# RUHR-UNIVERSITÄT BOCHUM

# A numerical study of precession driven flow using the pseudo-spectral approach combined with an immersed boundary method

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#### **General Introduction**

### Goal:

- Numerical study of precession driven flow in a cylindrical cavity spinning around two axes  $\rightarrow$  development of C++ code SpedDyn
- ► Focus on flows suitable for magnetic dynamo action in the DRESDYN dynamo experiment
- Special attention on baffles with adjustable penetration depth at the cylinder end caps with the aim to make the precession forcing more

#### **Numerical Methods**

#### **Pseudo-Spectral Method:**

Compute spatial derivatives in Fourier and multiplications in real space

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- Speed advantage due to efficiency of the FFT
- Pressure Poisson equation can be directly solved in Fourier space
- Implies periodic boundaries

#### **Immersed Boundary Method**

Boundary represented by triangular mesh

efficient



Figure 1: Schematic drawing of the precessing cylinder of radius R with (fully extended) baffles at the end caps. It spins around the symmetry axis with the angular velocity  $\Omega_S$  and is simultaneously in orthogonal precession with angular velocity  $\Omega_P$ . The maximum penetration depth of the baffles is  $d_0 = 0.065 R$ . They have a width of w = 0.04 R and their inner and outer radius is  $r_i = 0.49 R$ and  $r_o = 0.645 R$ .

#### **Problem:**

$$\partial_{t} \mathbf{u} = \underbrace{\mathbf{u} \times (\nabla \times \mathbf{u})}_{\text{self advection}} + \underbrace{\frac{1}{Re} \Delta \mathbf{u}}_{\text{diffusion}} - \underbrace{\nabla p}_{\text{modified pressure}} + \underbrace{2\mathbf{u} \times \Omega}_{\text{Coriolis force}} + \underbrace{\mathbf{r} \times \dot{\Omega}}_{\text{Poincaré force}} \quad (1)$$

$$\nabla \cdot \mathbf{u} = 0 \quad (2) \quad \mathbf{u} = 0 \text{ (at the boundaries)} \quad (3)$$

#### **Analytical predictions:**

- ► No-slip boundary condition is enforced by linear interpolation between mesh and adjacent fluid nodes on Cartesian grid  $\rightarrow O([\Delta x]^2)$
- $\blacktriangleright$  Does not commute with pressure calculation  $\rightarrow$  pressure corrector method  $(\epsilon_u \sim \mathcal{O}(\Delta t^2), \epsilon_p \sim \mathcal{O}(\Delta t^{3/2}))$



Figure 3: Schematic drawing presenting the linear interpolation of the immersed boundary to the IB nodes in two dimensions for a convex surface (a) and a concave corner (b). Shown are the three different types of nodes: fluid (blue), solid (red) and IB (black). The boundary is represented by a thick black line.

Results



 $\blacktriangleright$  Inviscid fluid in steadily and rapidly rotating cylinder  $\rightarrow$  Kelvin modes

$$\mathbf{u}(r,\phi,z,t) = A_{mnk} \mathbf{R}_{mnk}(r) \exp[i(m\phi + n\pi z + \omega_{mnk} t)]$$
(4)

- ► Weak precession: Resonance of Poincaré force with mode (m, n, k) = (1, 1, 1)
- ► Higher precession makes forced mode break down due to non-linear interactions
- $\blacktriangleright$  Experimental finding of excitation of double roll structure (0, 2, k) around  $Po \sim 0.095$
- ► Kinematic dynamo simulations at *HZDR* show critical role of double roll for the onset of dynamo action.



Figure 6: Azimuthal and temporal averaged velocity field Figure 4: Contour plot of the time averaged axial velocity  $V_z$  in the in the turntable frame for extended baffles. mantle frame of the cylinder. The presented flow is dominated by the double roll structure.



Figure 5: The non-axialsymmetric energy as a function of Po in the vicinity of the critical Poincaré number  $Po_F^c$  for different baffle lengths d.



Figure 7: Critical Poincaré numbers of the non-axialsymmetric energy  $(Po_F^c)$ , the double roll structure  $(Po_{DF}^c)$  and the turbulent energy  $(Po_{II}^{c})$  as functions of the baffle length.

Computational domain is split into pencils  $\rightarrow$ one dimension is completely available on each process

Processor grid is transposed using *MPI\_Datatypes* and the API function MPI\_Alltoallw



Figure 2: Scaling of *SpecDyn* for 100 time steps on a grid of size  $512^3$  and  $1024^3$ . The computation was performed on the JUWELS CPU cluster with 48 CPUs per node. The speed up is normalized to 3 nodes for  $N = 512^3$  and to 12 nodes for  $N = 1024^3$ .

#### Outlook

- Dynamical MHD simulations already in progress
- Determination of critical magnetic Reynolds numbers for different Poincaré numbers and of the general magnetic field structure.

#### **Publications**

[1] M. Wilbert, A. Giesecke, and R. Grauer.

Numerical investigation of the flow inside a precession driven cylindrical cavity with additional baffles using an immersed boundary method. POF, in press, 2022.



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