



PSEUDO MAJORANA FUNCTIONAL RENORMALISATION GROUP FOR QUANTUM SPIN SYSTEMS

Johannes Reuther

Freie Universität Berlin, Helmholtz Zentrum Berlin

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The group



Johannes Reuther, Tim Bauer, Cecilie Glittum, Matias Gonzalez, Noah Hassan, Anna Fancelli, Yannik Schaden, Julius Tietz, Sarah Simon, Sanne Kleinheinz

Introduction

(Spin liquids and their excitations)


Ordered states of matter

Usually, matter “freezes” at sufficiently low temperatures
 \Rightarrow it form a more ordered state

Examples:

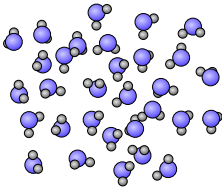


water

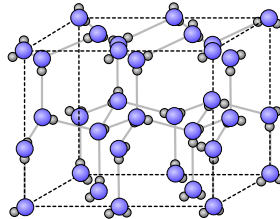
$T < 0^{\circ}\text{C}$ 



ice



disordered liquid



regular crystal

pictures: pixabay.com

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 \Rightarrow it form a more ordered state

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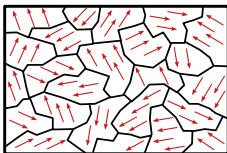


iron, paramagnetic

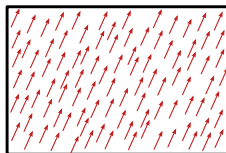
$T < 768^\circ\text{C}$
 \longrightarrow



iron, ferromagnetic



many domains



single domain

pictures: pixabay.com

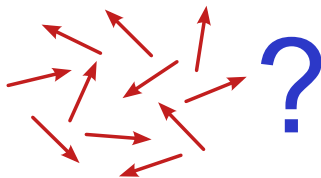
Spin liquids

Our research interest:

Magnetic states of matter and materials which remain **liquid** (disordered, dynamic) down to $T = 0K$.

Spin liquids

⇒ Special mechanism to keep the spins disordered are needed.



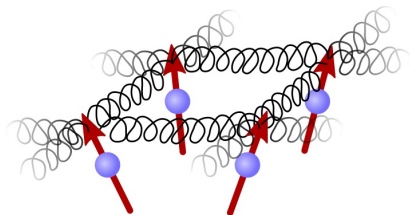
Basics: Microscopic situation

Consider interacting electron spins (spin-1/2) on a lattice.

Simplest (and most realistic) type of interaction:

Heisenberg couplings (spin isotropic)

$$H = \sum_{ij} J_{ij} \vec{S}_i \vec{S}_j$$

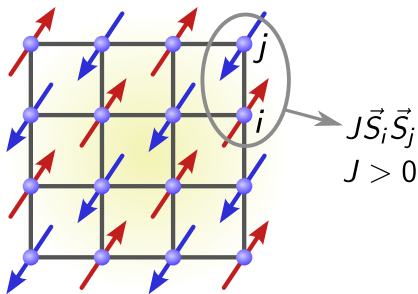


Basics: Antiferromagnetic interactions

Heisenberg couplings naturally arise in **insulating** materials with **strong electron interactions** (Mott insulator).

⇒ Electrons cannot move ⇒ **spin** degree of freedom important!

⇒ Interactions are **antiferromagnetic**, i.e. spins want to align antiparallel.



Simplest case:

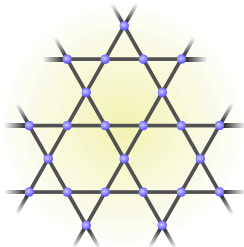
Spin-1/2 Heisenberg antiferromagnet on **square lattice**

$$H = \sum_{\langle ij \rangle} J \vec{S}_i \vec{S}_j$$

Magnetic order: **Néel state**

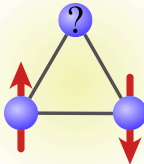
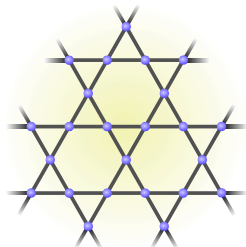
Magnetic frustration

What happens for antiferromagnetic **triangles**? (e.g. Kagome lattice)



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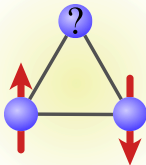
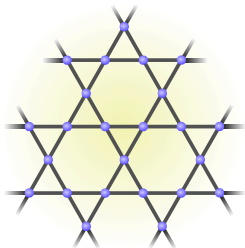


⇒ **Magnetic frustration**

Terms in H cannot all be simultaneously minimized!

Magnetic frustration

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⇒ **Magnetic frustration**

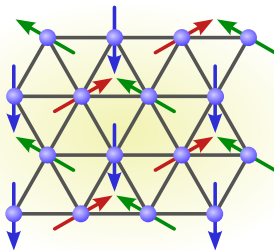
Terms in H cannot all be simultaneously minimized!

Possibility 1:

Spins find a compromise.

Example:

120° Néel order on triangular lattice.



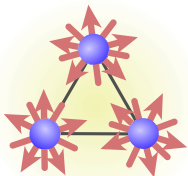
Many other interesting types of orders exist (e.g. spin spirals).

Magnetic frustration

Possibility 2:

Use properties from [quantum mechanics](#):

Spin don't have to point in any specific direction. The directions can [fluctuate](#).



⇒ **Novel spin state:**

Quantum spin liquid

No magnetic order: $\langle \Psi | \vec{S}_i | \Psi \rangle = 0$

Quantum mechanical description?

Spin singlets

Simplest spin state with $\langle \Psi | \vec{S}_i | \Psi \rangle = 0$?

Two spins in a **singlet** $|\Psi\rangle = \frac{1}{\sqrt{2}} (|\uparrow_1\downarrow_2\rangle - |\downarrow_1\uparrow_2\rangle)$

Is also **ground state** of $H = J\vec{S}_1\vec{S}_2$ for $J > 0$



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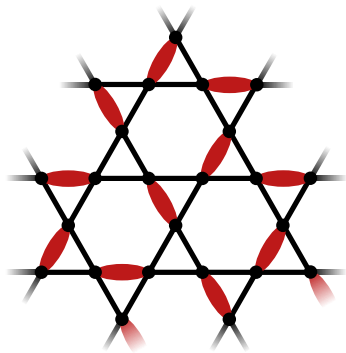
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Pair up spins on the lattice into singlet bonds:



Non-magnetic state



But still **static**!



Resonating valence bond state

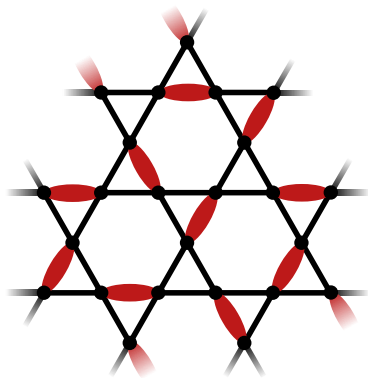
Solution: Superimpose macroscopically many singlet configurations!

$$|\psi\rangle = a_1 |\text{Diagram 1}\rangle + a_2 |\text{Diagram 2}\rangle + a_3 |\text{Diagram 3}\rangle + \dots$$
The equation shows three diagrams of a triangular lattice, each representing a different singlet configuration. In each diagram, a subset of the bonds is highlighted in red, forming a pattern of singlets (pairs of sites connected by a red bond). The first diagram shows a specific arrangement of red bonds, the second shows a different arrangement, and the third shows a third arrangement. The diagrams are superimposed to illustrate the concept of a resonating valence bond state.

= resonating valence bond (RVB) state. (Anderson, Baskaran, 1987)

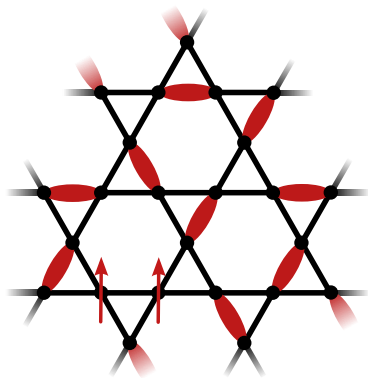
Magnetic excitations: Spinons

Break up a singlet dimer into two aligned spins (= spin-1 excitation):



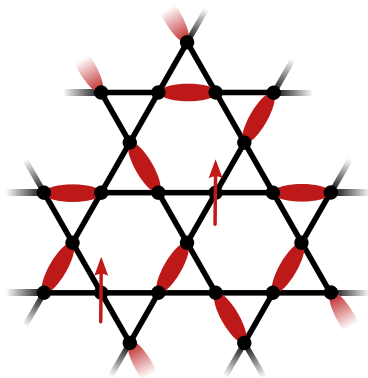
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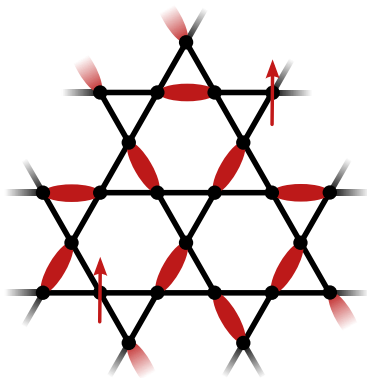
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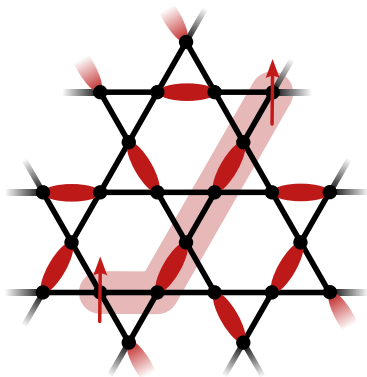
Isolated spins are effectively free (**deconfined**).

⇒ New quasiparticles: **spinons**. Carry spin $\frac{1}{2}$!

Fractional excitations

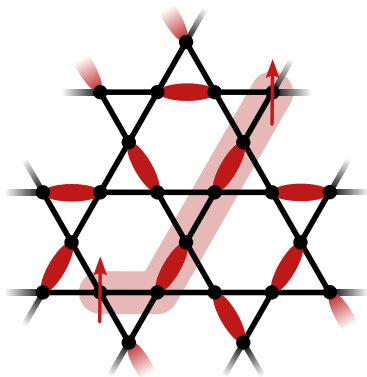
Rearranged dimer background

Fractional spinons are connected by a **string** of rearranged dimers:



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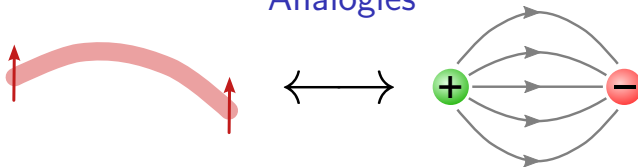


Basic ingredients for a spin liquid at low energies:

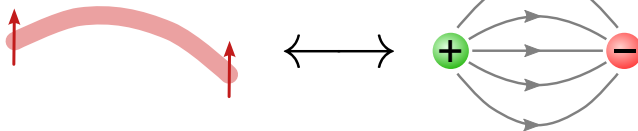
- Spinons.
- Strings connecting them.



Analogies



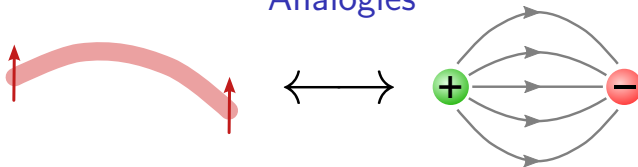
Analogies



Spinons	\longleftrightarrow	Charged particles (“matter fields”, e.g. electrons)
Strings	\longleftrightarrow	Field lines (“gauge fields”, e.g. electric field)

Spin liquid \longleftrightarrow **Theories from fundamental physics**

Analogies



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(“matter fields”, e.g. electrons)

Strings \longleftrightarrow Field lines
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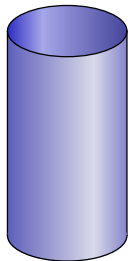
Spin liquid \longleftrightarrow Theories from fundamental physics

A spin liquid realizes theories otherwise only known from fundamental physics \rightarrow playground for finding new physics in exotic states of matter.

- Emergent electrodynamics (spin ice materials).
- Majorana fermions (half electrons) in Kitaev spin liquids.
- Effective gravity in spin-1 magnets.

Topological properties of spin liquids

Fractional excitations and topological ground state degeneracy:

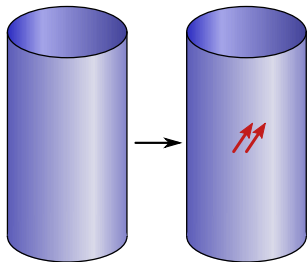


$|1\rangle$

spin liquid in ground-state $|1\rangle$

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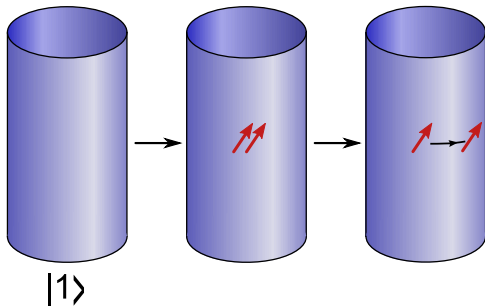
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→ create pair of spinons

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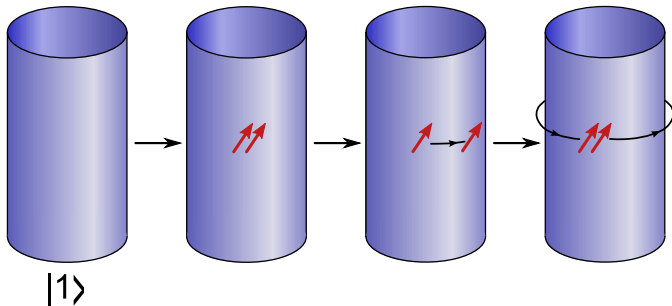
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→ create pair of spinons

→ wind one spinon around cylinder

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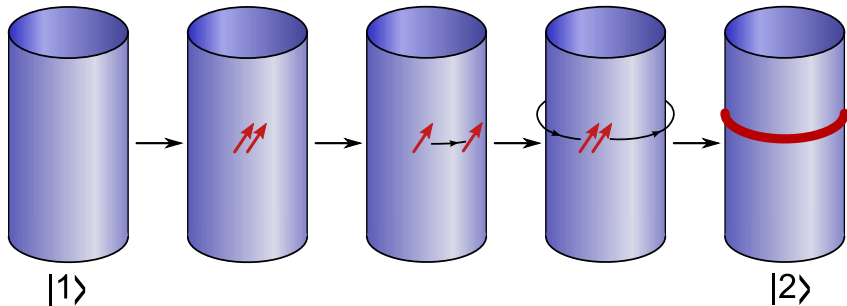
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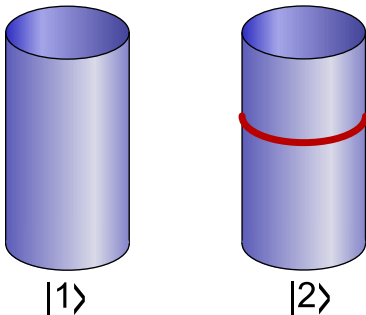
→ **wind** one spinon around cylinder

→ **annihilate** spinons, end up in ground state $|2\rangle \neq |1\rangle$

Possible quantum information applications

Transition from $|1\rangle$ to $|2\rangle$ requires **non-local operation**.

→ **quantum information is topologically protected!**



Idea: Use the states $|1\rangle$ and $|2\rangle$ as **stable qubit states**.

Quantum spin liquids in experiments

Spinons can only be created in **pairs**!

Do **spin-1 excitation** with energy E and momentum q :

- **Decomposition** into spinon 1 with energy ϵ_1 and momentum k_1
and spinon 2 with energy ϵ_2 and momentum k_2 .
- **Continuum** of possibilities such that $E = \epsilon_1 + \epsilon_2$ and $q = k_1 + k_2$

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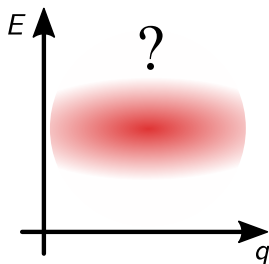
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Excitation spectrum:

Quantum spin liquid



two-spinon continuum

Quantum spin liquids in experiments

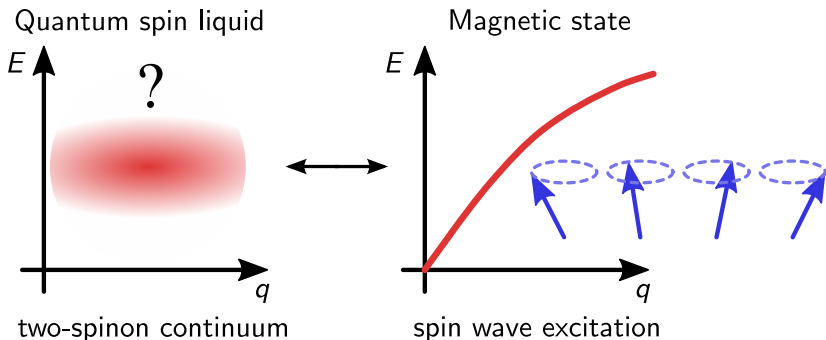
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Excitation spectrum:



Numerical methods: PFFRG and PMFRG

Review article: T. Müller, D. Kiese, N. Niggemann, B. Sbierski, JR, S. Trebst, R. Thomale, and Y. Iqbal, *Pseudo-fermion functional renormalization group for spin models*, Rep. Prog. Phys. 87, 036501 (2024)

Pseudo fermions

Starting point: $H = \sum_{ij} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$

Express spin operators in terms of pseudo fermions: $\mathbf{S} = \frac{1}{2} \sum_{\alpha\beta \in \uparrow, \downarrow} f_{\alpha}^{\dagger} \boldsymbol{\sigma}_{\alpha\beta} f_{\beta}$
where $\{f_{\alpha}, f_{\beta}^{\dagger}\} = \delta_{\alpha\beta}$.

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Be careful: Enlarged Hilbert space with unphysical non- and doubly occupied states.

Experience: Unphysical states are excitations. Application at $T = 0$.

Exceptions: B. Schneider, D. Kiese, and B. Sbierski, PRB 106, 235113 (2022)

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$$H = \sum_{ij} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j \longrightarrow \frac{1}{4} \sum_{\alpha\beta\gamma\delta} \sum_{ij} J_{ij} \left(f_{i\alpha}^{\dagger} \boldsymbol{\sigma}_{\alpha\beta} f_{i\beta} \right) \cdot \left(f_{j\gamma}^{\dagger} \boldsymbol{\sigma}_{\gamma\delta} f_{j\delta} \right)$$

Diagrammatics in the fermions:

Progagator: $G_0(i\omega) = \frac{1}{i\omega} = \text{---}\blacktriangleleft$ Interaction vertex: $\Gamma_0 = \text{---}\blacktriangleright \cdots \blacktriangleright \sim J$

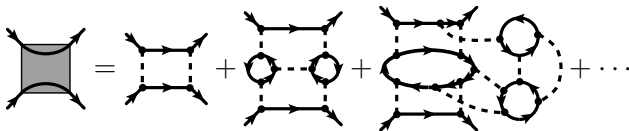
Functional renormalization group

Magnetic susceptibility, spin-spin correlations, spin structure factor:



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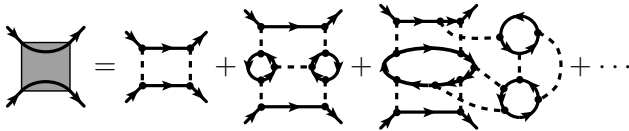
Sampling problem for Feynman diagrams! Efficient resummation scheme:

Functional renormalization group (FRG)

Metzner, Salmhofer, Honerkamp, Meden, Schönhammer, RMP 84, 299 (2012)

Functional renormalization group

Magnetic susceptibility, spin-spin correlations, spin structure factor:



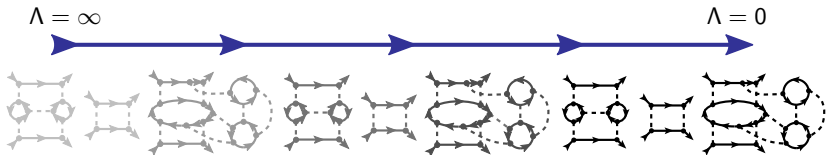
A diagrammatic equation representing the magnetic susceptibility. On the left is a shaded square with four external lines. This is equal to a series of diagrams: a square with dashed internal lines, a square with two internal loops, a square with two internal loops and a central dashed line, and a square with two internal loops and a central dashed line with a bubble, followed by an ellipsis.

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Metzner, Salmhofer, Honerkamp, Meden, Schönhammer, RMP 84, 299 (2012)

Sum up Feynman diagrams during the flow of an RG parameter Λ :



Functional renormalization group

Example: Implementation of Λ as a low frequency cutoff in $G_0(i\omega)$:

$$G_0(i\omega) = \frac{1}{i\omega} \longrightarrow \frac{\Theta(|\omega| - \Lambda)}{i\omega}$$

Functional renormalization group

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$$G_0(i\omega) = \frac{1}{i\omega} \longrightarrow \frac{\Theta(|\omega| - \Lambda)}{i\omega}$$

FRG formulates **differential equations for all m -particle vertex functions**:

$$\frac{d}{d\Lambda} \text{ (2-point vertex) } = - \text{ (2-point vertex with self-energy loop) }$$

$$\frac{d}{d\Lambda} \text{ (4-point vertex) } = \text{ (4-point vertex with two internal 2-point vertices) } + \text{ (4-point vertex with two internal 4-point vertices) } - \text{ (4-point vertex with two internal 2-point vertices and a loop) } + \text{ (4-point vertex with two internal 2-point vertices and a loop) } + \text{ (4-point vertex with two internal 2-point vertices and a loop) } + \text{ (4-point vertex with two internal 2-point vertices and a loop) }$$

$$\frac{d}{d\Lambda} \text{ (6-point vertex) } = \dots$$

Functional renormalization group

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FRG formulates **differential equations for all m -particle vertex functions**:

$$\frac{d}{d\Lambda} \text{ (circle with dot) } = - \text{ (cylinder) }$$

$$\frac{d}{d\Lambda} \text{ (square) } = \text{ (two squares) } + \text{ (two squares) } - \text{ (cylinder) } + \text{ (cylinder) } + \text{ (cylinder) } + \text{ (hexagon with X) }$$

$$\frac{d}{d\Lambda} \text{ (hexagon with X) } = \dots$$

Truncation needed! A. A. Katanin, PRB **70**, 115109 (2004)

Pseudofermion functional renormalization group (PFFRG)

J. Reuther and P. Wölfle, PRB **81**, 144410 (2010)

Pseudo Majorana functional renormalization group

Alternatively: Express spin operators in terms of **Majorana fermions**:

$$S^\mu = -\frac{i}{4} \sum_{\nu\sigma \in x,y,z} \epsilon_{\mu\nu\sigma} c^\nu c^\sigma \text{ with } (c^\mu)^2 = 1.$$

Advantage: **No unphysical states, only redundant Hilbert spaces!**
Applicable at $T > 0$.

Pseudo Majorana functional renormalization group (PMFRG)

N. Niggemann, B. Sbierski, and JR, PRB 103, 104431 (2021)

N. Niggemann, JR, B. Sbierski, SciPost Phys. 12, 156 (2022)

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Higher efficiency with temperature-flow approach where $\Lambda = T$!

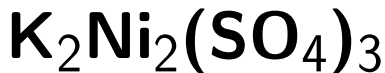
C. Honerkamp and M. Salmhofer, PRB 64, 184516 (2001)

Facts about the code

- Main challenge: Solving large systems of coupled ODEs but does not require much memory.
- Benefits greatly from parallelization (openMP and MPI parallelized).
- Language: Julia (older code in C++), public package:
<https://github.com/NilsNiggemann/PMFRG.jl>
- ~ 15 years of code development.
- Possibility to generate lattices automatically from VESTA files → easy to work with complicated models.

Application example:

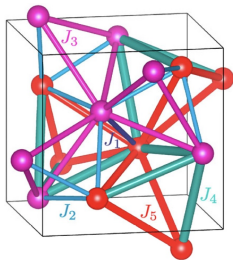
Spin liquid candidate



M. Gonzalez, JR et al., Nature Communications 15, 7191 (2024)

Lattice structure of $\text{K}_2\text{Ni}_2(\text{SO}_4)_3$

$\text{K}_2\text{Ni}_2(\text{SO}_4)_3$ has spin-1, Ni^{2+} , two interpenetrating trillium lattices!



● Ni^{2+}

● Ni^{2+}

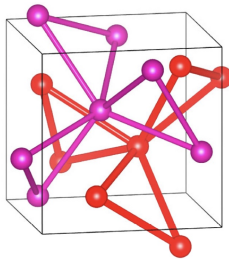
$J_1 = 0.364(2) \text{ K}$

$J_2 = -0.144(1) \text{ K}$

$J_3 = 0.798(2) \text{ K}$

$J_4 = 5.545(1) \text{ K}$

$J_5 = 2.657(1) \text{ K}$



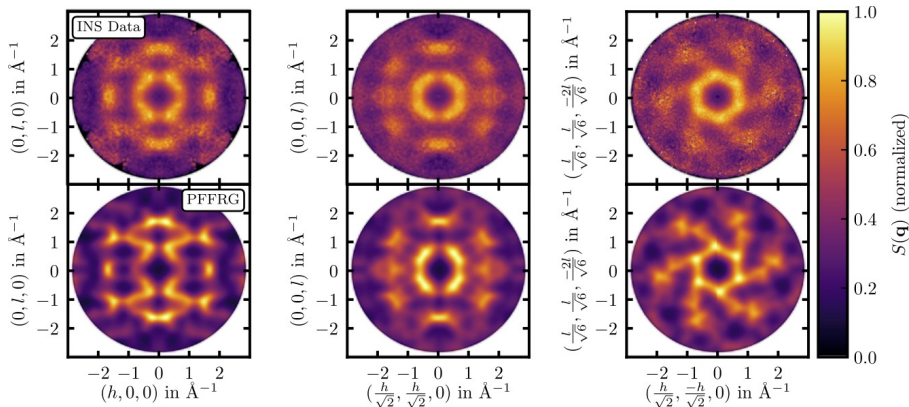
Interactions J_1, J_2, J_3, J_4, J_5 are strongest.

J_1, J_2, J_4 : Inter-trillium couplings, J_4 largest

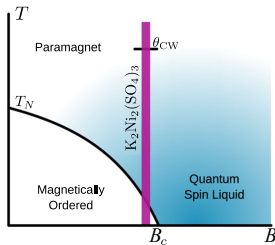
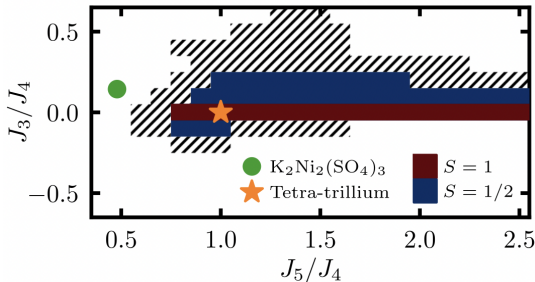
J_3, J_5 : Intra-trillium couplings

$\text{K}_2\text{Ni}_2(\text{SO}_4)_3$: Theory versus experiment

Spin structure factor $\mathcal{S}(\mathbf{q})$ from PFFRG and from neutron scattering:



$K_2Ni_2(SO_4)_3$: Phase diagram



- Phase diagram contains ‘island of liquidity’ close to $K_2Ni_2(SO_4)_3$.
- Model of $K_2Ni_2(SO_4)_3$ shows weak magnetic order (also seen in experiments).
- Small modifications, e.g. by magnetic field can induce spin liquid.

Type of spin liquid? To be continued....

Conclusions

- Spin liquids are exotic states of matter realizing emergent phenomena, e.g. theories from fundamental physics.
- The models describing spin liquids and possible candidate materials are among the hardest in physics research.
- PFFRG and PMFRG are methods that are capable of dealing with these complicated models.
- The methods are ideal for simulating magnetic quantum materials, and the results are in good agreement with measurements.

**Thank you for
your attention!**

