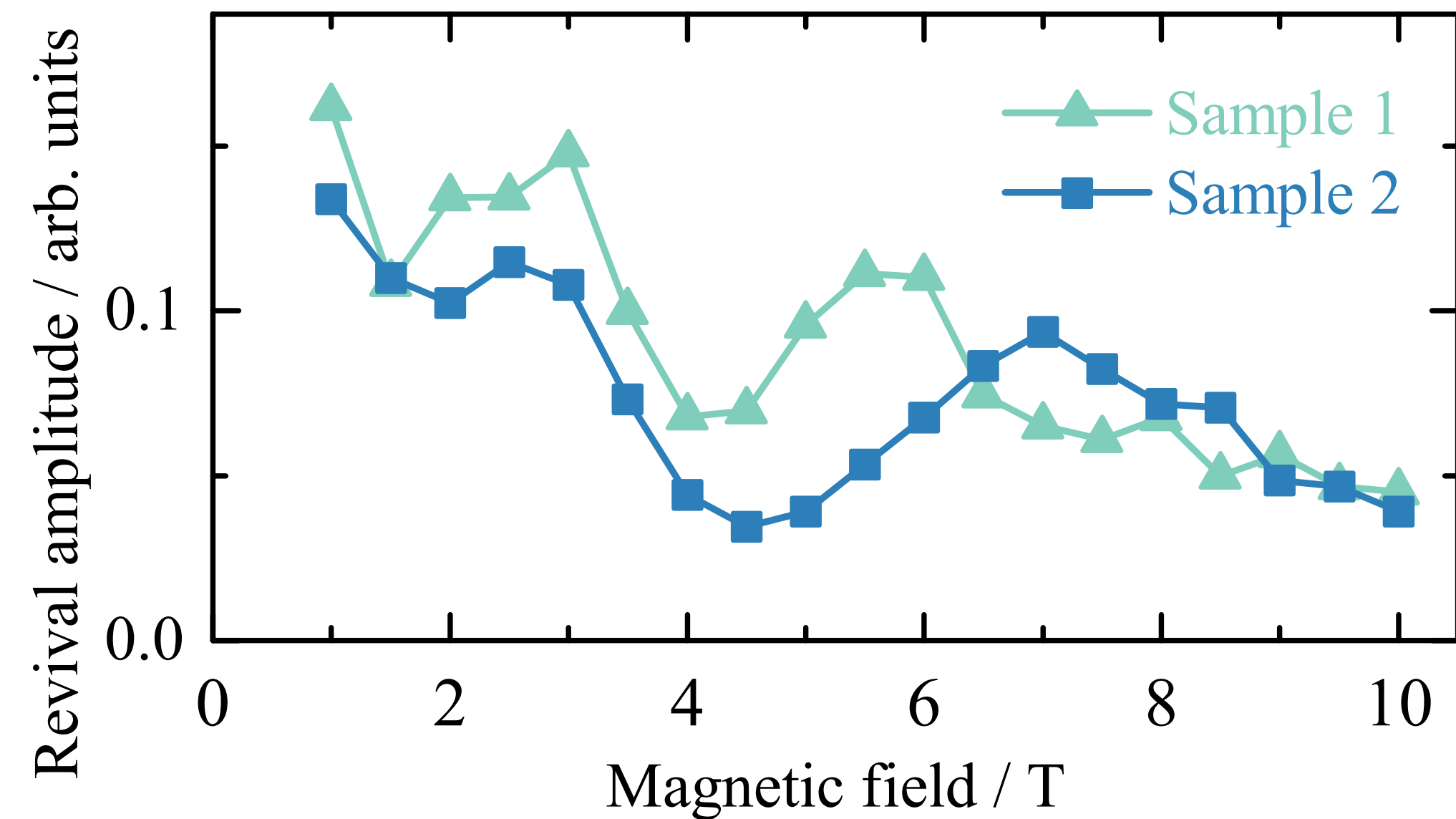
Analysis of the revival amplitude

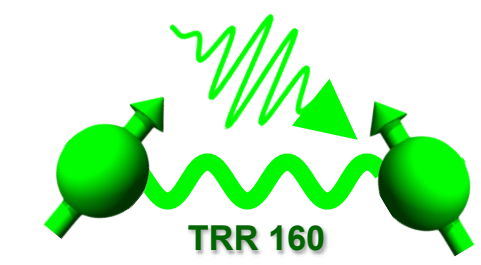


revival amplitude: magnetic field dependent



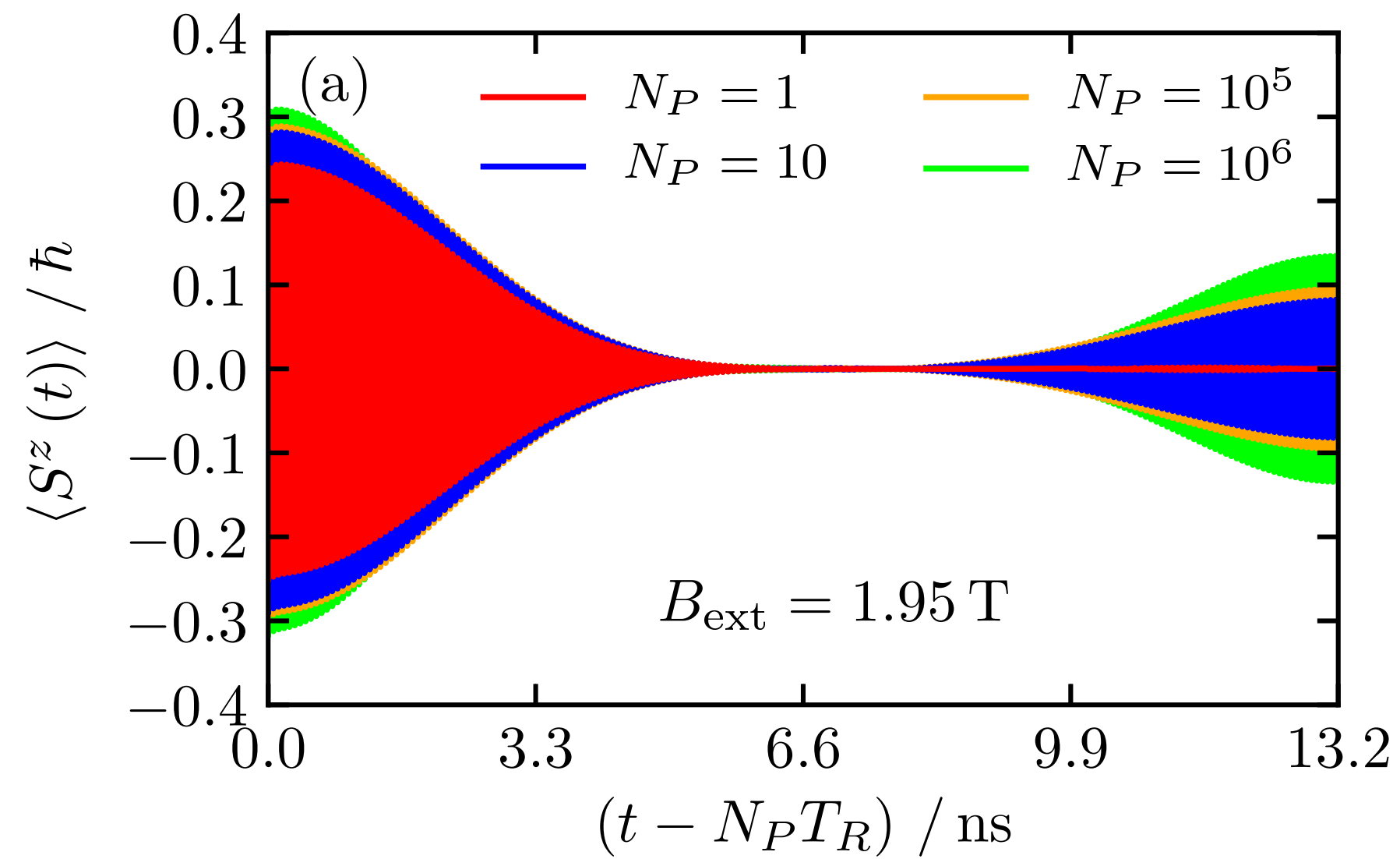
- can we understand the non-linear dependency of the revival amplitude
- what is the role of the nuclear spins?





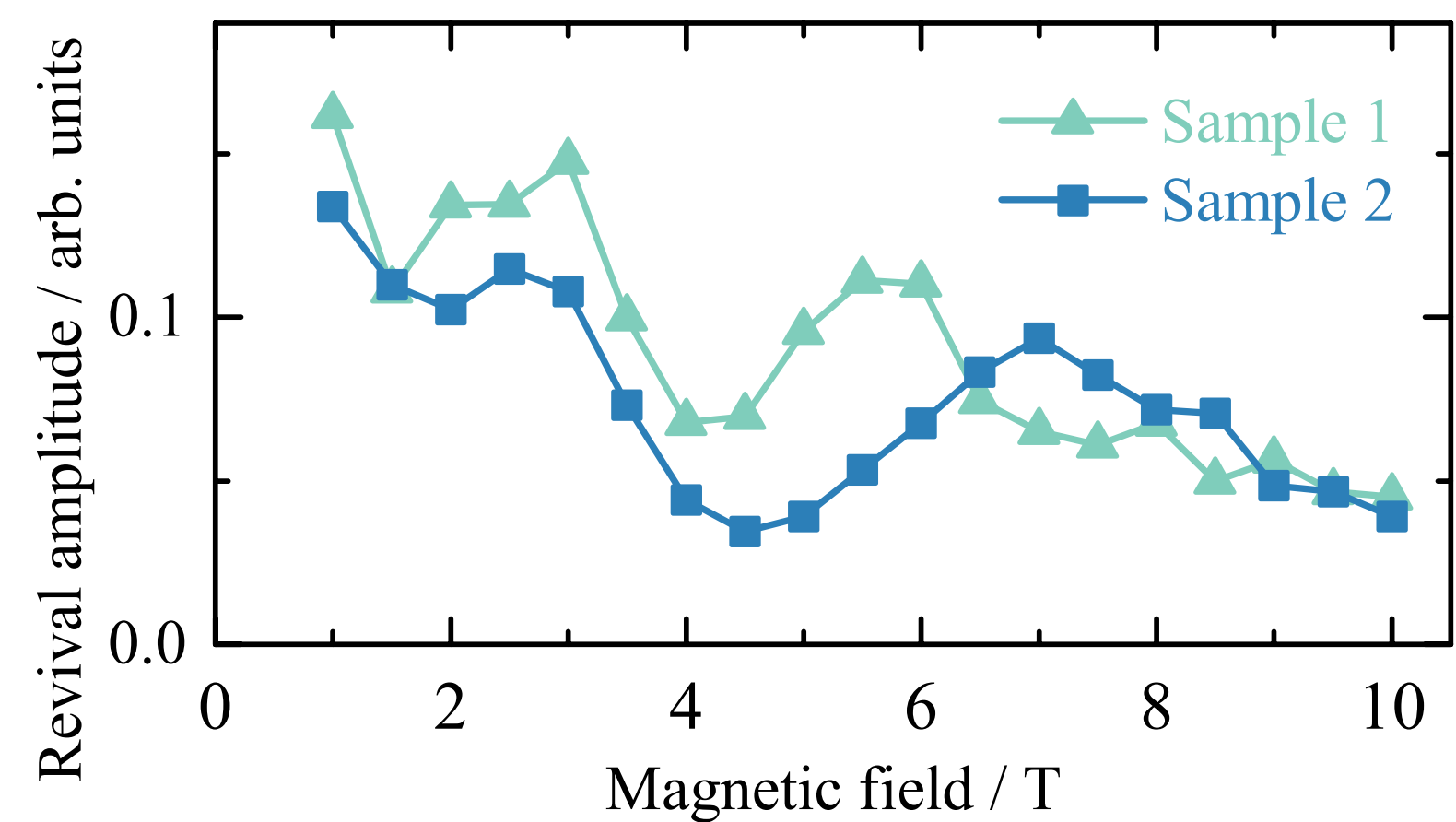
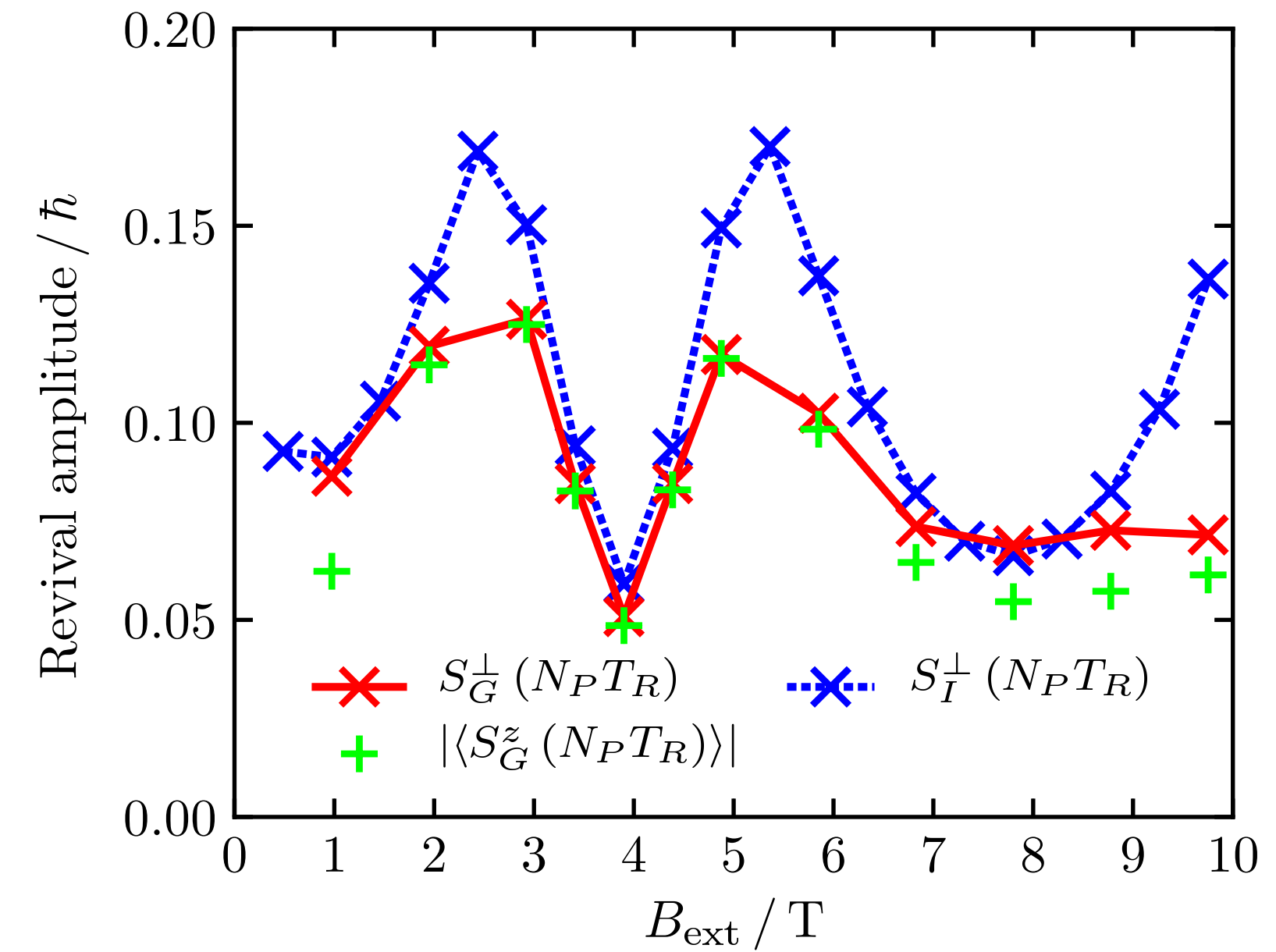
Analysis of the revival amplitude

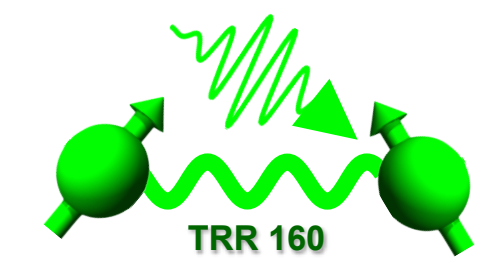
quantum mechanical simulation of 20mil pulses



- minimum: nuclear spin resonance condition
- experiment: different isotope mixtures

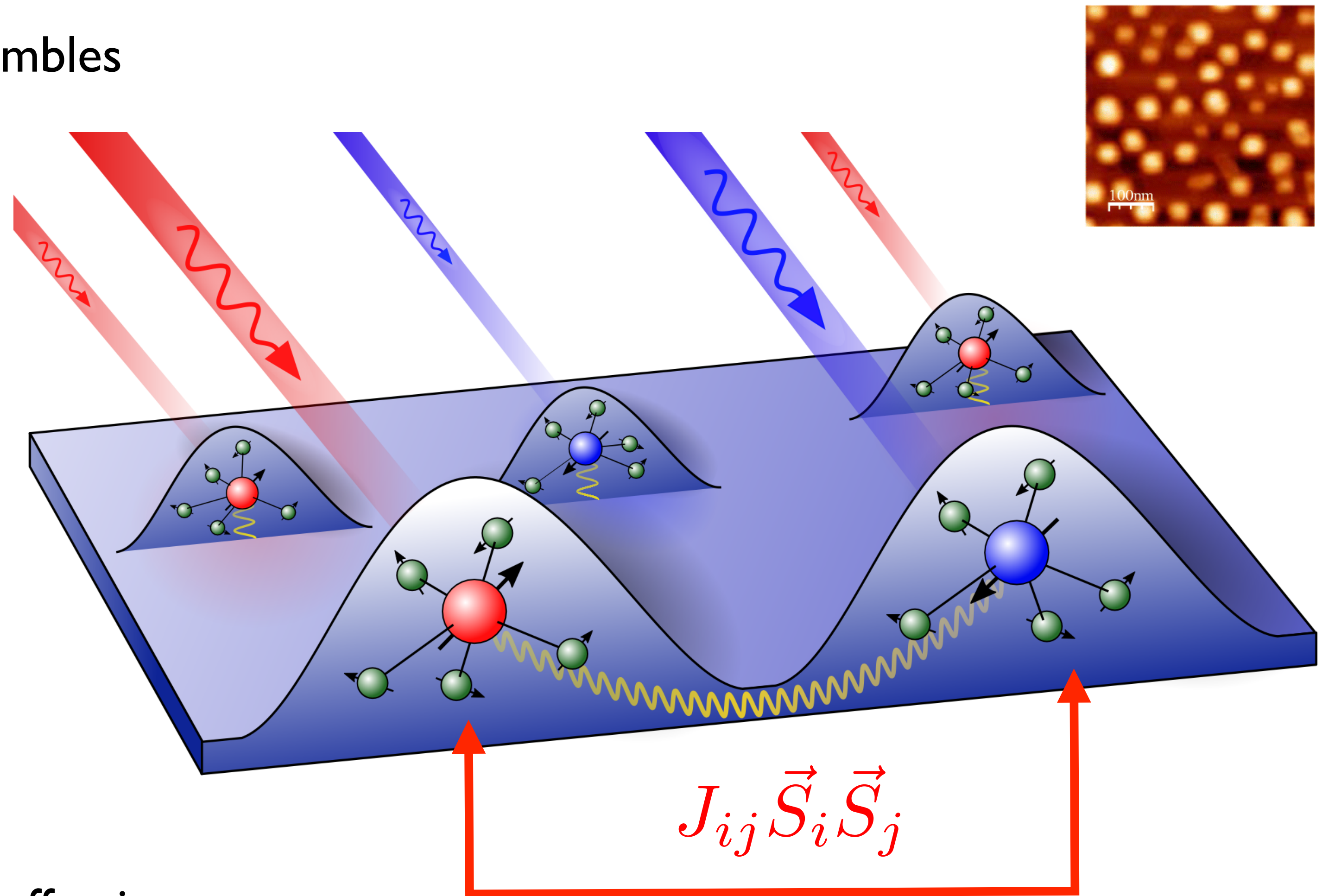
$$\omega_N T_R = \pi n$$





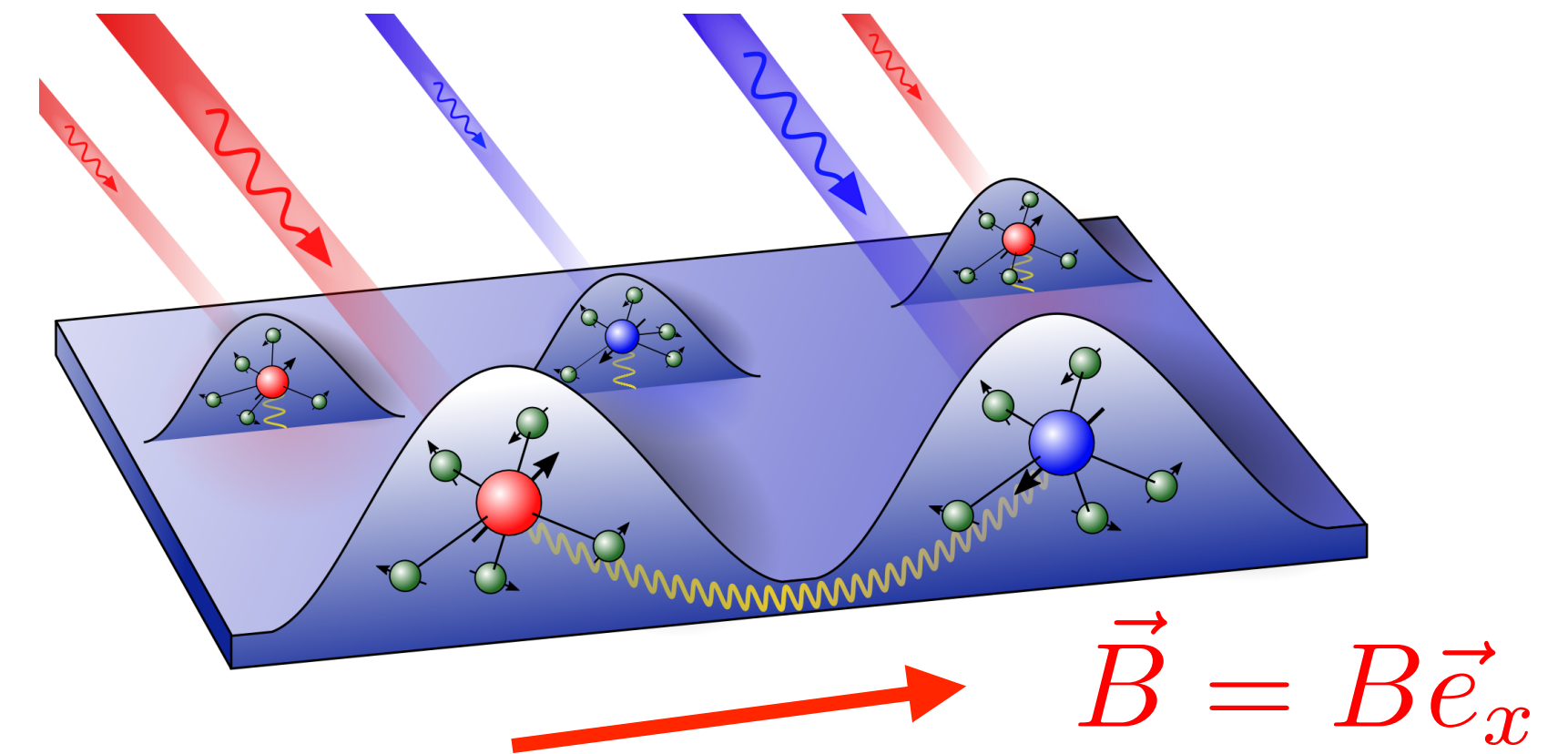
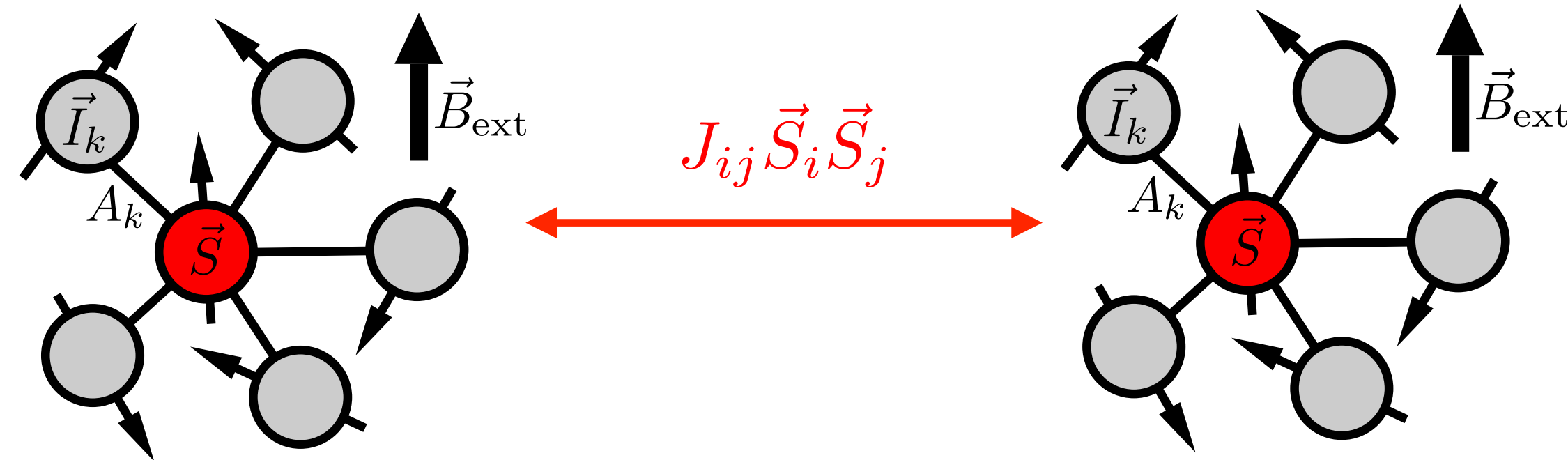
Multi-color pumping: addressing different QDs

- Two color pumping of quantum dot ensembles



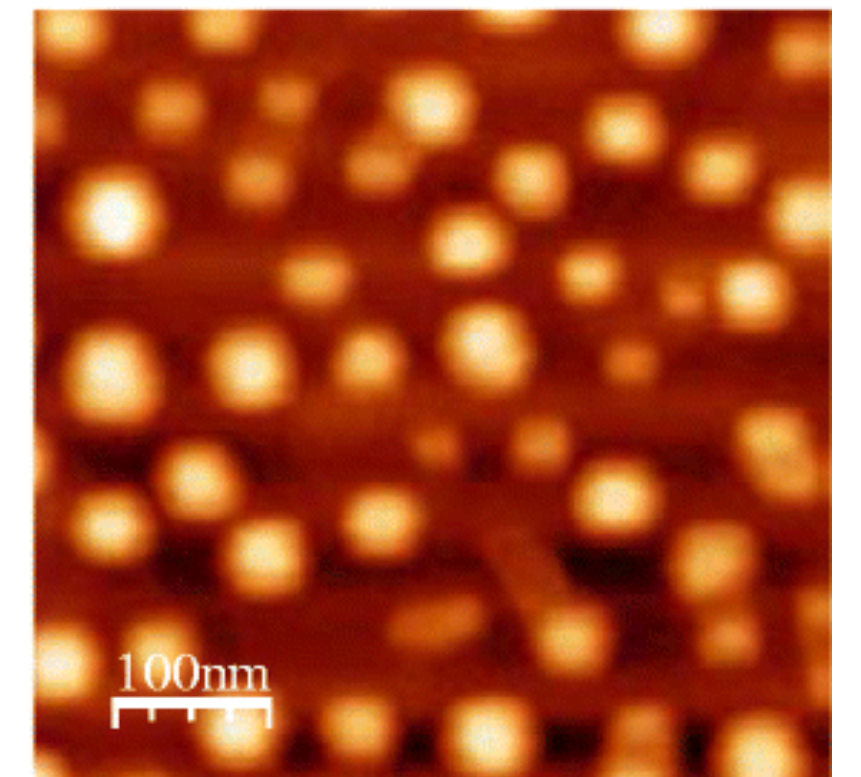
- ➔ experimental data is consistent with an effective Heisenberg coupling between two quantum dots
- ◎ RKKY via wetting layer?

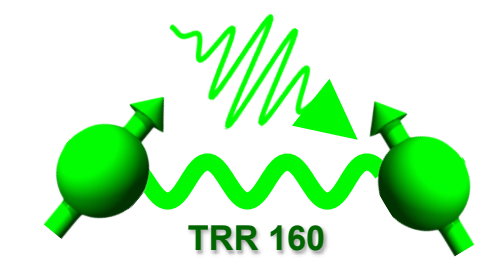
Multi-color pumping: addressing different QDs



$$\frac{d}{dt} \vec{S}_i(t) = \left(\vec{B}_N^i + \vec{B}_{\text{QDs}}^{\text{no-pump}} + \vec{B}_{\text{QDs}}^{\text{pumped}} + \vec{B}_{\text{ext}} \right) \times \vec{S}_i(t) + \gamma P_{T\mu}^i(0) \vec{e}_z e^{-2\gamma t}$$

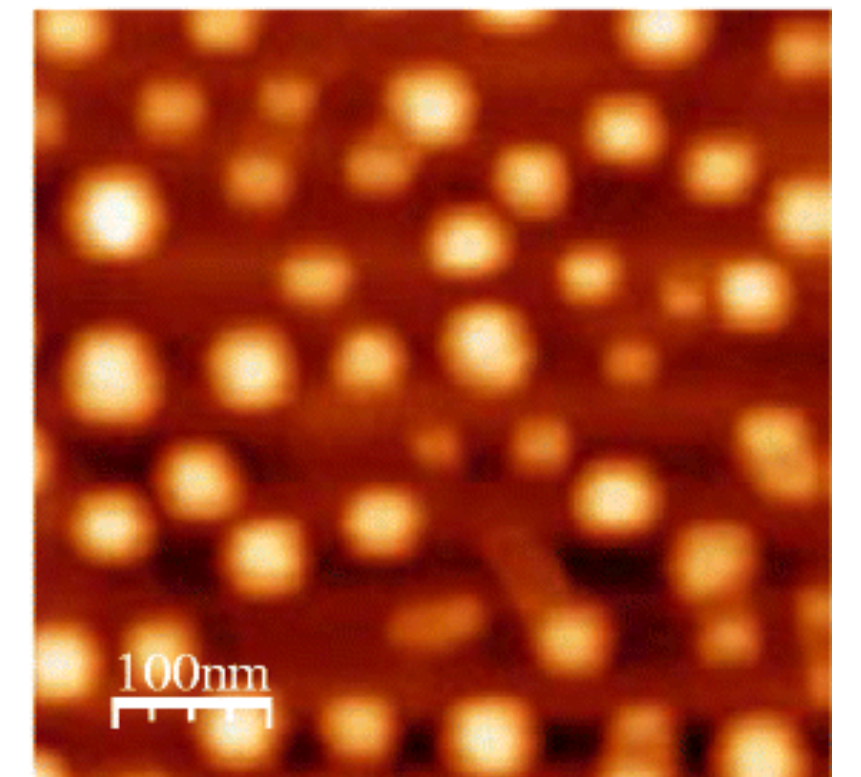
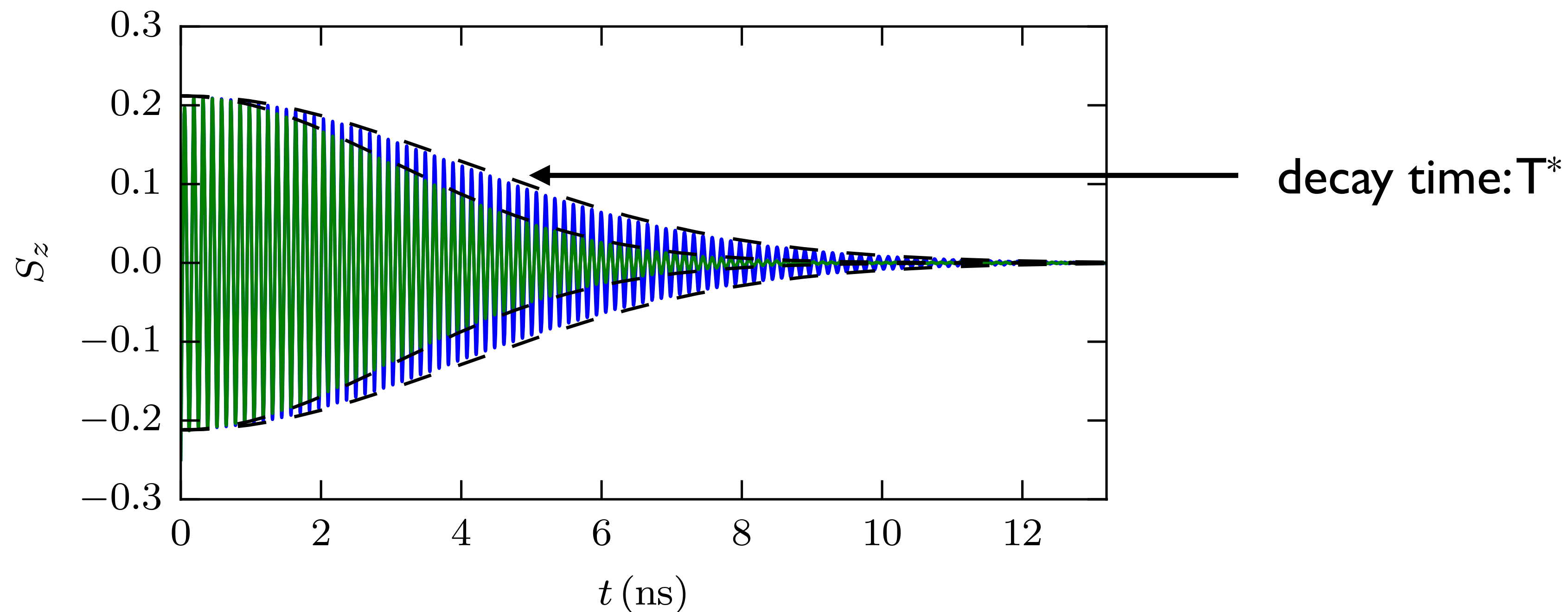
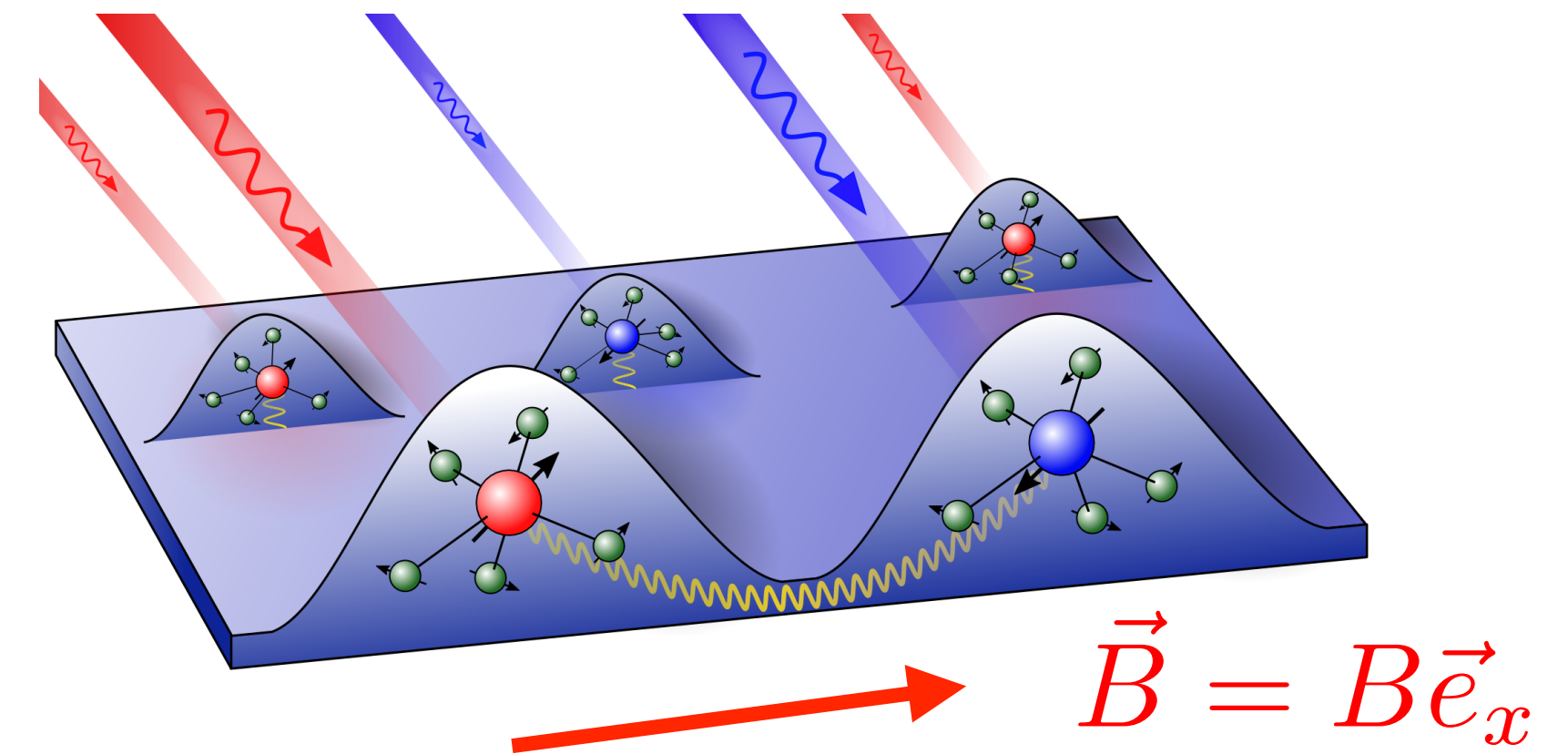
- two **random Gaussian** magnetic fields: the **Overhauser field B_N** and **$B_{\text{QDs}}^{\text{no-pumped}}$**
- coupled semiclassical differential equations plus QM light-matter interaction: massive parallelized on JURECA booster Intel Phi architecture

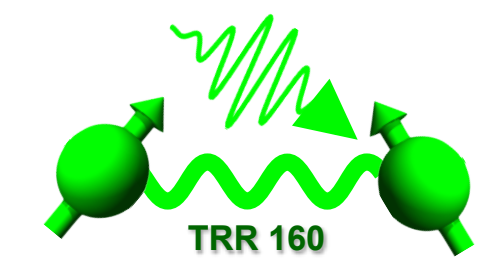




Multi-color pumping: addressing different QDs

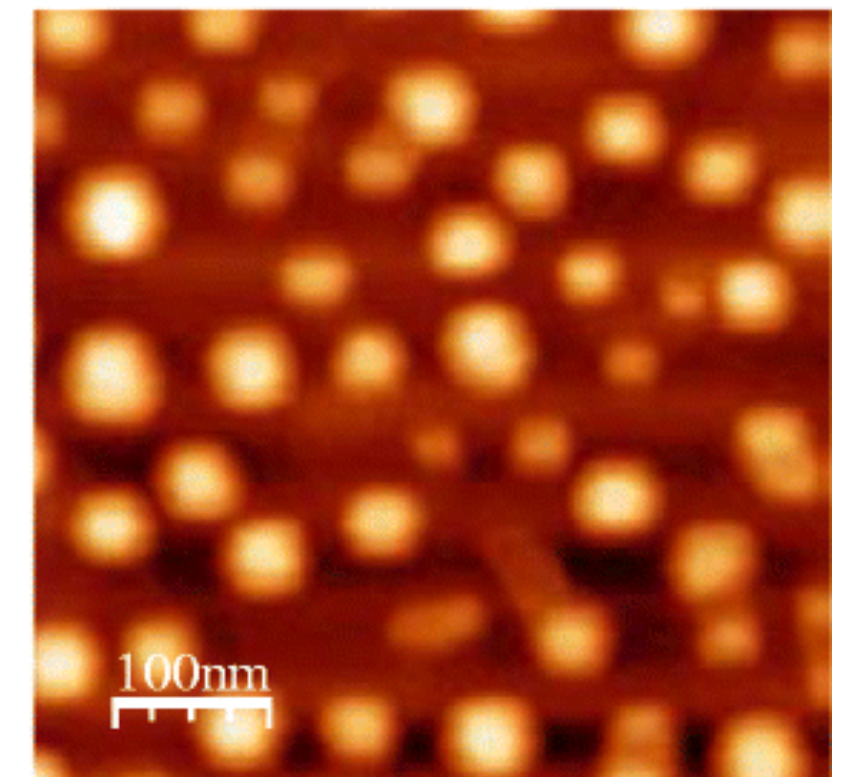
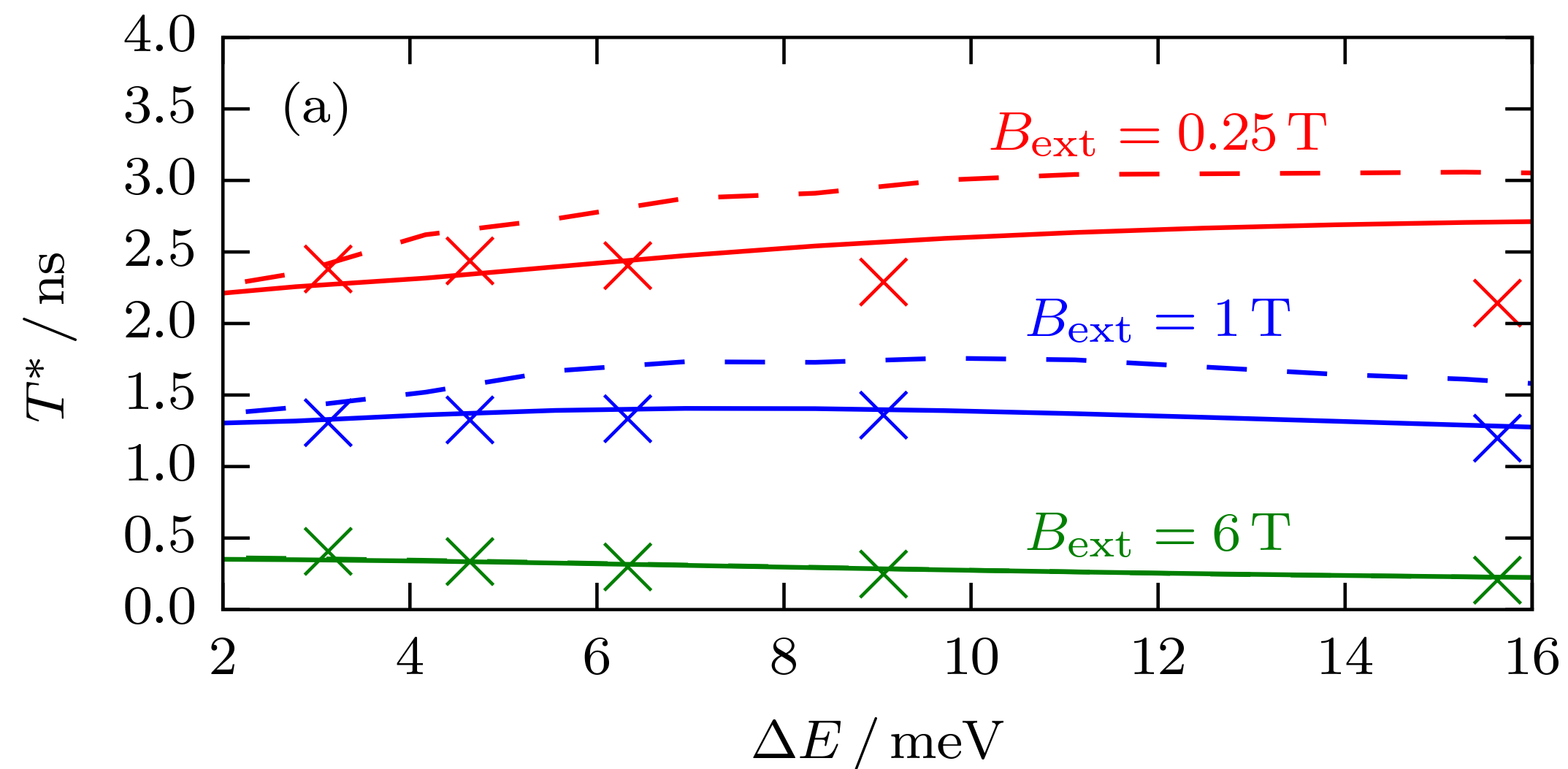
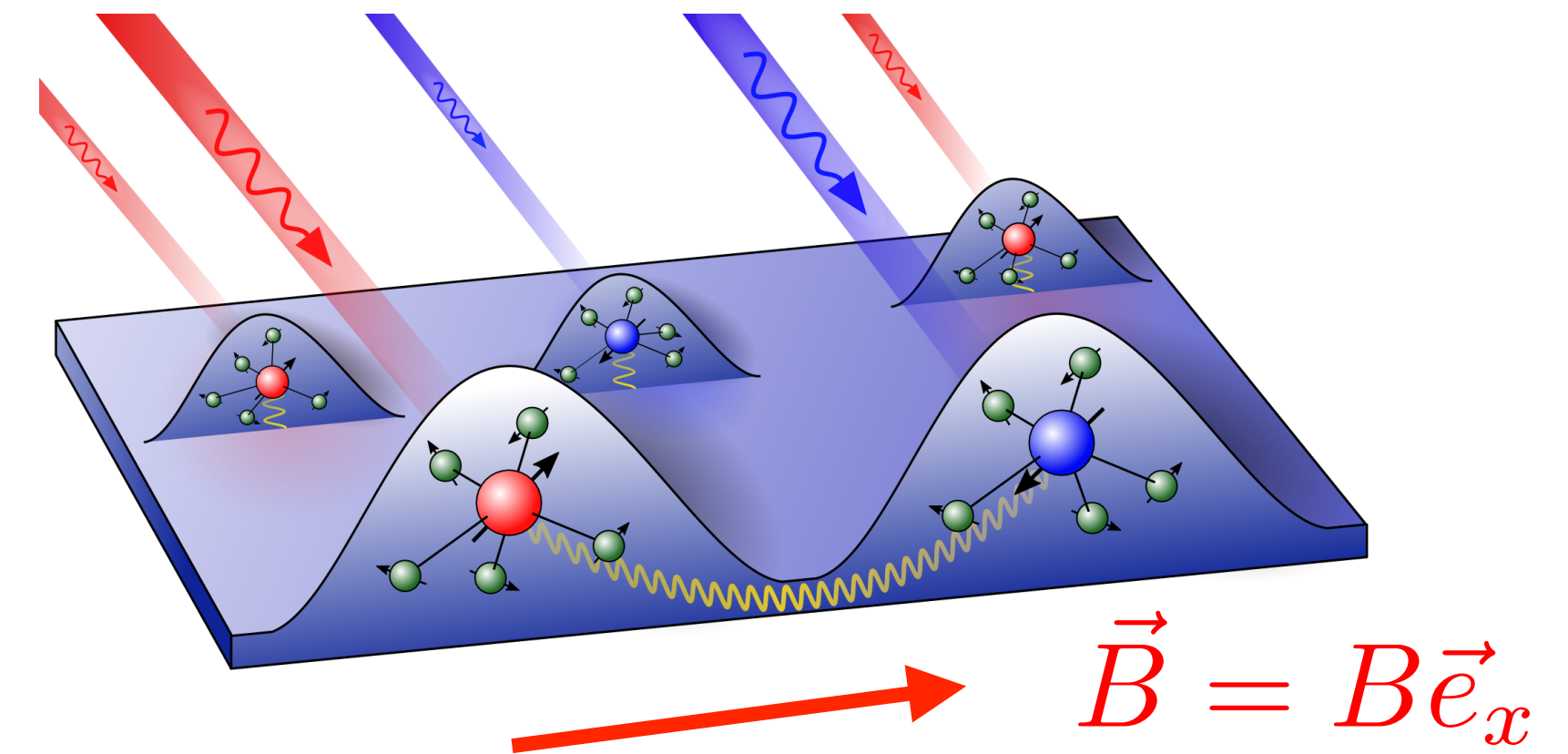
- resonance energy depends on the random shape of the QD
→ **sharp** laser frequency: **pumping** of only **one sub-ensemble**
- finite pulse:
 - ✓ finite width ΔE
 - ✓ pumping several sub-ensembles
 - ✓ synchronisation vs dephasing of unpumped QDs





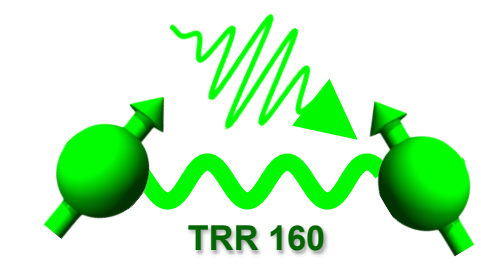
Multi-color pumping: addressing different QDs

- resonance energy depends on the random shape of the QD
→ sharp laser frequency: pumping of only one sub-ensemble
- finite pulse:
 - ✓ finite width ΔE
 - ✓ pumping several sub-ensembles
 - ✓ synchronisation vs dephasing of unpumped QDs



Fischer et al, PRB, Nov 2018

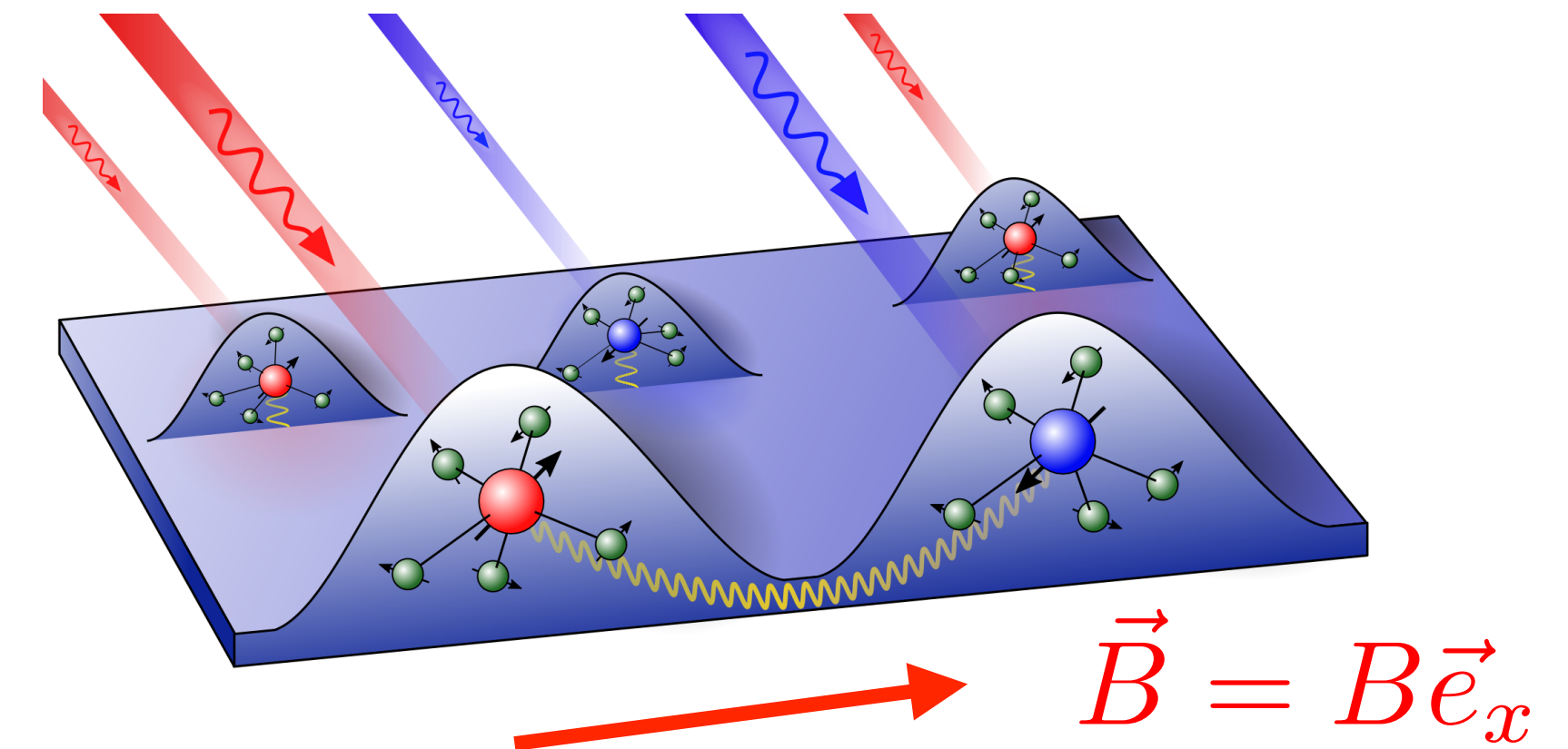
JURECA BOOSTER: Intel Phi



Multi-color pumping: addressing different QDs

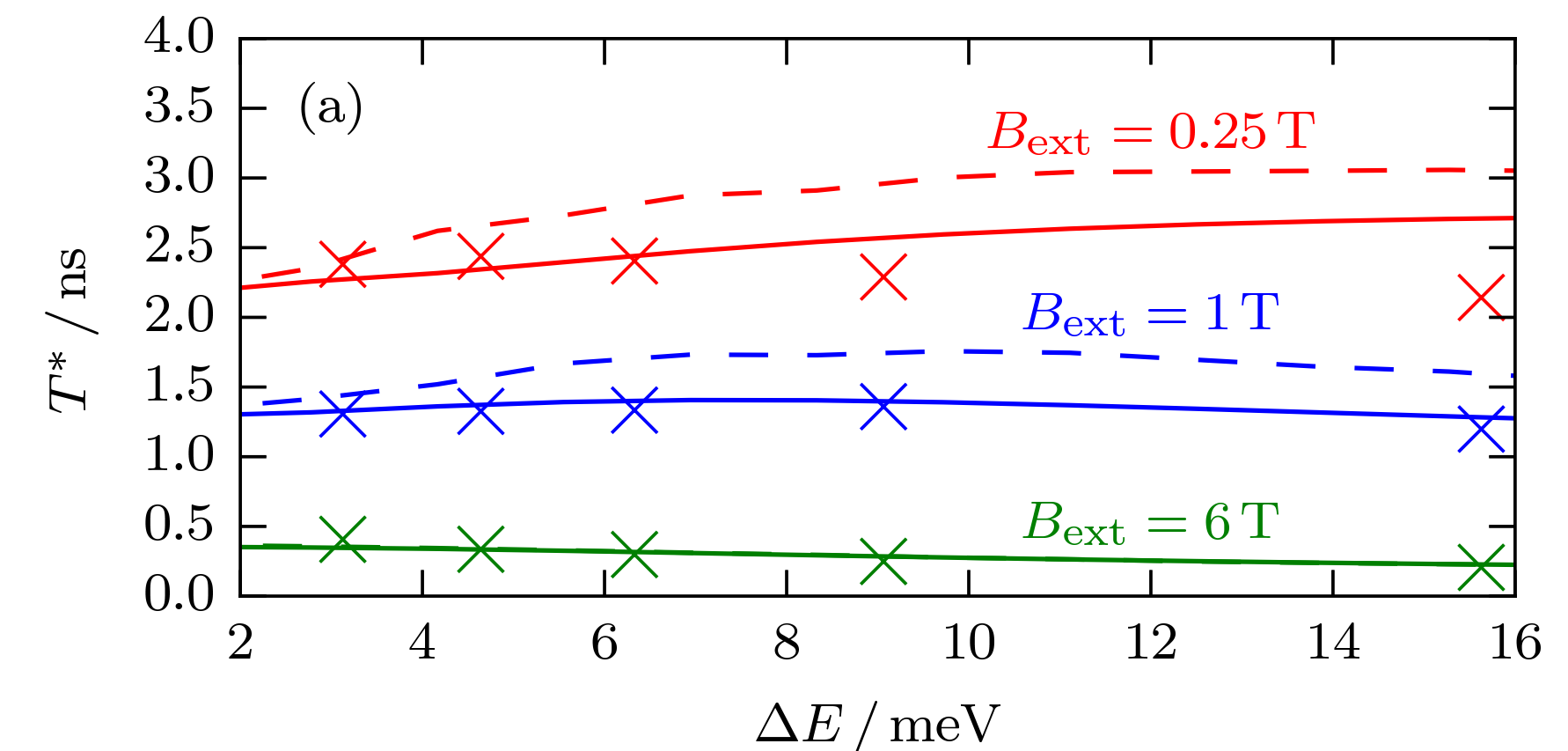
Summary of the multi-color pumping

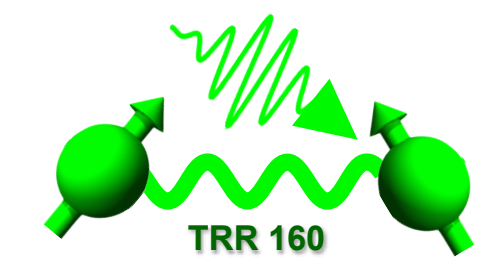
- surprising experimental finding: **decoherence time independent** of the LASER energy width ΔE
- simulation: consistent with spin-coupled QD spins



• Questions

- origin of the coupling?
- can we learn to manipulate this coupling?
- coupling of different subset: qubit interactions?
- different pulse sequences





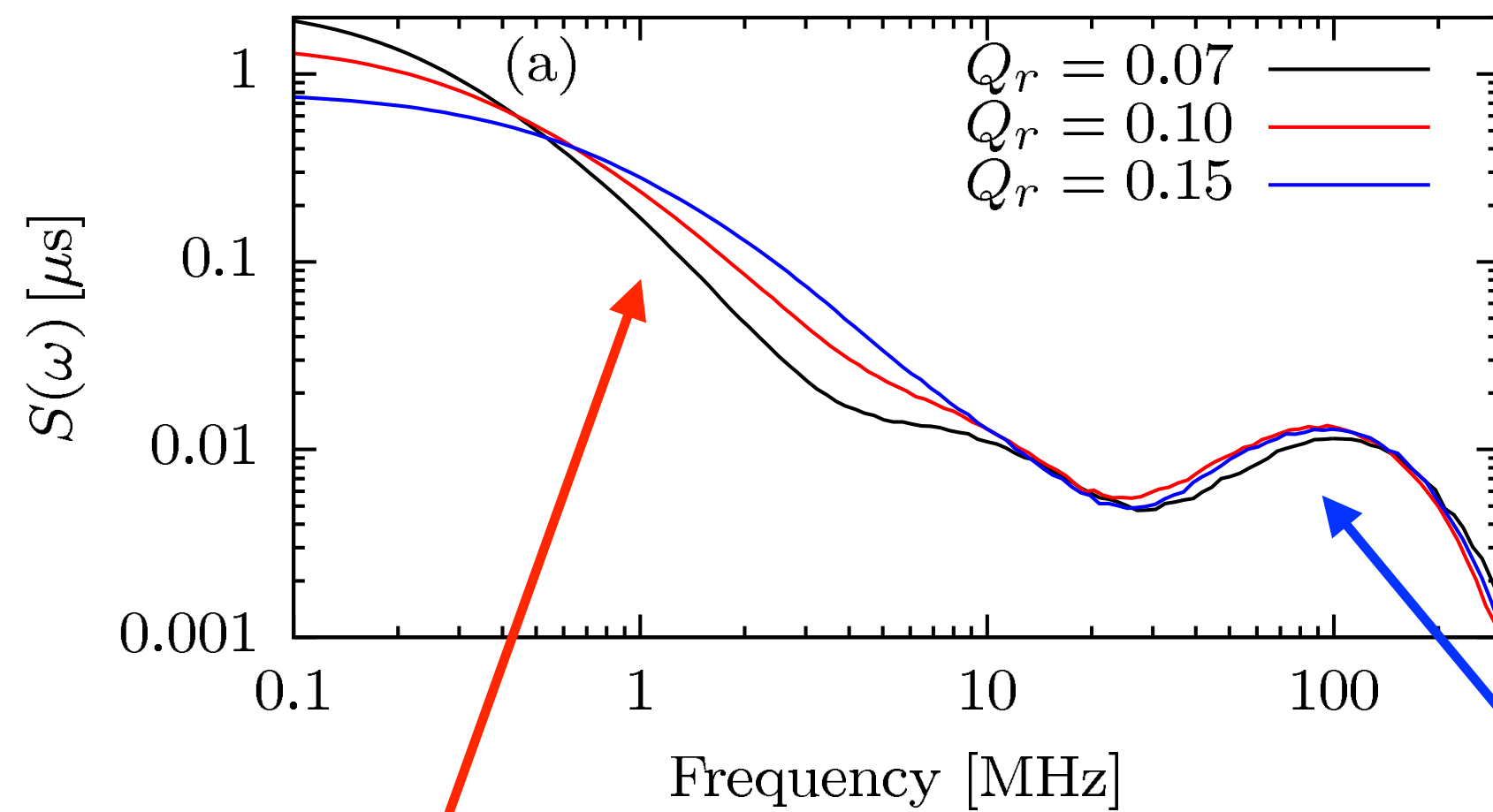
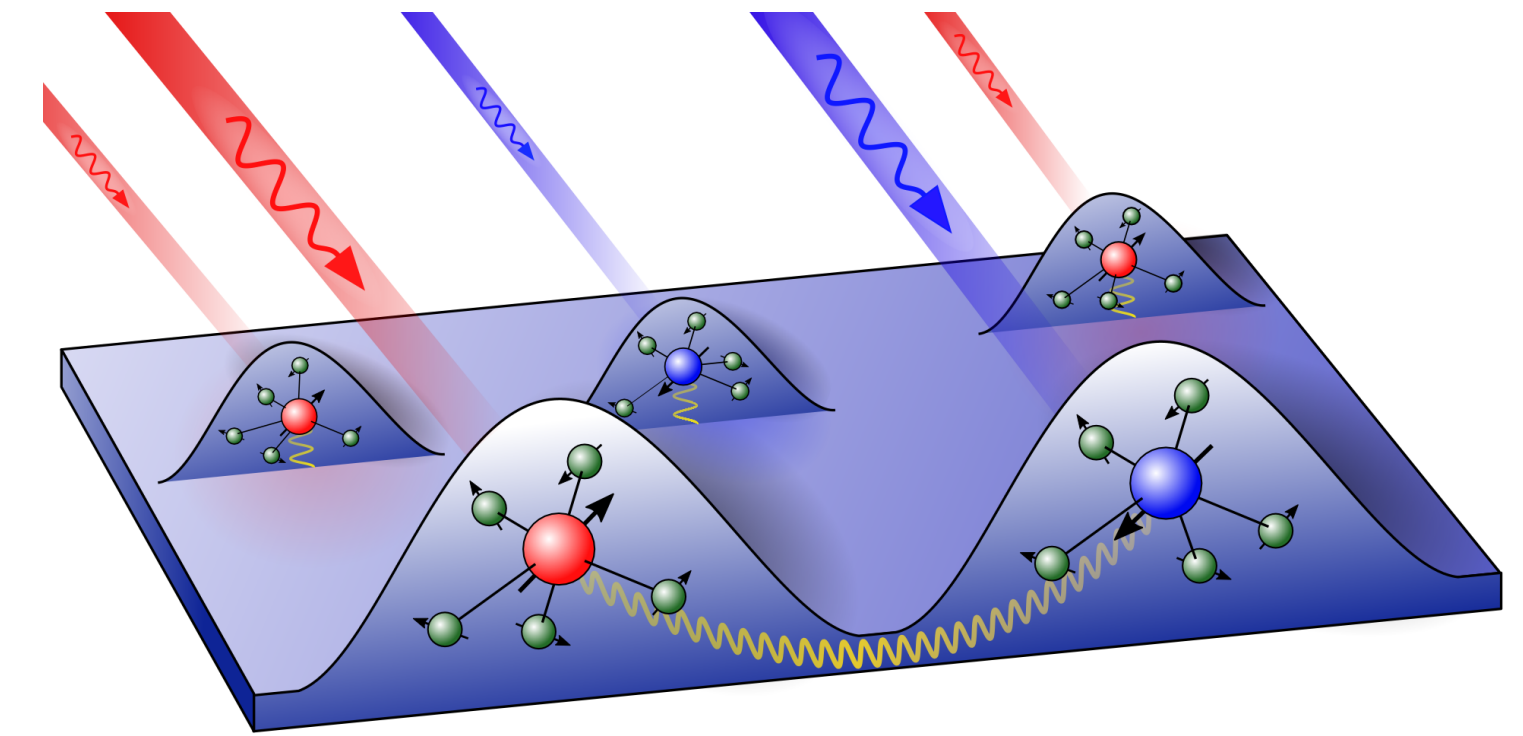
Higher-order spin noise in QDs

Spin noise

$$C_2(t) = \text{Tr} [\rho S_z(t) S_z]$$

- spin noise: Fourier transformation of $C_2(t)$

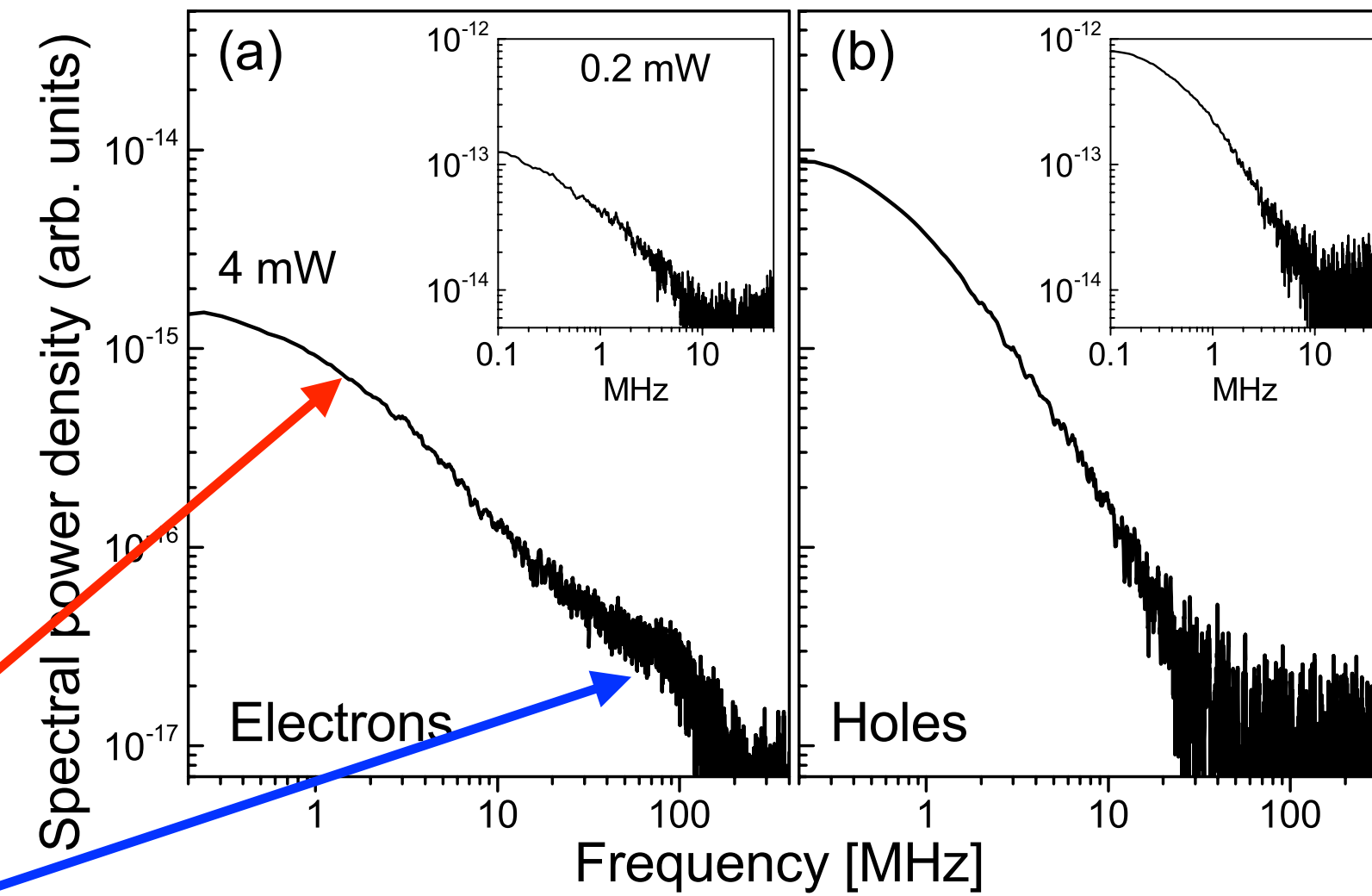
$$S(\omega) = C_2(\omega) = \langle |S_z(\omega)|^2 \rangle$$



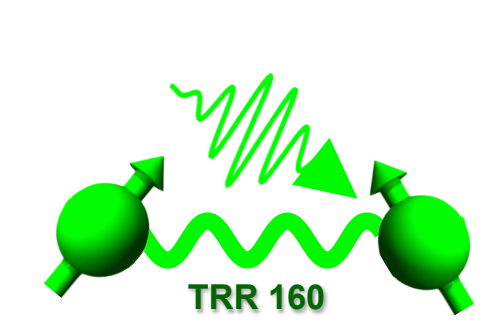
theory

Gaussian nuclear spin fluctuations

nuclear electric quadrupolar interaction



experiment



Higher-order spin noise in QDs

Spin noise

$$C_2(t) = \text{Tr} [\rho S_z(t) S_z]$$

- spin noise: Fourier transformation of $C_2(t)$

$$S(\omega) = C_2(\omega) = \langle |S_z(\omega)|^2 \rangle$$

- message: $C_2(t)$ is well understood

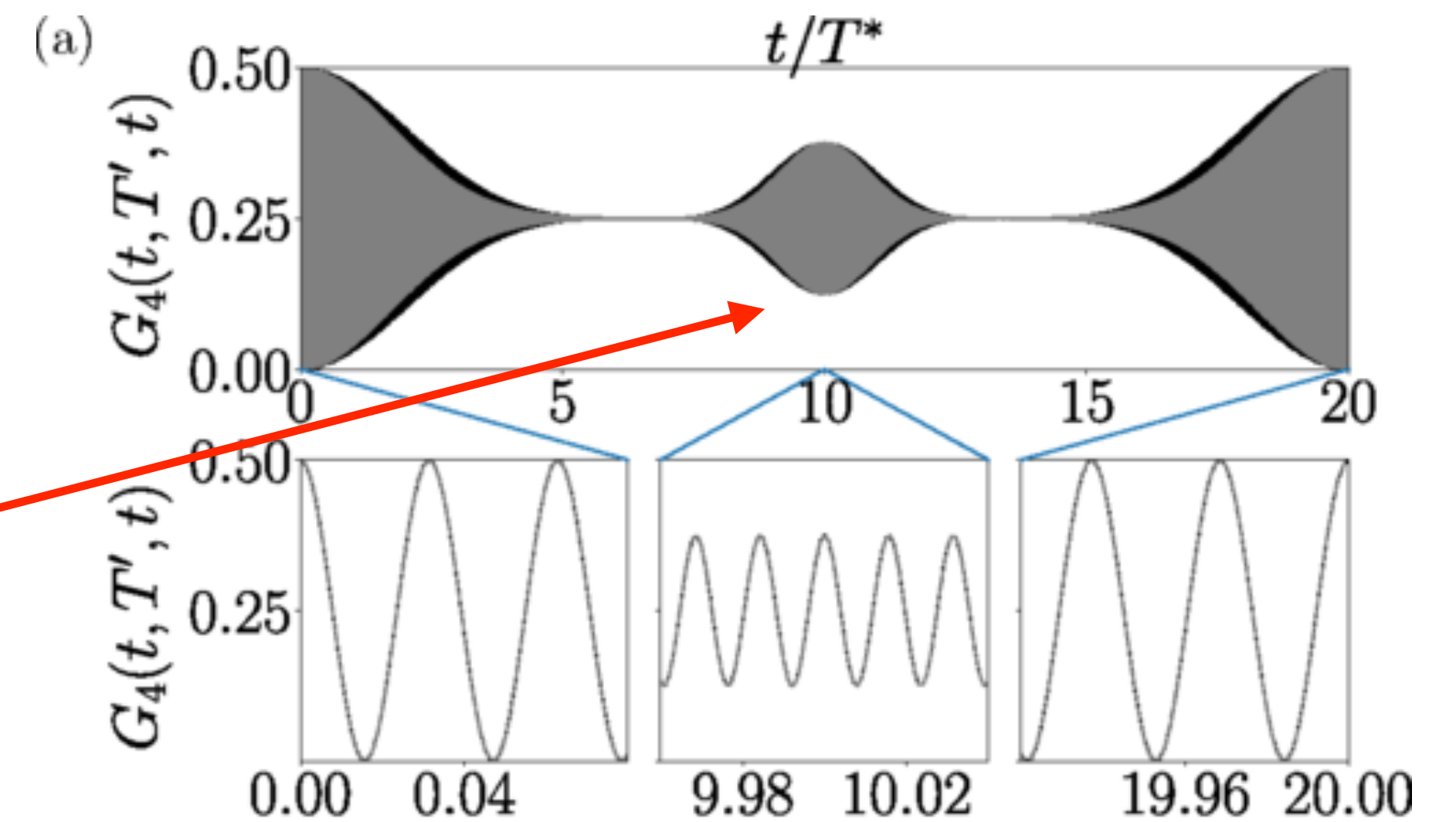
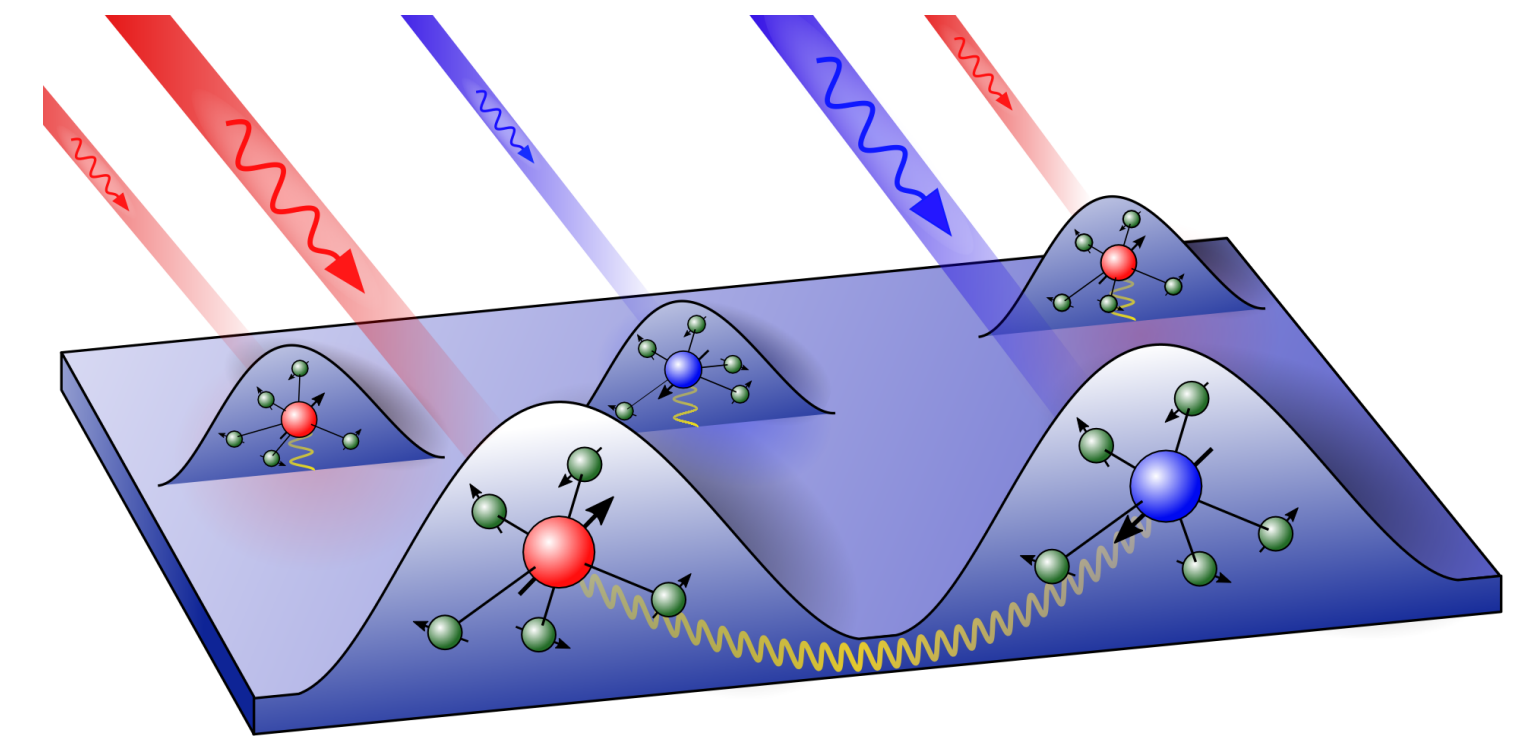
Higher order spin correlations

$$C_4(t_1, t_2, t_3) = \langle S_z(t_1) S_z(t_2) S_z(t_3) S_z \rangle$$

1. spin echo $C_4(t_1, t_1 + t_2, t_1)$

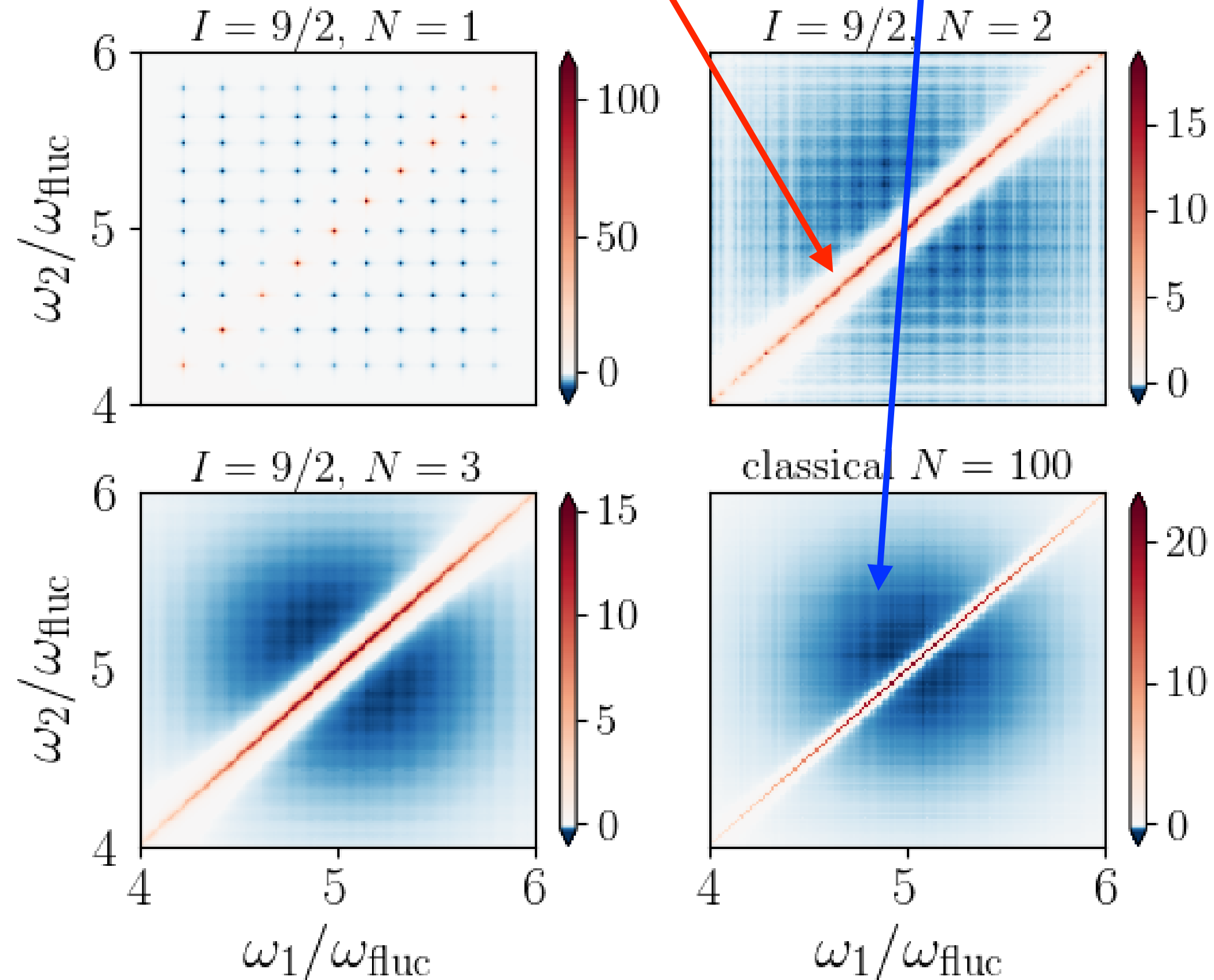
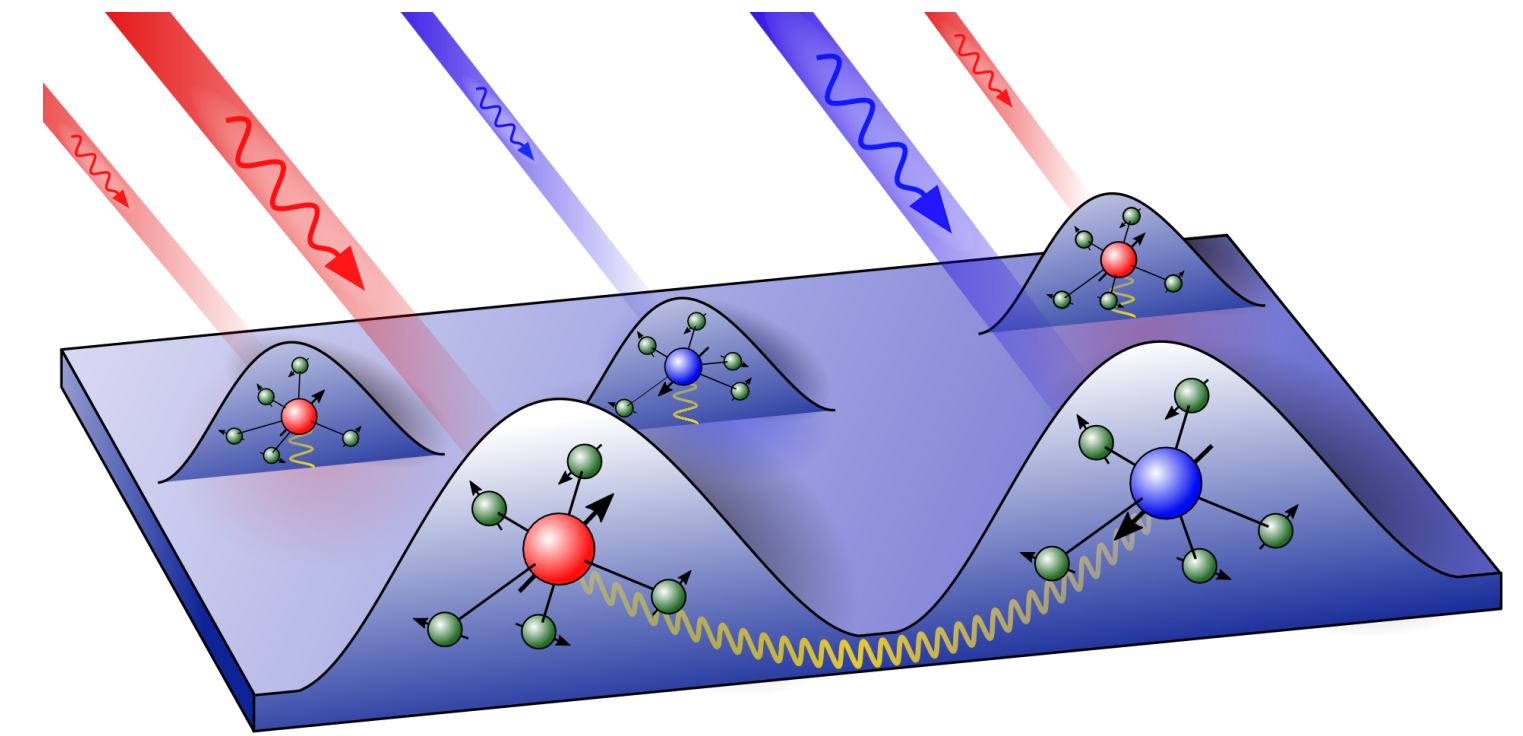
2. 4th-order spin noise

$$C_4(\omega_1, \omega_2) = \langle |S_z(\omega_1)|^2 |S_z(\omega_2)|^2 \rangle = \text{FT} \left[\lim_{T_m \rightarrow \infty} \frac{1}{T_m} \int_{-T_m}^{T_m} d\tau C_4(t_1 + \tau, \tau, t_2) \right]$$



4th-order spin noise

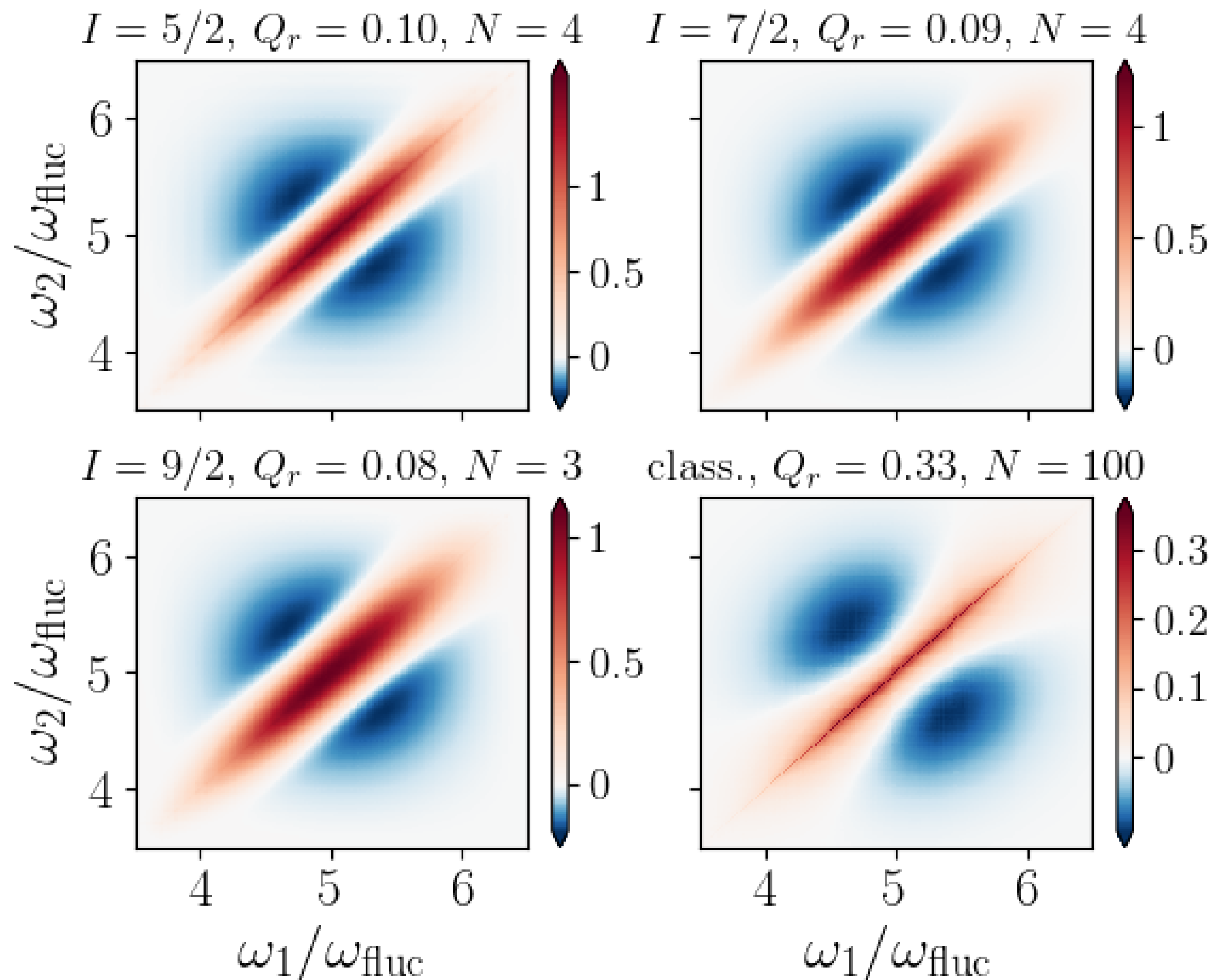
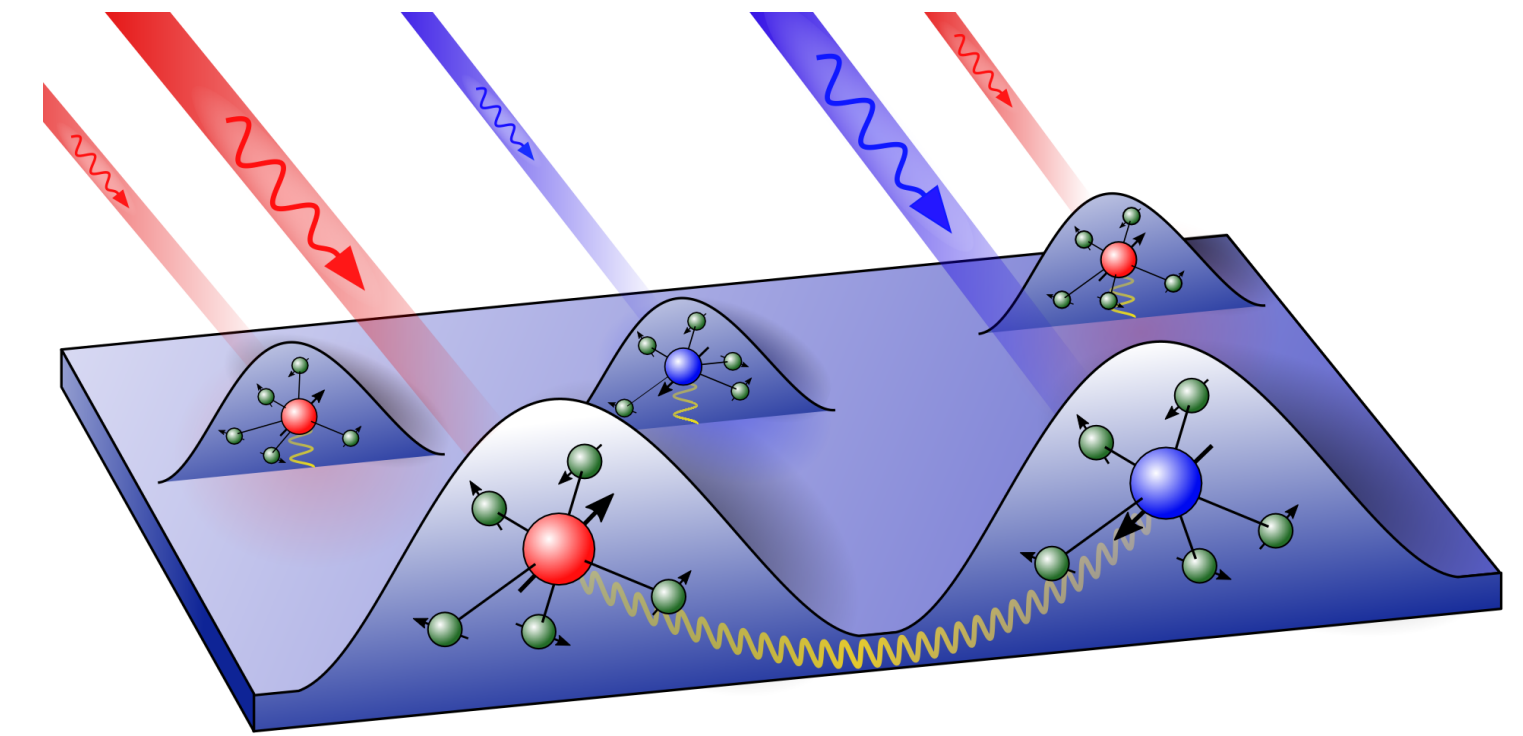
$$S_4(\omega_1, \omega_2) = C_4(\omega_1, \omega_2) - C_2(\omega_1)C_2(\omega_2)$$



- crossover from a single spin to a classic continuum
- 4th-order: only correlation on the diagonal for a simple model
- finite magnetic field: Gaussian anti-correlations

4th-order spin noise: include quadrupolar interactions

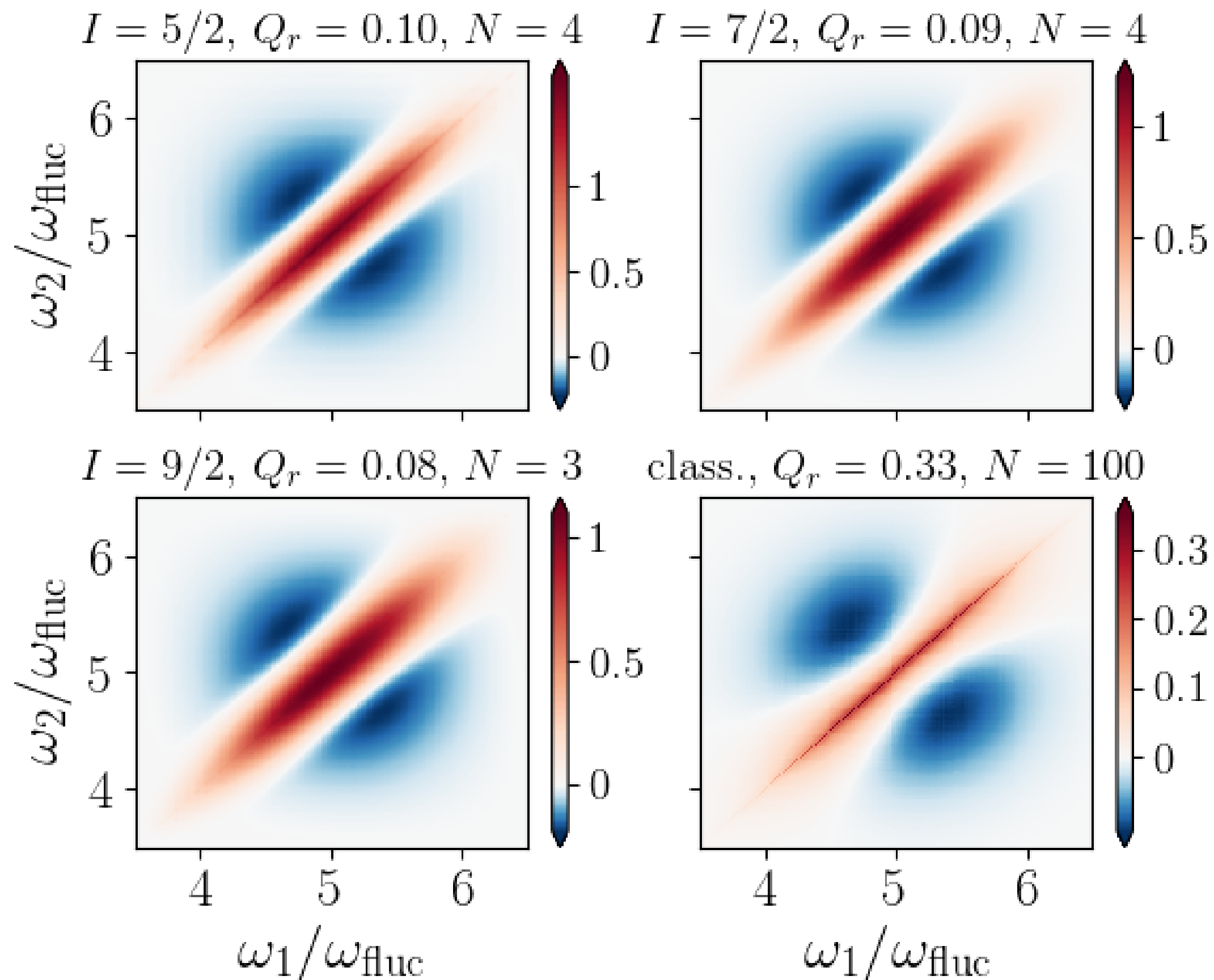
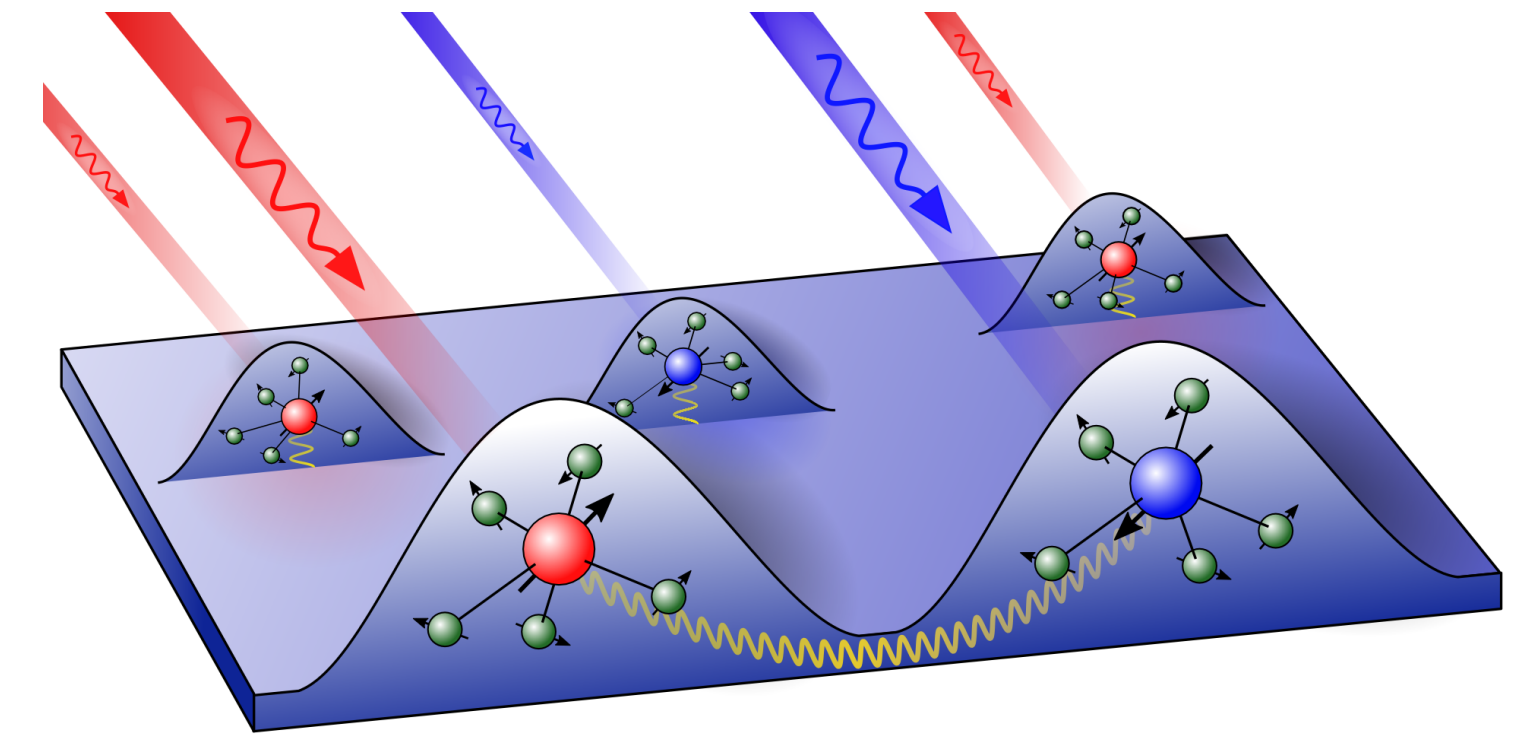
$$S_4(\omega_1, \omega_2) = C_4(\omega_1, \omega_2) - C_2(\omega_1)C_2(\omega_2)$$



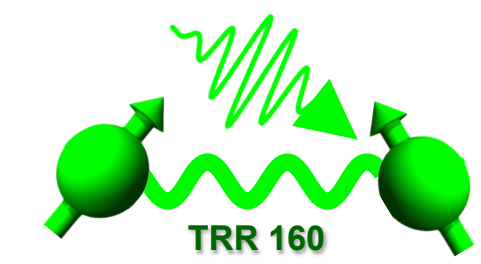
- off-diagonal noise: probe for additional interactions that causes decoherence
- modification of the Gaussian noise

4th-order spin noise: include quadrupolar interactions

$$S_4(\omega_1, \omega_2) = C_4(\omega_1, \omega_2) - C_2(\omega_1)C_2(\omega_2)$$



- off-diagonal noise: probe for additional interactions that causes decoherence
- modification of the Gaussian noise



1. Numerical simulation of periodically pulsed quantum dot ensembles: understanding the synchronisation of the nuclear spin bath
2. Long-range spin interaction between quantum dots: coupling of subsystems pulsed with different laser colors
3. Proposal: higher-order spin correlation functions can reveal additional weak interactions



Thank you very much!